

# Neutrino-Driven Turbulent Convection in Stalled Supernova Cores

David Radice



Collaborators: E. Abdiakamalov, S. Couch, R. Haas, C. Ott, E. Schnetter

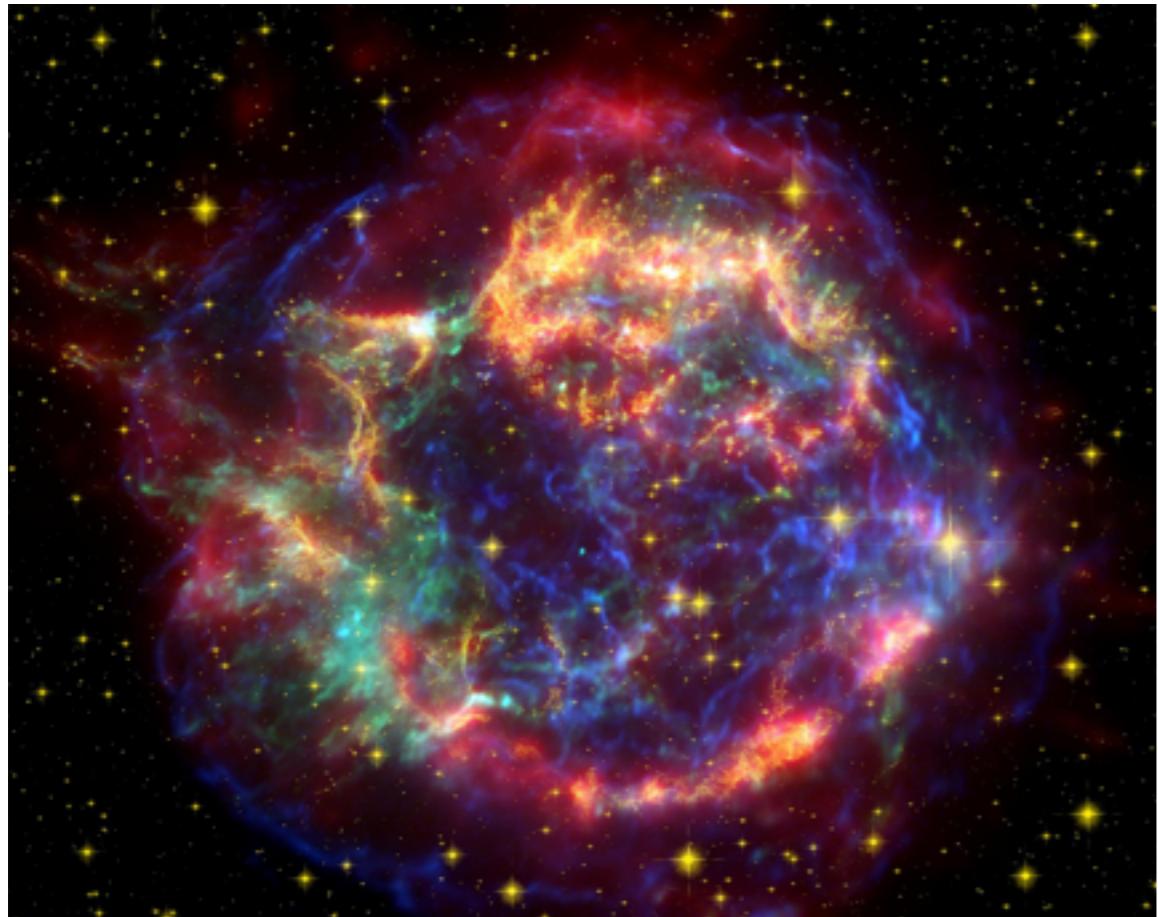
# Contents

- 1.Turbulence in core-collapse supernovae
- 2.Numerical simulations
- 3.Conclusions

# Contents

- 1.Turbulence in core-collapse supernovae**
- 2.Numerical simulations
- 3.Conclusions

# The Supernova Problem



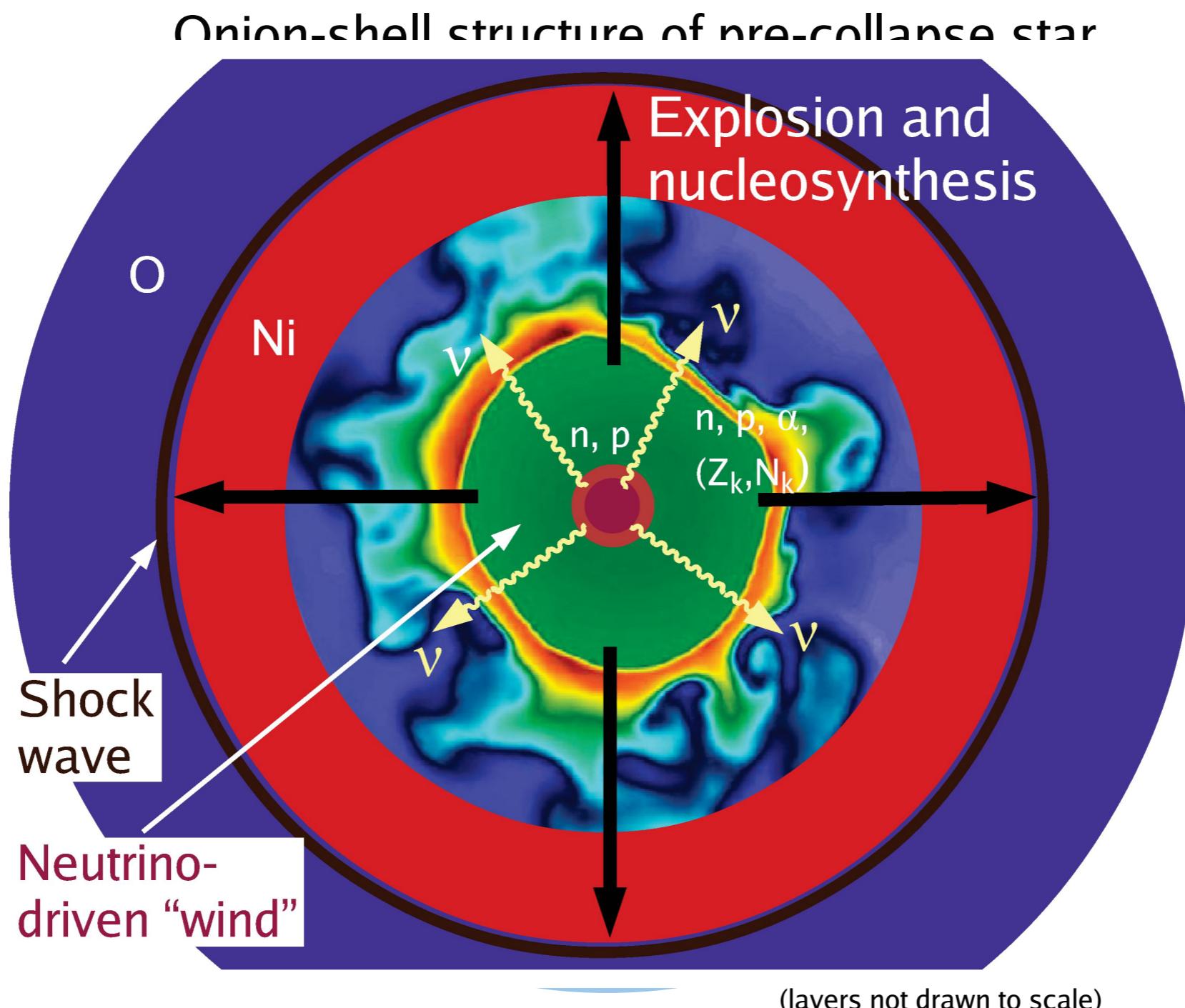
Cassiopeia-A

Core-Collapse Supernovae:

- End of massive stars
- Birthplace of heavy elements, neutron stars, black holes ...
- Regulate star formation
- ...

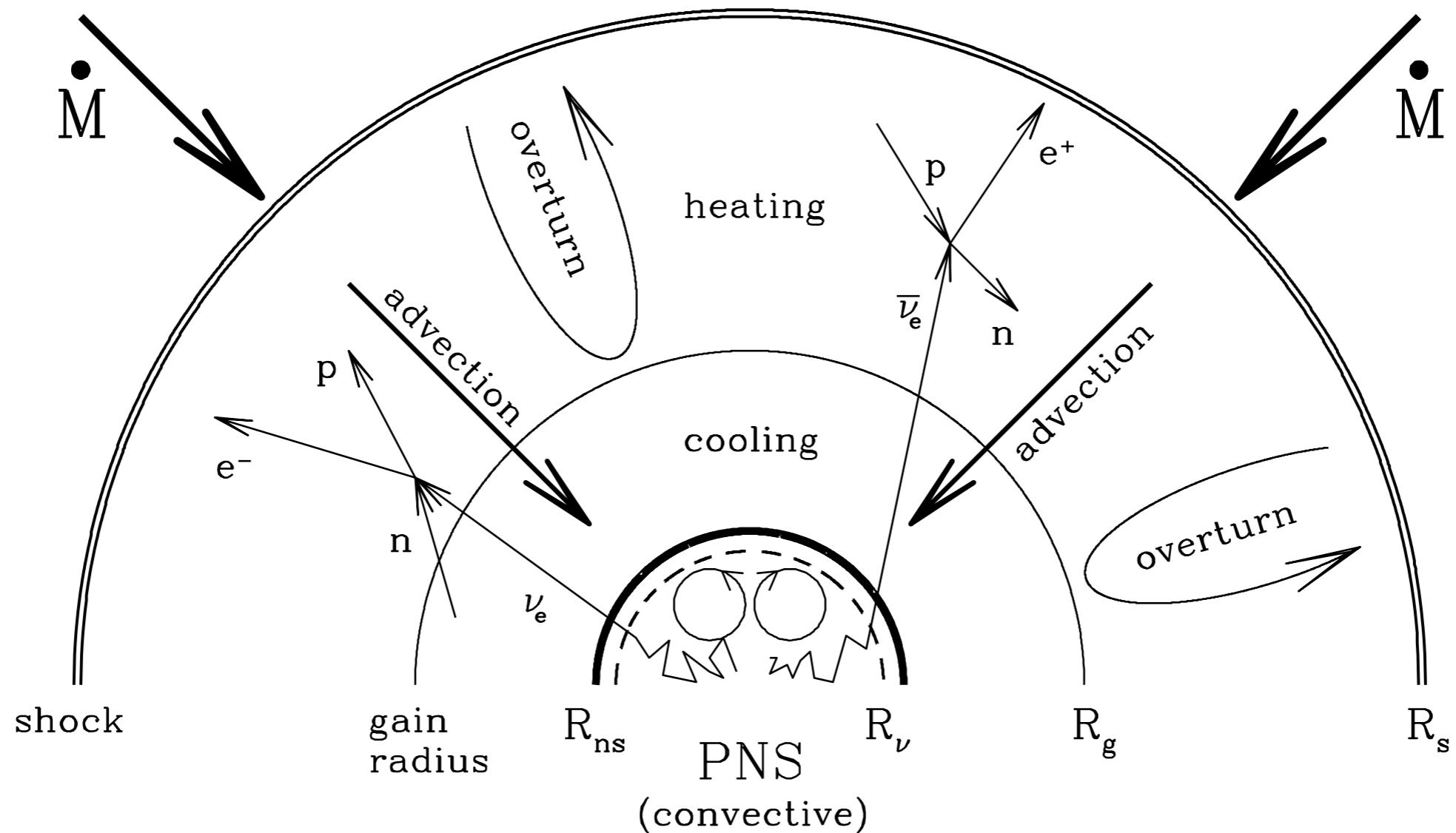
Problem: how do they explode?

# Core-Collapse Supernovae



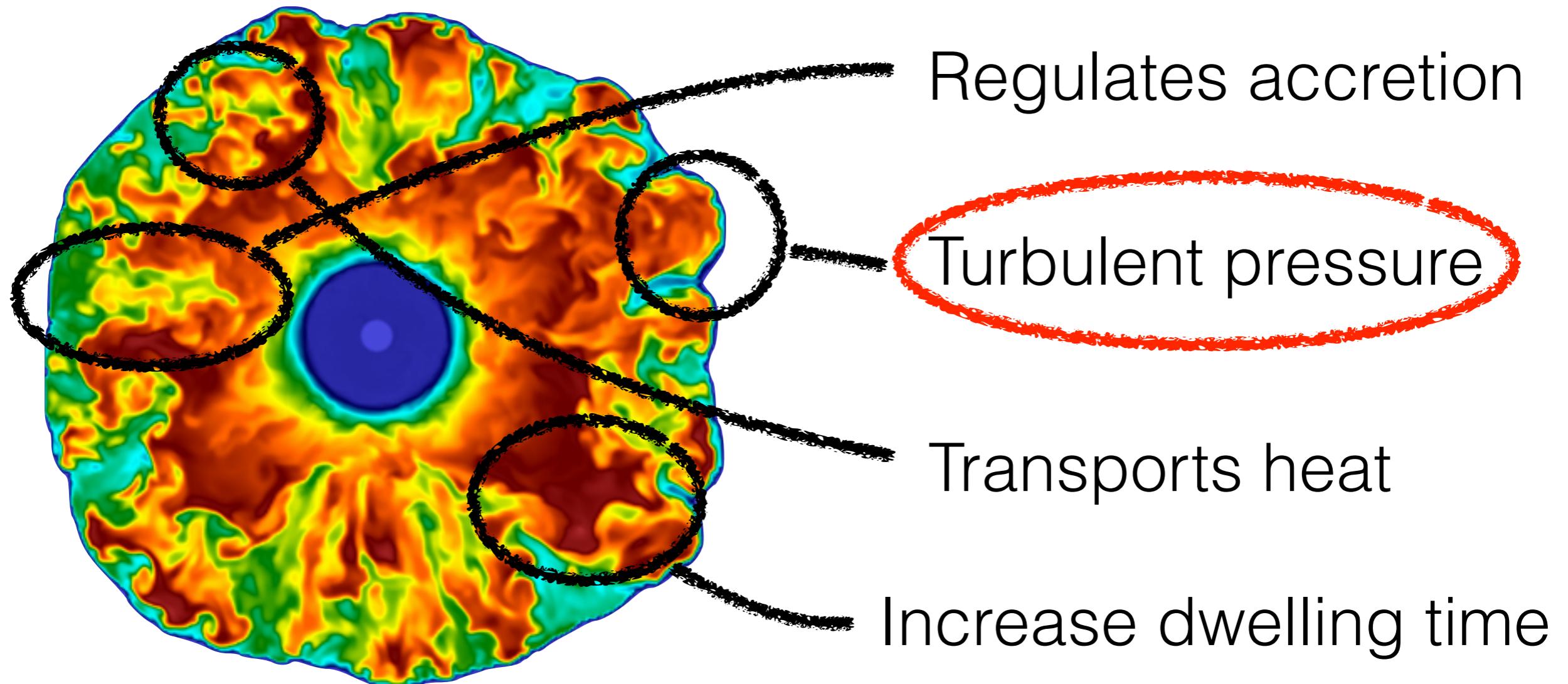
From Janka et al. 2012

# Shock Revival by Neutrinos



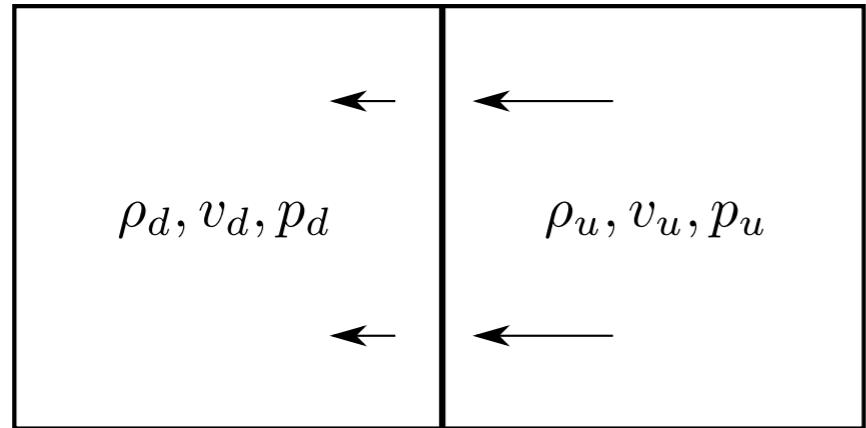
From Janka 2001

# The Roles of Turbulence



Difficult to simulate!

# Turbulent Pressure



Rankine-Hogoniot jump condition:

$$\rho_d v_d^2 + p_d = \rho_u v_u^2 + p_u$$

EOS:

$$p_d = (\gamma_{\text{th}} - 1)\rho_d \epsilon_d \quad \gamma_{\text{th}} \simeq \frac{4}{3}$$

Effect of downstream turbulence (Murphy et al. 2013):

$$\rho_d v_d^2 + p_d \rightarrow \rho_d \bar{v}_d^2 + \rho_d (\delta v)_d^2 + p_d$$

$\gamma_{\text{turb}} > \gamma_{\text{th}}$ !

Turbulence can be modeled with an effective EOS

$$\rho_d (\delta v)_d^2 \leftrightarrow (\gamma_{\text{turb}} - 1)\rho_d \epsilon_{\text{turb}} \quad \gamma_{\text{turb}} \simeq 2$$

Jump conditions for a shock with downstream turbulence:

$$\rho_d \bar{v}_d^2 + (\gamma_{\text{turb}} - 1)\rho_d \epsilon_{\text{turb}} + (\gamma_{\text{th}} - 1)\rho_d \epsilon_d = \rho_u v_u^2 + p_u$$

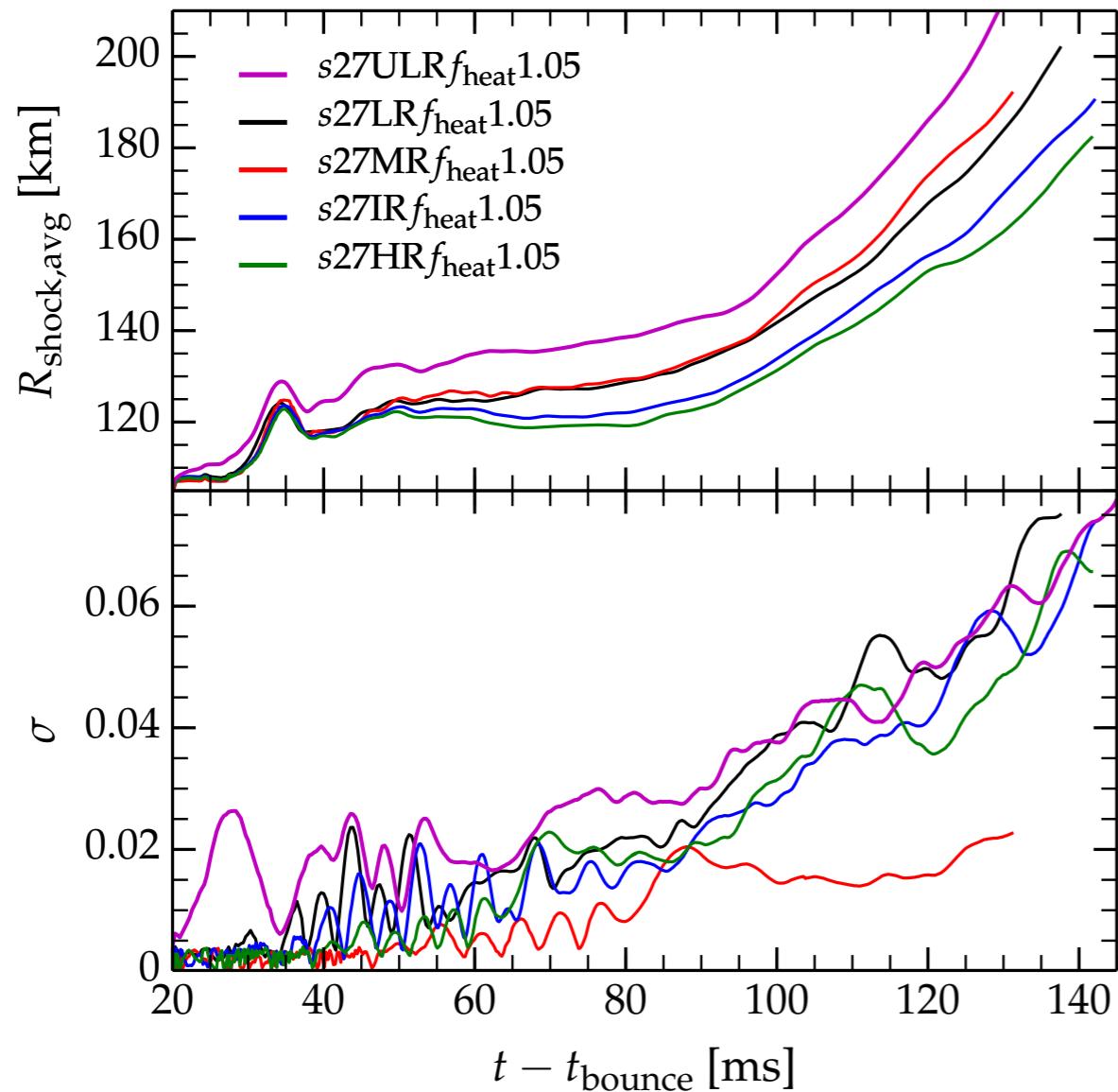
# Contents

1.Turbulence in core-collapse supernovae

**2.Numerical simulations**

3.Conclusions

# Resolution Dependence



ULR	3.78 km
LR	1.89 km
MR	1.42 km
IR	1.24 km
HR	1.06 km

Resolutions

Explosion more difficult at higher resolution!

# Why?

- Lower resolution favors the formation of larger, longer lived structures
- Secondary instabilities (Kelvin-Helmholtz) is suppressed by numerical viscosity
- When is the resolution good enough?

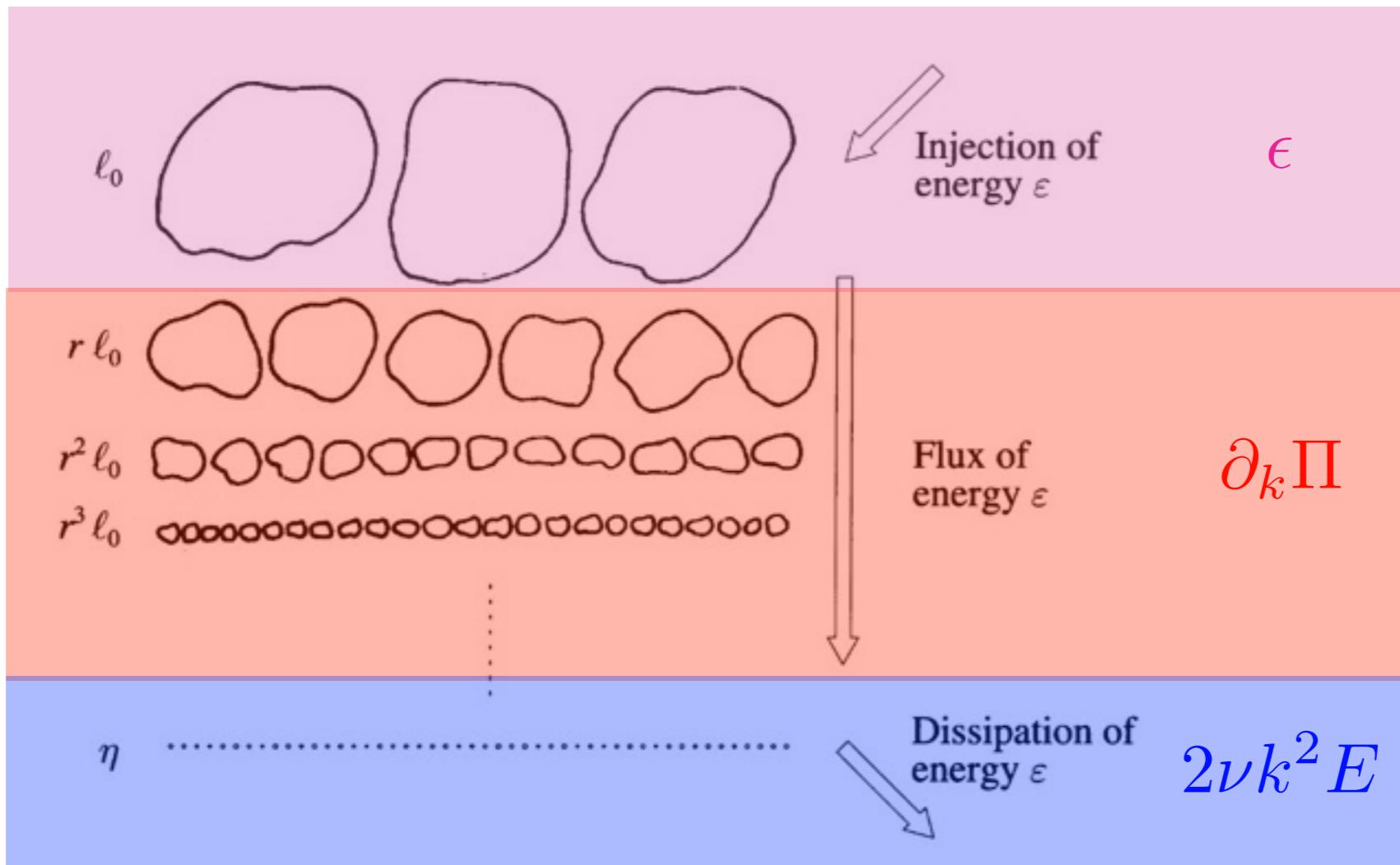
# Turbulent Cascade I

$$\partial_t E + \partial_k \Pi = -2\nu k^2 E + \epsilon$$

Diagram illustrating the energy balance equation for a turbulent cascade:

- Energy flux** (red arrow pointing down):  $\partial_k \Pi$
- Energy injection** (pink arrow pointing down):  $\epsilon$
- Specific kinetic energy** (purple arrow pointing up):  $E$
- Energy dissipation** (blue arrow pointing up):  $-2\nu k^2 E$

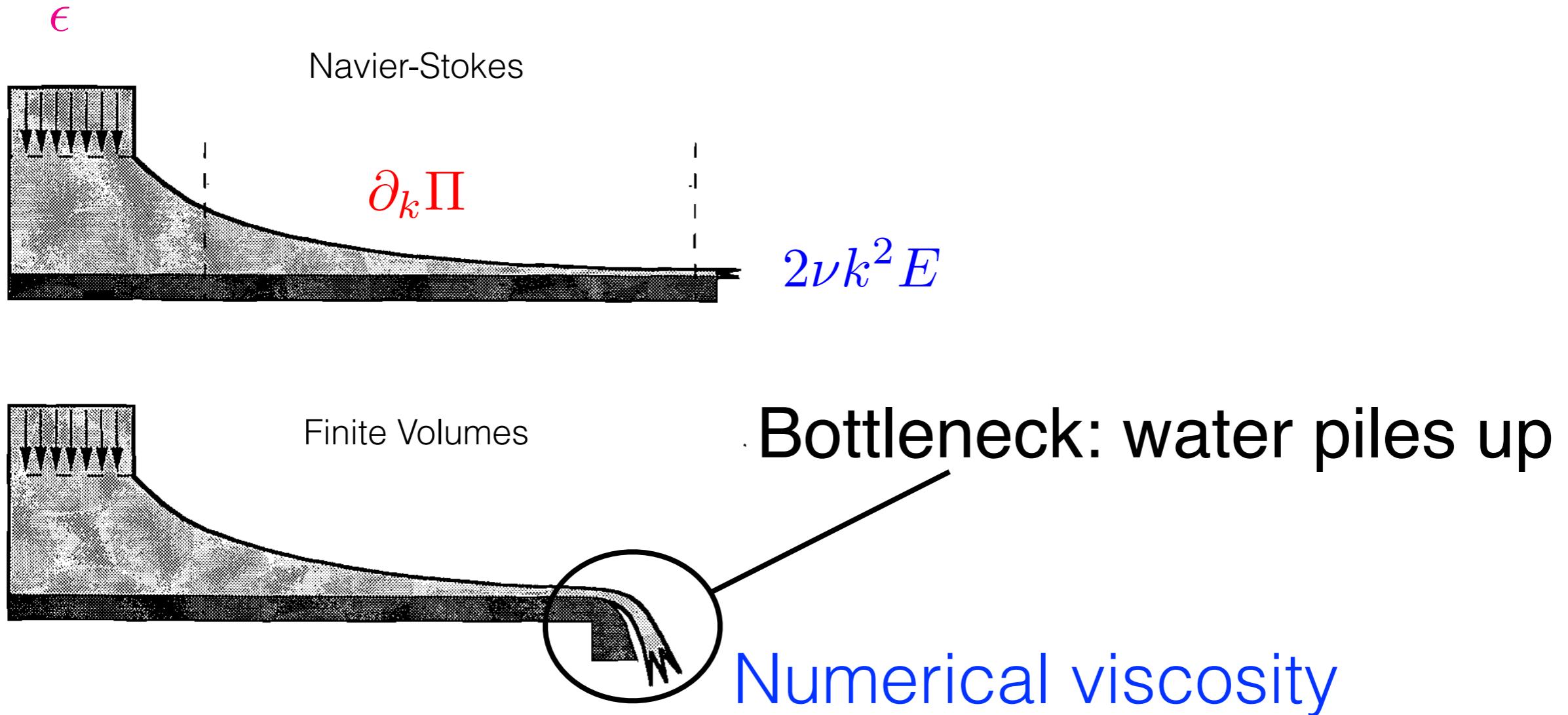
# Turbulent Cascade II



Adapted from Frisch 1996

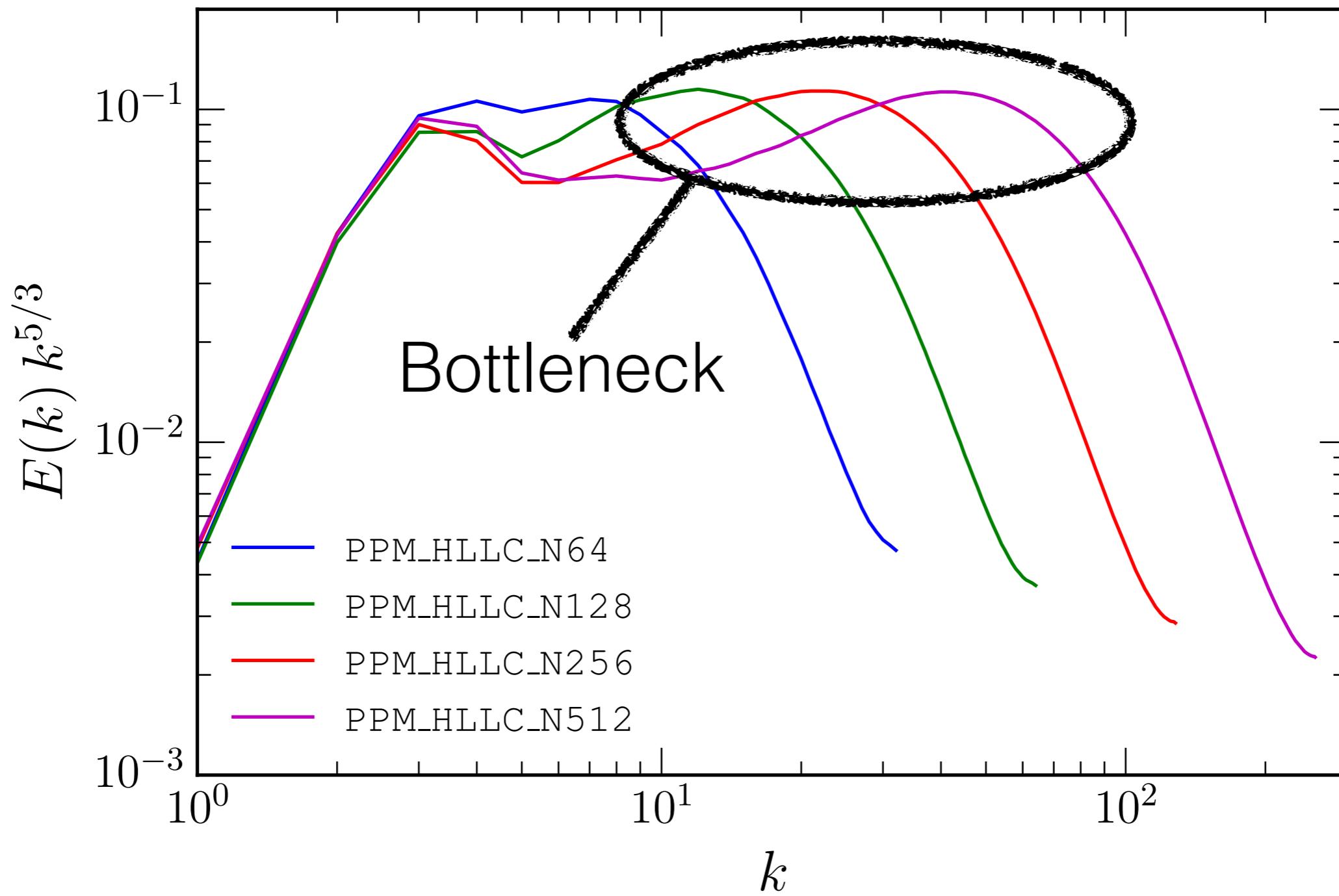
Kolmogorov 1941:  $\Pi \simeq \text{const}$   $\implies E \sim k^{-5/3}$

# The Water-Spill Analogy

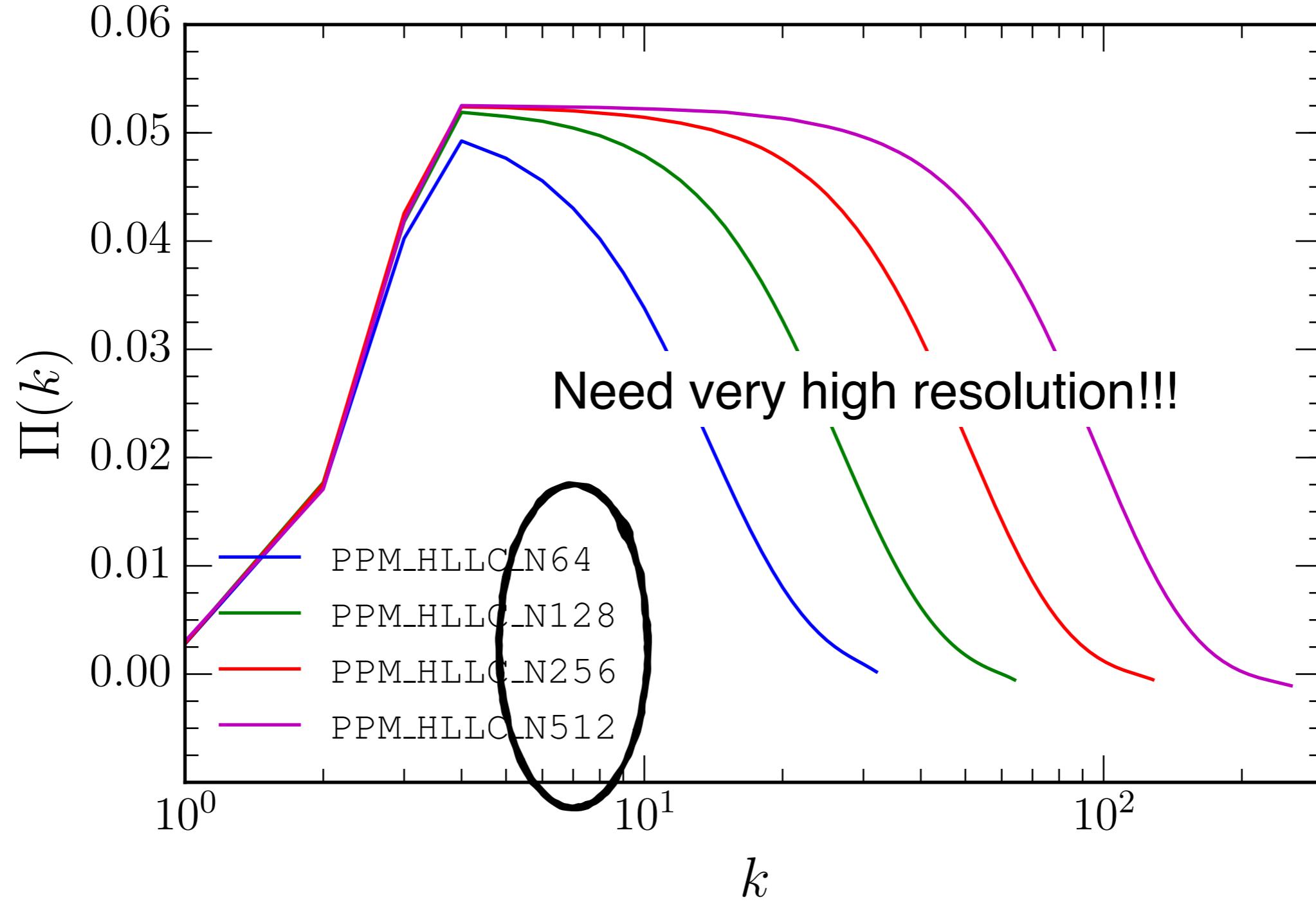


Adapted from Boris 1992

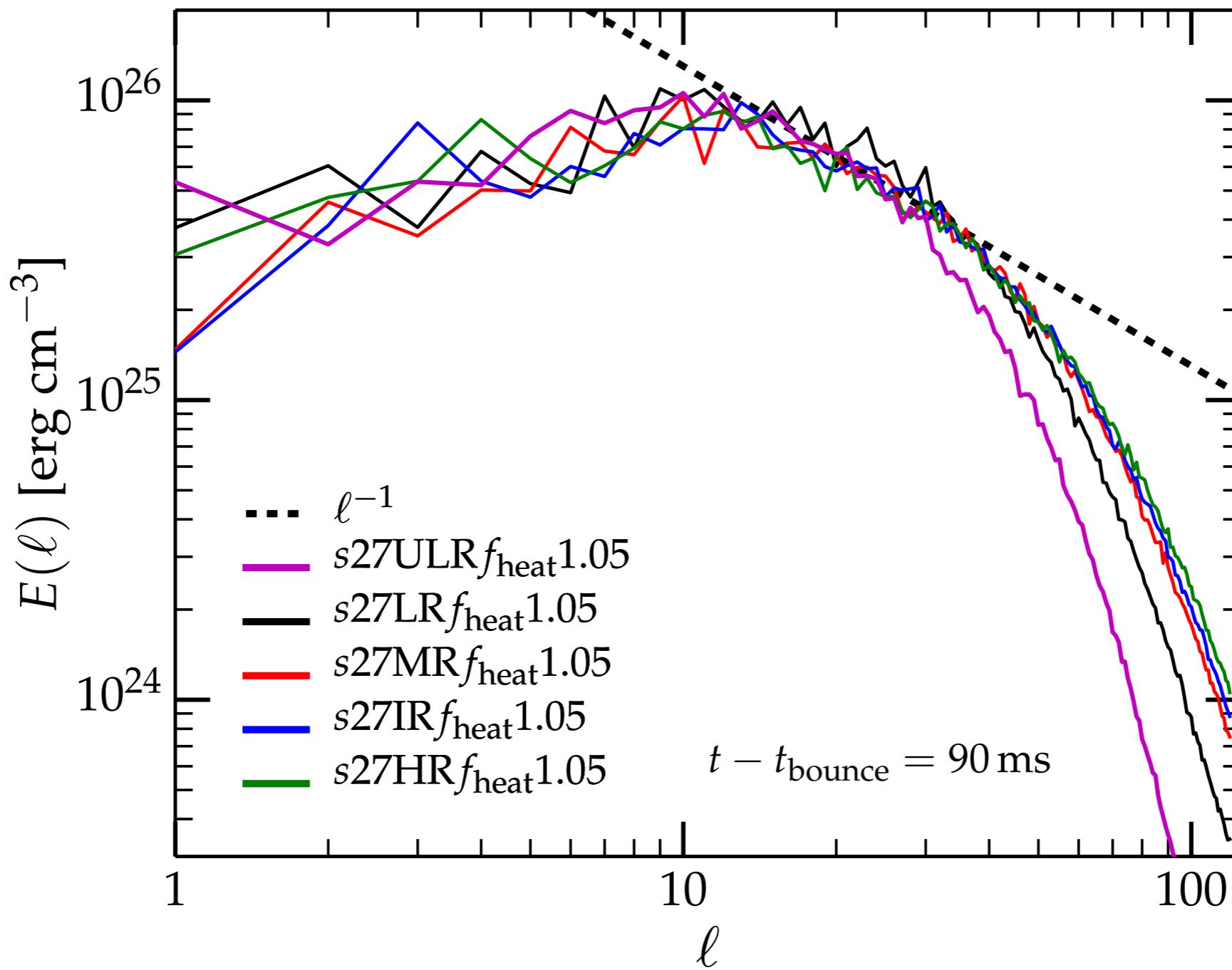
# The Bottleneck Effect



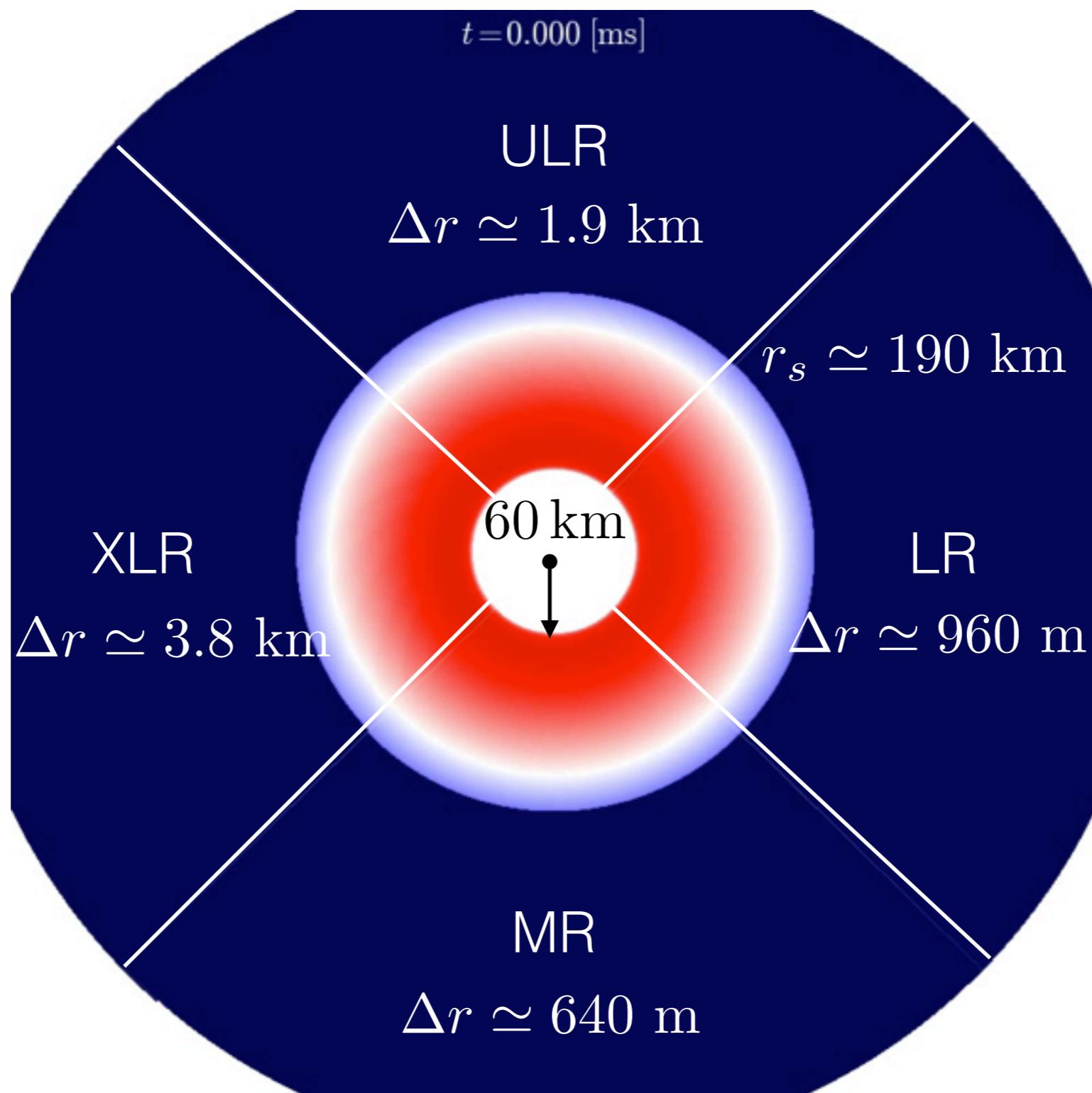
# Energy Cascade: PPM

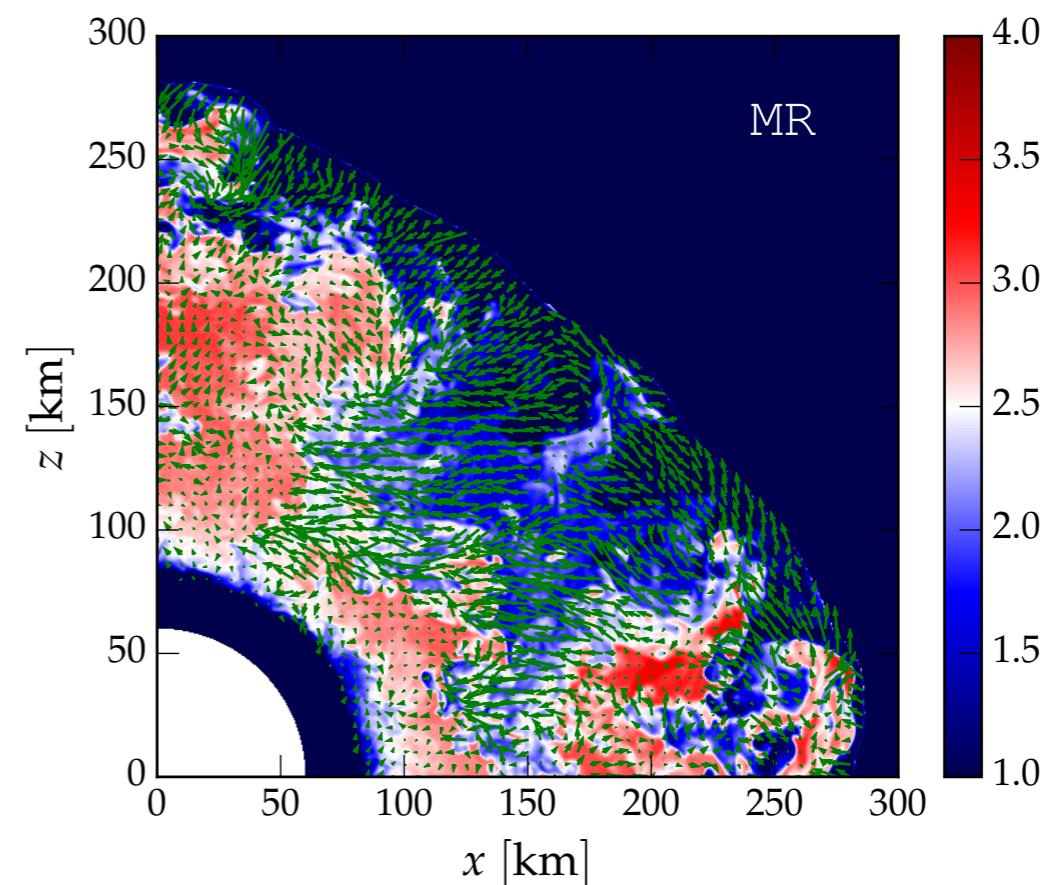
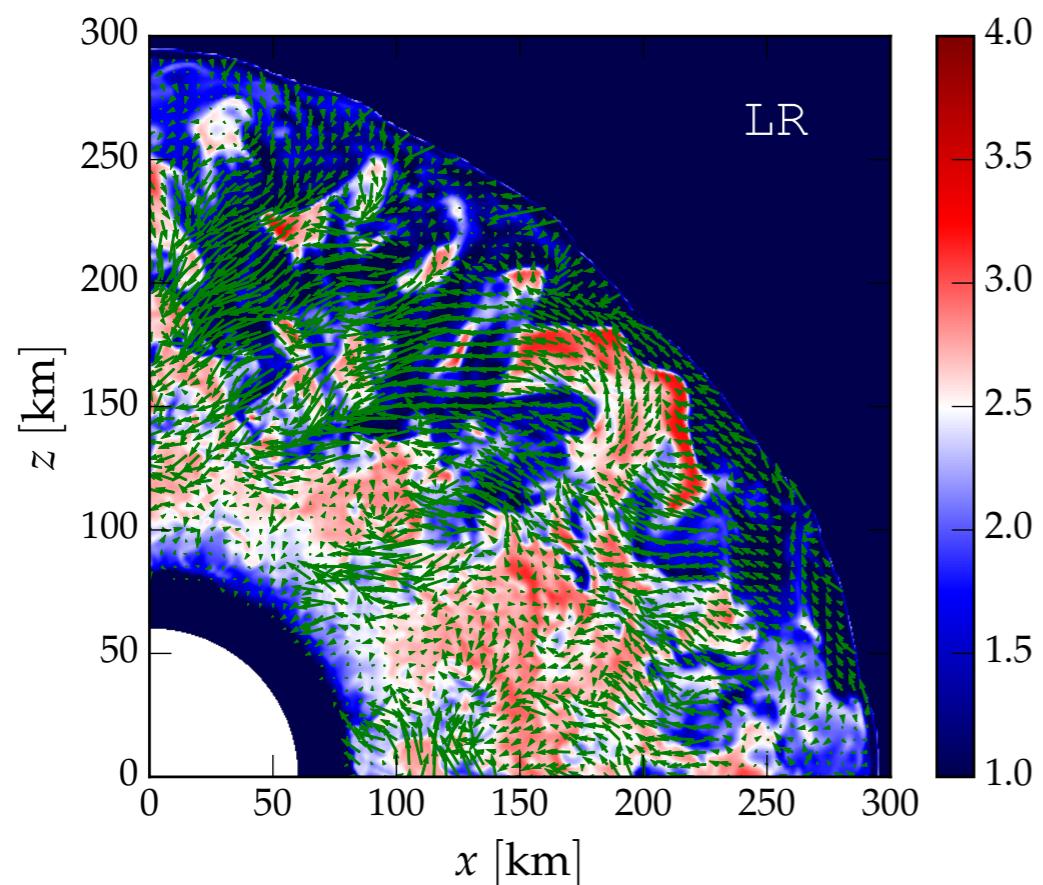
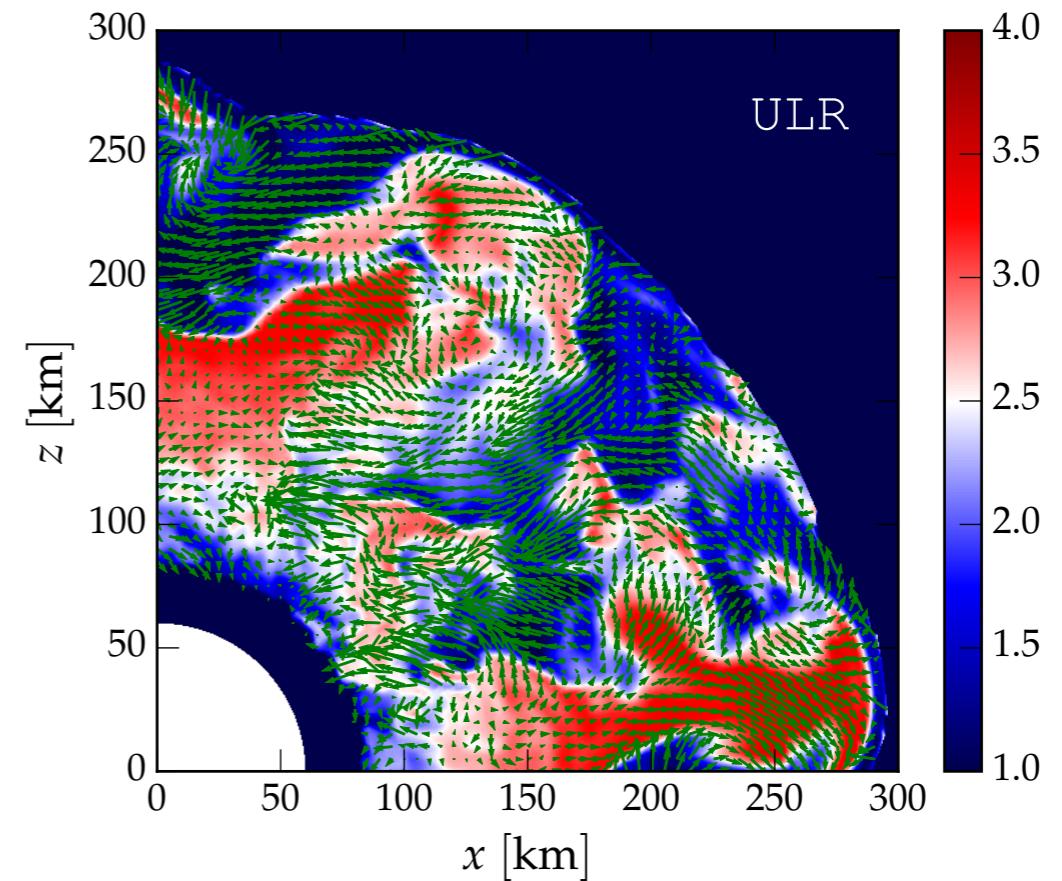
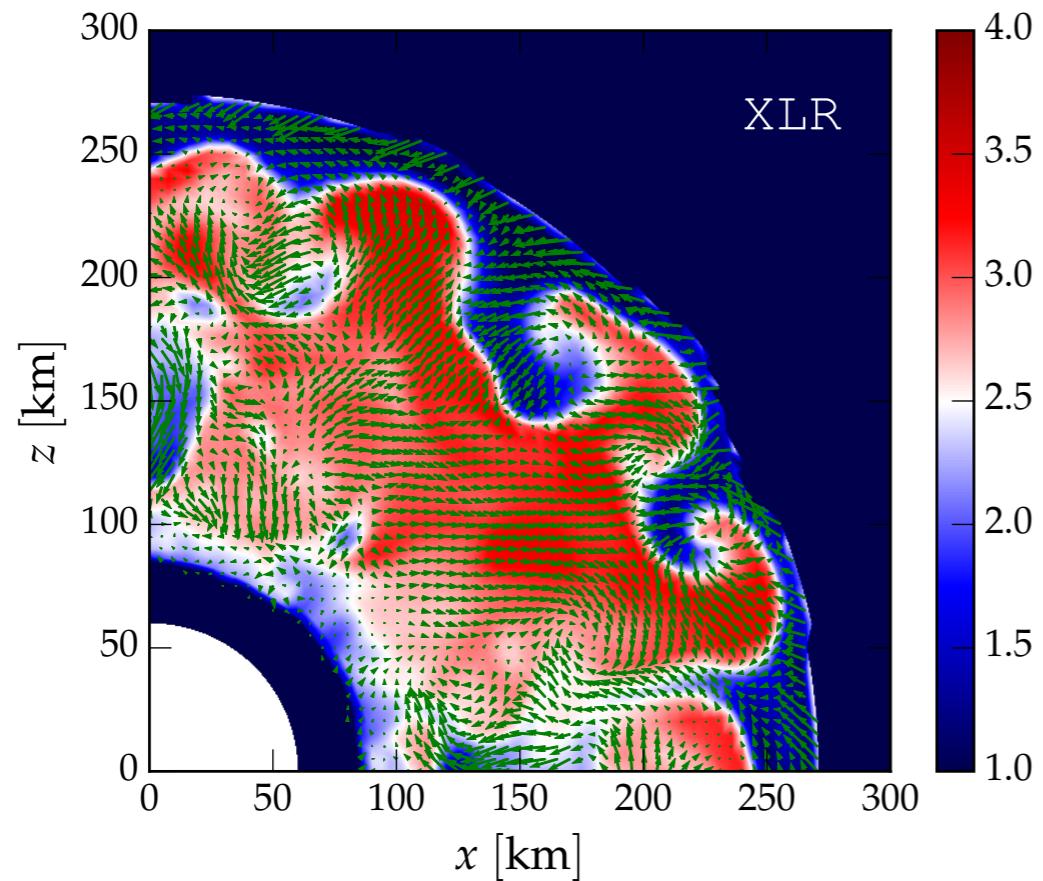


# Turbulent Energy Spectrum

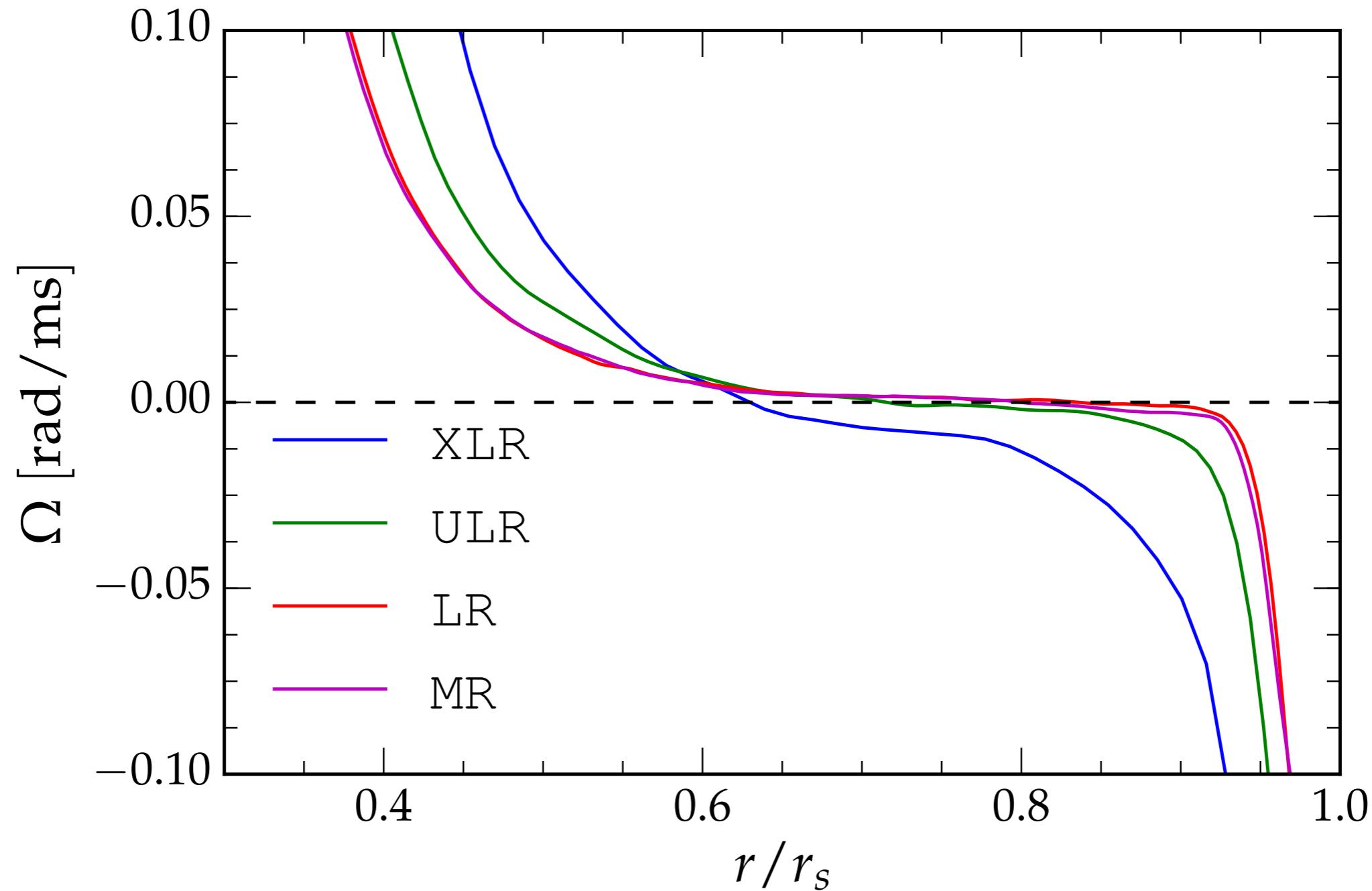


# Semi-Global Convection Study

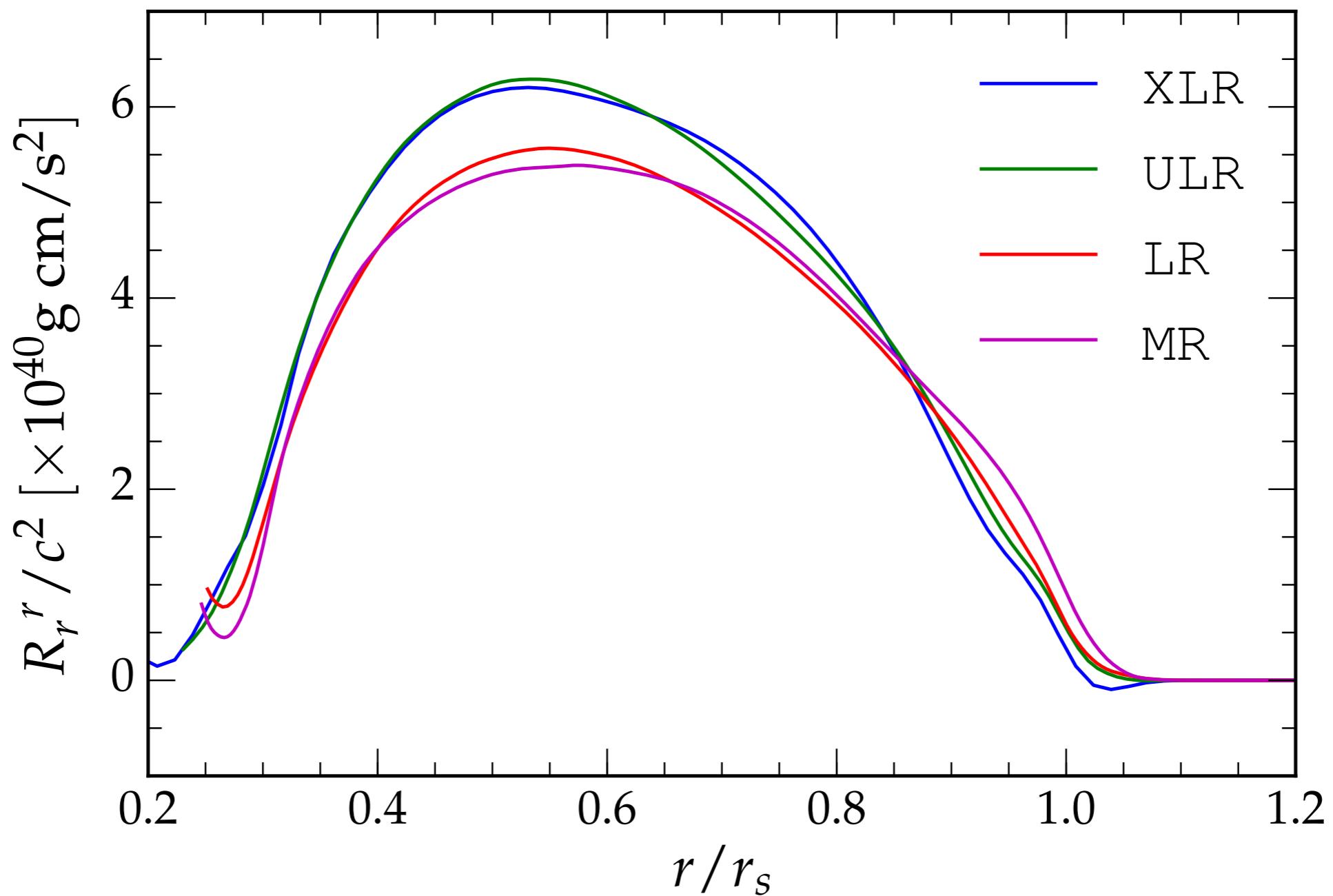




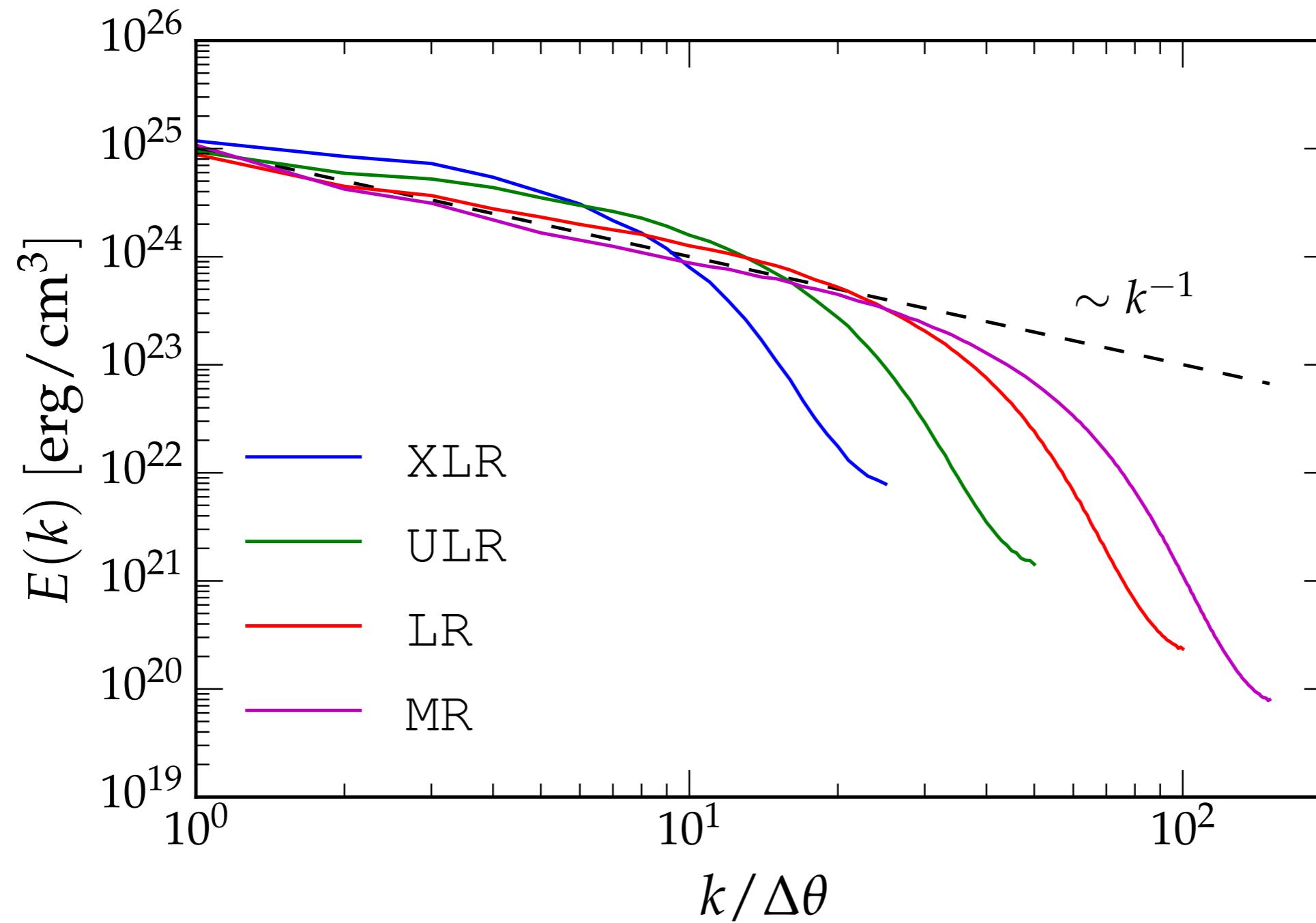
# Convective Instability



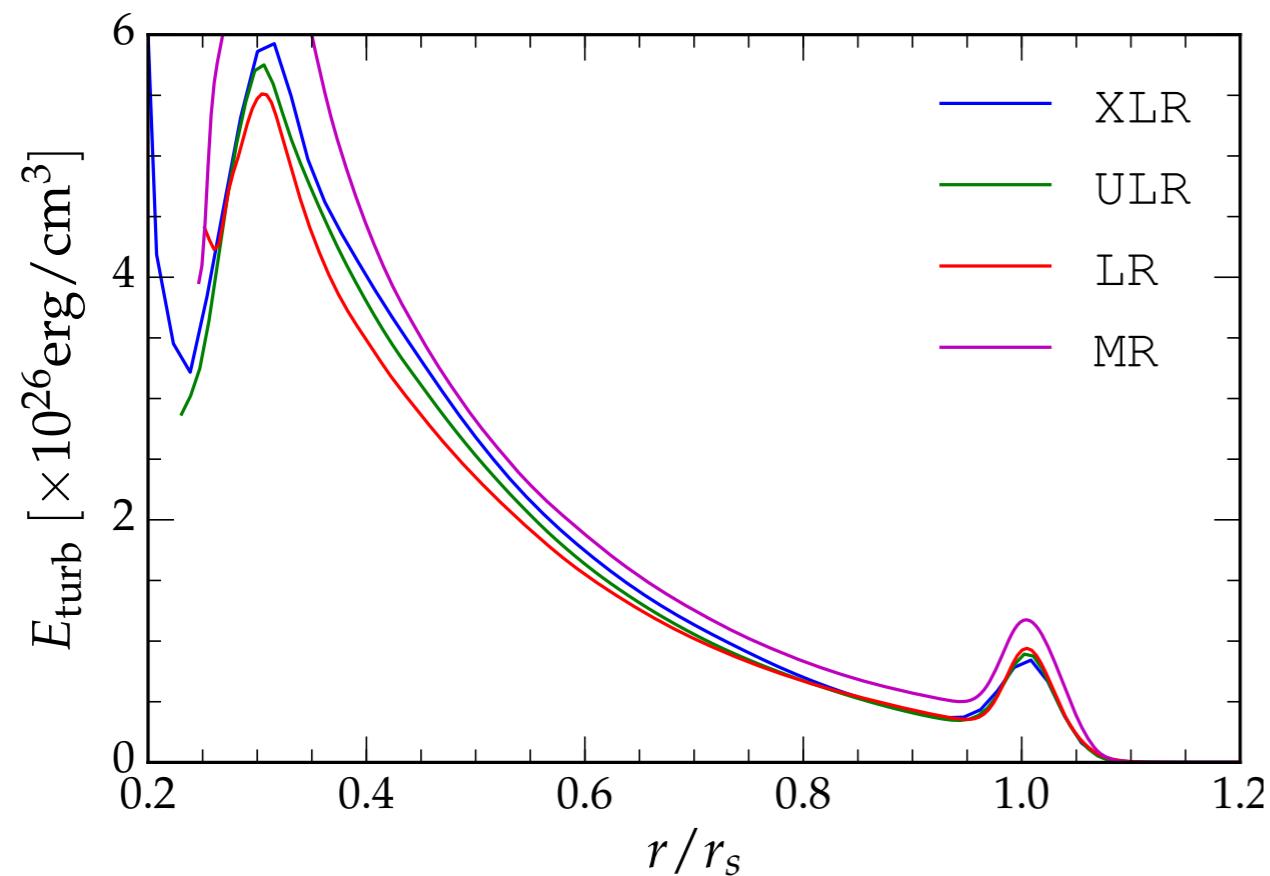
# Radial Reynolds Stresses



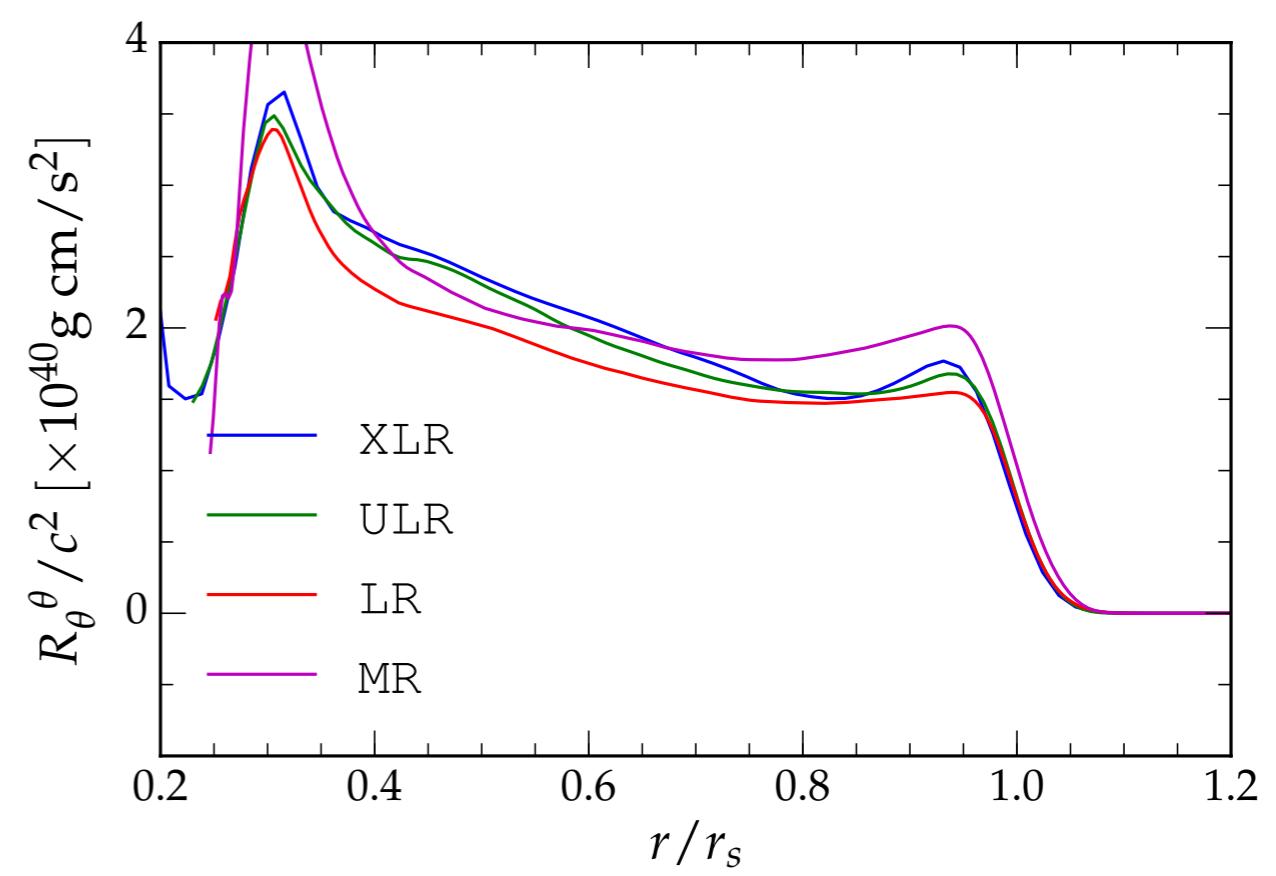
# Not Quite There Yet



# A New Ingredient: Intermittency I

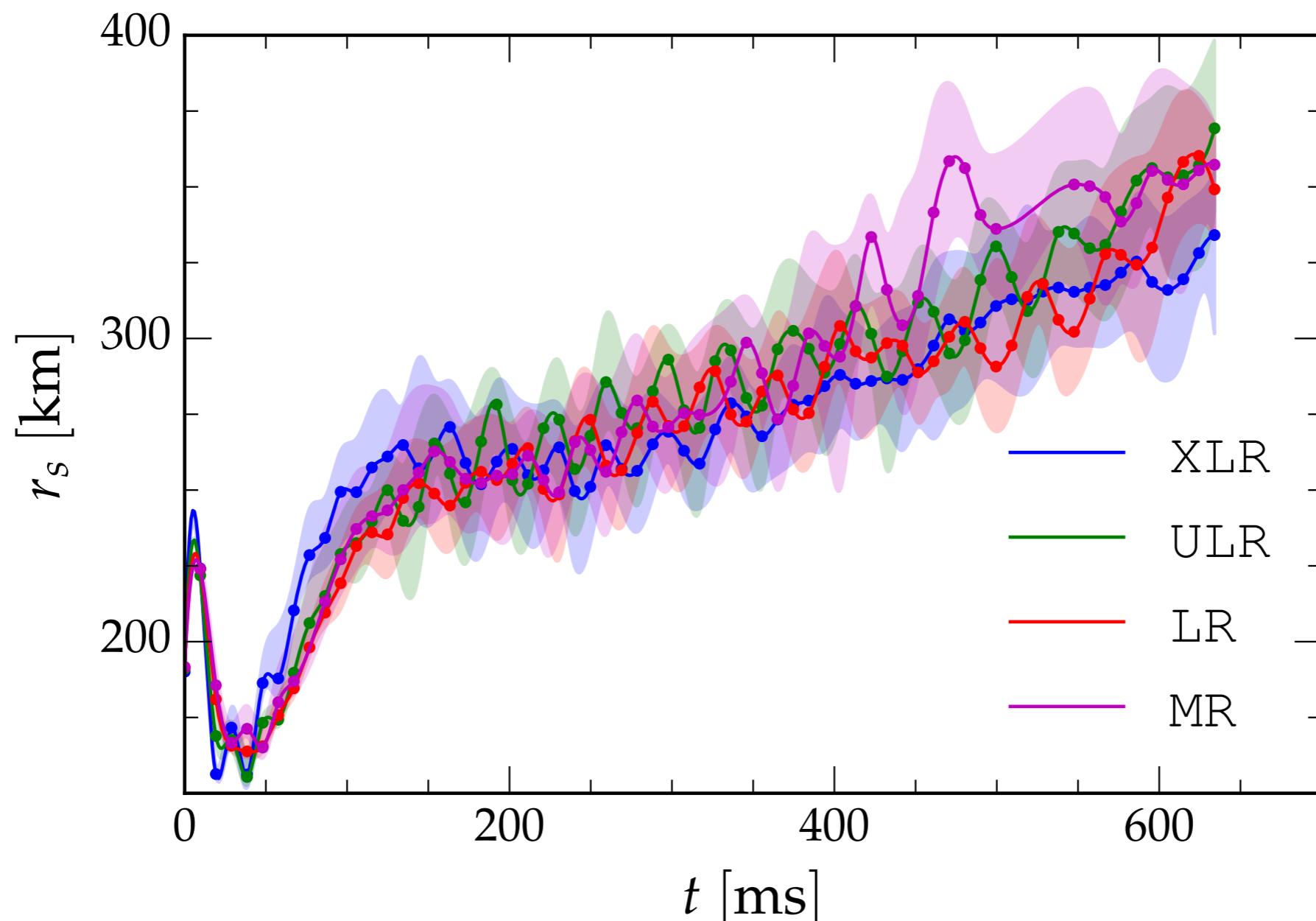


Turbulent energy density



Tangential Reynolds stress

# A New Ingredient: Intermittency II



Shock radius evolution

# Contents

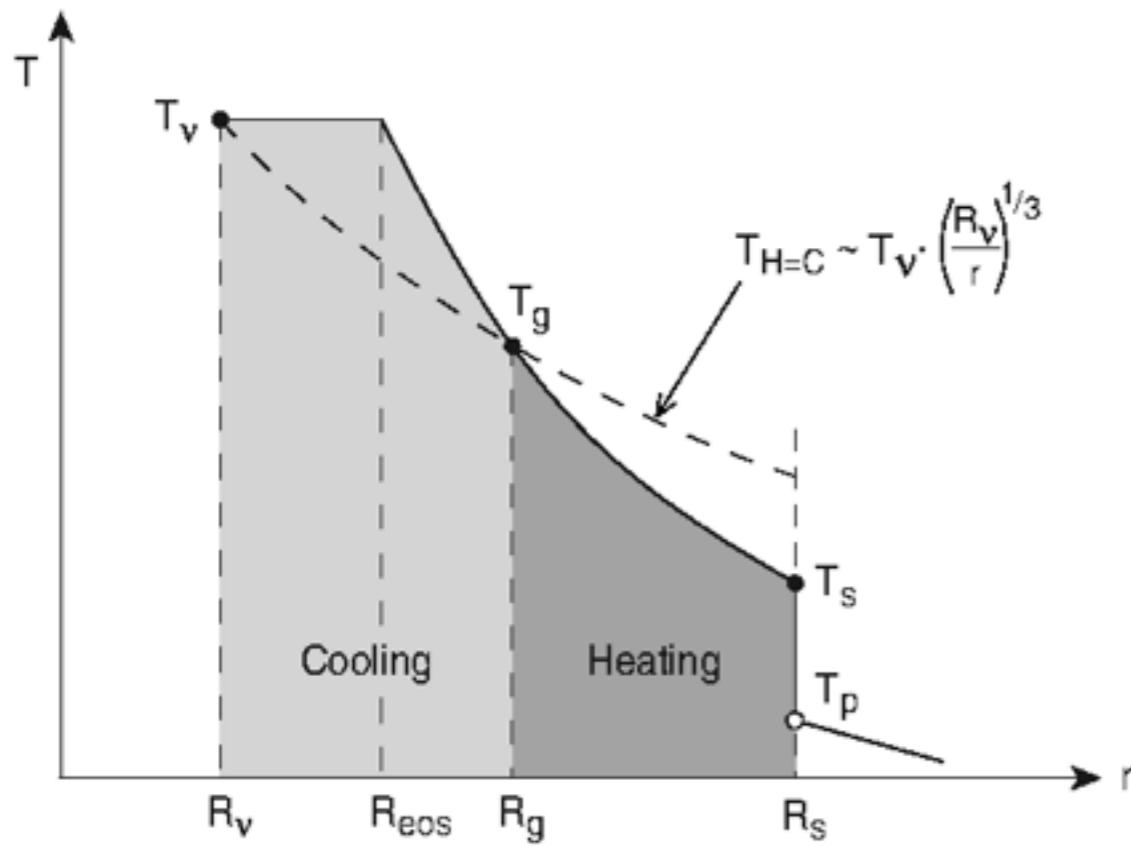
- 1.Turbulence in core-collapse supernovae
- 2.Numerical simulations
- 3.Conclusions**

# Conclusions

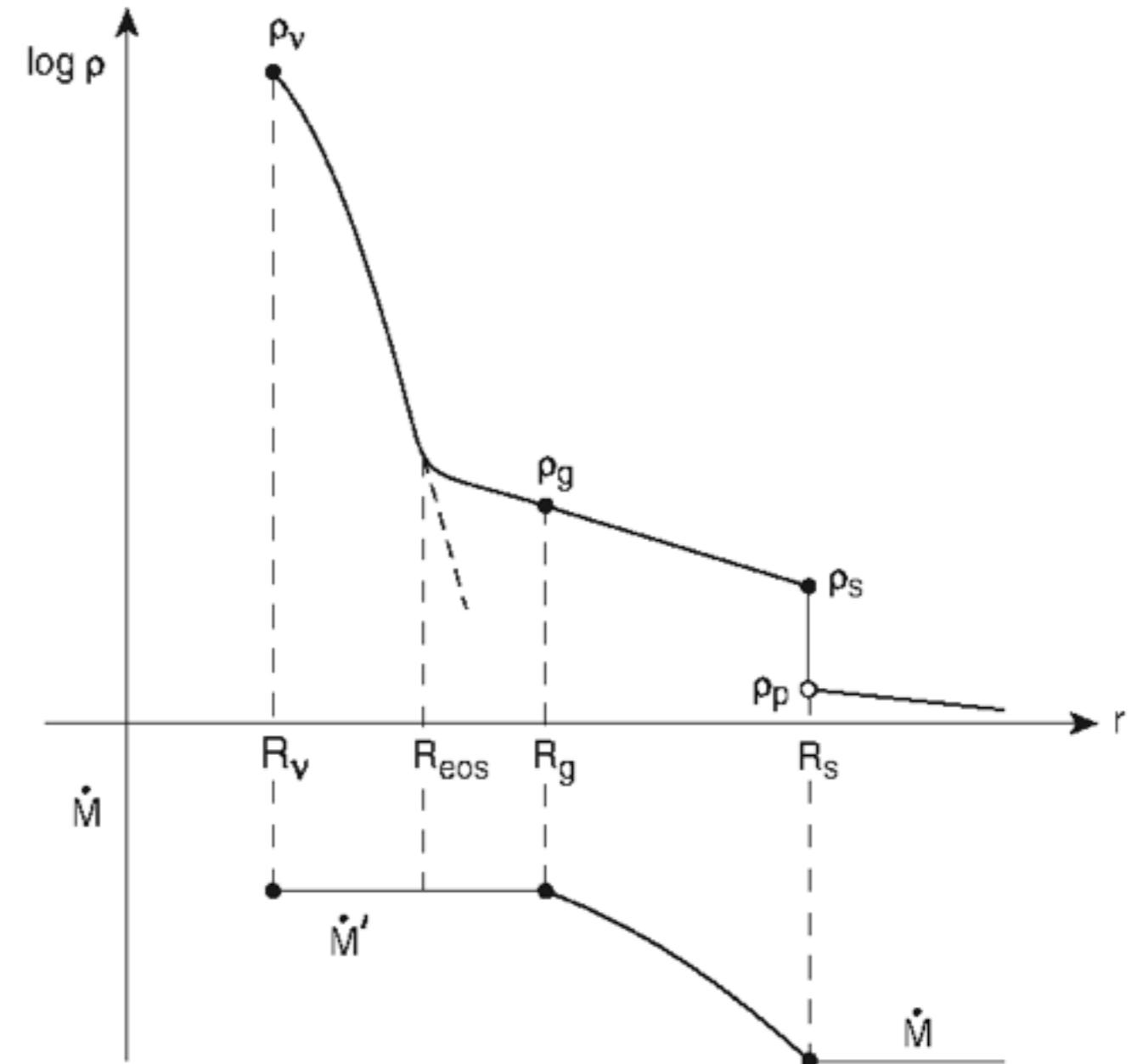
- Turbulence: crucial role for supernova explosions
- Local simulations: very high resolution is needed
- Idealized global simulations: rich dynamics of turbulent convection



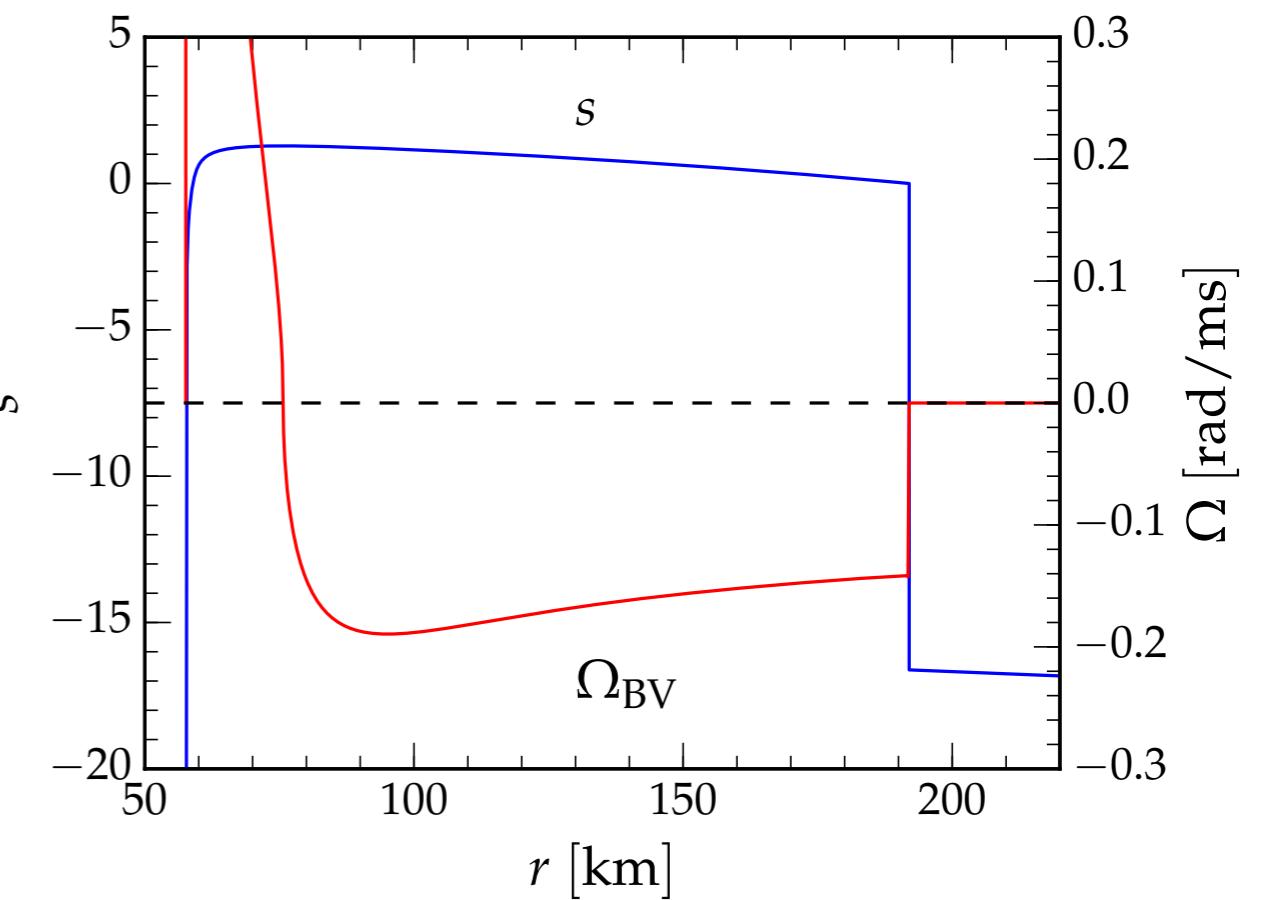
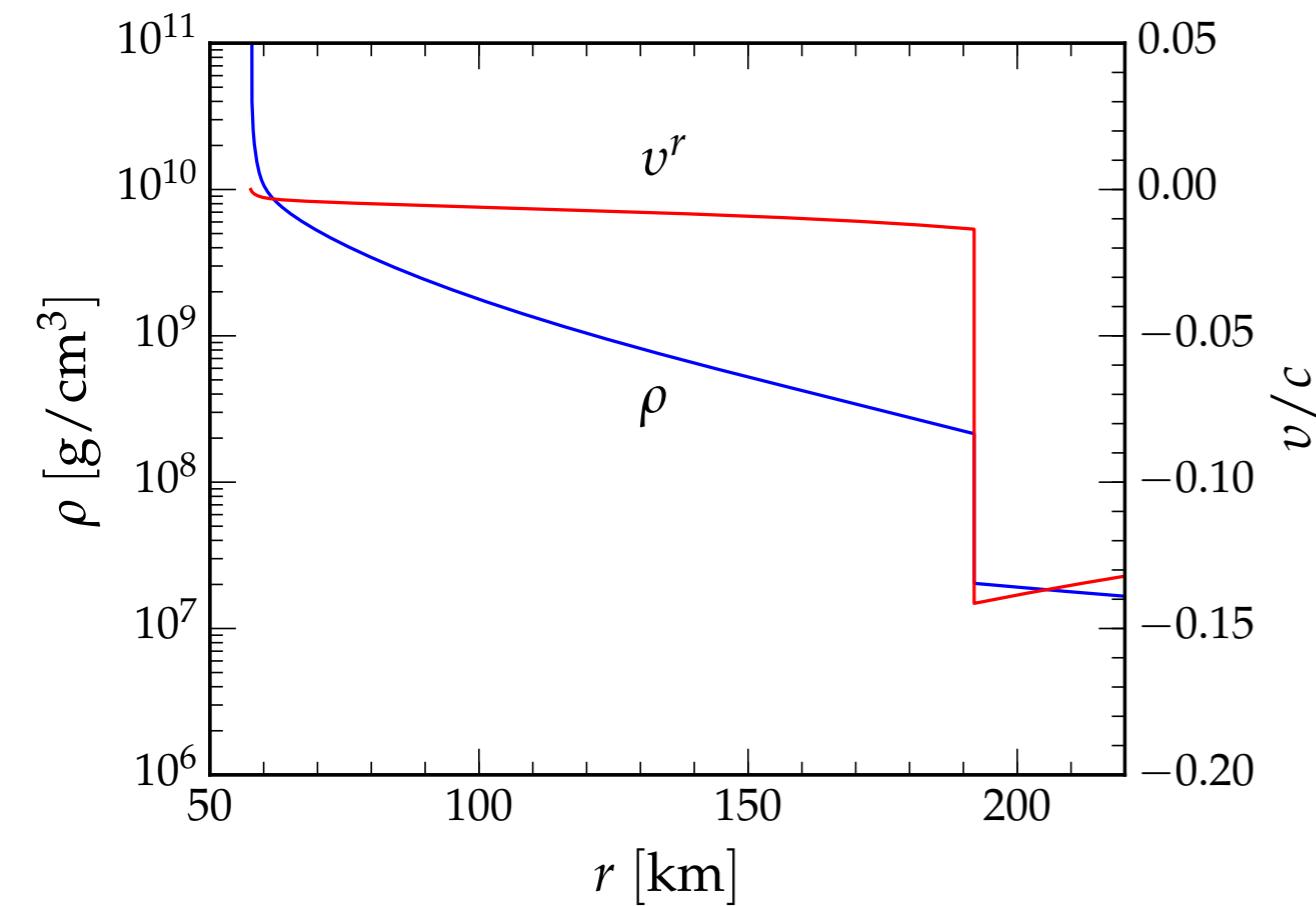
# The Standing Shock Flow



From Janka 2001



# Initial Data



Stationary initial data