### Entanglement in the sky

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Entanglement in the sky

Plan:

\* The vacuum

- entanglement, squeezing and the CMB
- origin of primordial entanglement
- \* Inflation and pre-inflationary initial conditions
  - PLANCK constraints on models
  - gravity-driven inflation, initial conditions and entanglement

# The Vacuum



No particles  $a(\vec{k}) |0\rangle = 0$ 

Vanishing expectation value

$$\langle 0 | \varphi(\vec{x}) | 0 \rangle = 0$$

Non-vanishing correlations

$$\langle 0 | \varphi(\vec{x}) \varphi(\vec{y}) | 0 \rangle = \int \frac{d^3 \vec{k}}{(2\pi)^3} \frac{1}{|\vec{k}|} \cos\left(\vec{k} \cdot (\vec{x} - \vec{y})\right) = \frac{1}{2\pi^2} \frac{1}{|\vec{x} - \vec{y}|^2}$$

Fluctuations of the field averaged over a region of size  $\ R$ 

$$(\Delta \varphi_R)^2 \equiv \langle 0 | \varphi_R \varphi_R | 0 \rangle - (\langle 0 | \varphi_R | 0 \rangle)^2 \sim \frac{1}{R^2}$$







## EPR pair: $\pi^0 \rightarrow e^- + e^+$

### Entanglement Einstein-Podolsky-Rosen (1935)





0



### Ground state

|0
angle|0
angle

#### Entangled oscillators



#### Ground state

without the spring |0
angle|0
angle

with the spring



#### Entangled oscillators



without the spring |0
angle|0
angle

with the spring





#### Entangled oscillators



without the spring |0
angle|0
angle

with the spring

$$|\Omega\rangle = \frac{1}{\cosh r} \sum_{n=0}^{\infty} (\tanh r)^n |n\rangle |n\rangle$$



# The vacuum state of a quantum field is highly entangled





### Squeezing the vacuum

#### examples:

- Schwinger effect
- Hawking radiation
- Primordial density fluctuations



Squeezing the vacuum and entanglement entropy growth Two new results:

(I) Entanglement growth at instabilities

$$S_A(t) \equiv -\text{Tr}_A(\rho(t)\log\rho(t)) \sim \left(\sum_i \lambda_i\right) t$$

<u>linear growth: rate given by the sum of the positive Lyapunov exponents</u> Bianchi-Hackl-Yokomizo Phys.Rev D (2015)



Squeezing the vacuum and entanglement entropy growth Two new results:

(2) Entanglement growth at parametric resonances

$$S_A(t) \equiv -\text{Tr}_A(\rho(t)\log\rho(t)) \sim \left(\sum_i \mu_i\right) t$$

### linear growth: rate given by the sum of the positive Floquet exponents

Bianchi-Hackl-Yokomizo to appear (2016)



### Squeezing the vacuum

#### examples:

- Schwinger effect
- Hawking radiation
- Primordial density fluctuations

### Schwinger effect (1951)

Vacuum fluctuations are unstable in the presence of a strong electric field



Extraction of entangled pairs of electrons and positrons from vacuum fluctuations



EB, Hackl and Yokomizo (2016)

## Hawking radiation (1974)

Vacuum fluctuations are unstable in the presence of a strong gravitational field





## The vacuum of a quantum field before inflation



$$P(k) = \frac{1}{k}$$

## The vacuum of a quantum field after inflation

Mukhanov-Chibisov (1981)

$$P(k) = \frac{(aH)^2}{k^{3n_s}}$$
$$n_s \lesssim 1$$



### Mechanism: amplification of vacuum fluctuations by instabilities



Harmonic oscillator with time-dependent frequency

$$H(t) = \frac{1}{2}p^2 + \frac{1}{2}(k^2 - f(t))q^2$$

## The vacuum of a quantum field before inflation



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The anisotropies of the Cosmic Microwave Background as observed by Planck



Planck Collaboration, arxiv.org/abs/1502.02114 ``Planck 2015 results. XX. Constraints on inflation"



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### Origin of primordial entanglement

BKL conjecture (Belinsky-Khalatnikov-Lifshitz 1970)

In classical General Relativity, the spatial coupling of degrees of freedom is suppressed in the approach to a space-like singularity

Quantum BKL conjecture (Bianchi-Hackl-Yokomizo 1512.08959)

In quantum gravity, correlations between spatially separated degrees of freedom are suppressed in the approach to a Planck curvature phase

$$\begin{cases} \hat{H}\Psi[g_{ij}(x),\varphi(x)] = 0\\ \lim_{K \to 1/\ell_P^4} \Psi[a,\phi,\delta g_{ij}(x),\delta\varphi(x)] = \prod_{\vec{x}} \psi(\phi,\delta g_{ij}(x),\delta\varphi(x)) \end{cases}$$

Scenario: correlations present at the beginning of inflation are produced during a pre-inflationary BKL phase via vacuum squeezing

Motivated by the "architecture conjecture" (Bianchi-Myers, CQG 2013)

### Origin of primordial entanglement

Pre-inflationary scenario:

BKL phase with entanglement entropy growing from a zero law to an area law



EB, Hackl and Yokomizo 1512.08959

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PLANCK 2013 constraints on inflationary models

[1303.5082] 0.4 *Planck*+WP+BAO:  $\Lambda$ CDM + r quadratic potential *Planck*+WP+BAO:  $\Lambda$ CDM + r + (d $n_s$ /d ln k) Tensor-to-scalar ratio (r)0.1 0.2 0.3 plateau potential 0.0 0.94 0.96 0.98 1.00[Starobinsky 1979] Primordial tilt  $(n_s)$ 





Slow-roll initial conditions in  $R + \alpha R^2$  gravity Action  $S = \frac{1}{16\pi G} \int (R + \alpha R^2) \sqrt{-g} d^4 x$ 

Classical FLRW background + quantum perturbation

$$g_{\mu\nu}(\vec{x},t) = \bar{g}_{\mu\nu}(t) + \delta g_{\mu\nu}(\vec{x},t)$$

Einstein equations

$$G_{\mu\nu} + \alpha \,\mathcal{H}_{\mu\nu} = 0$$

Gravity-driven inflation

$$\begin{cases} A_s = \frac{1}{18\pi^2} \frac{\ell_P^2}{\alpha} N^2 \\ n_s = 1 - \frac{2}{N} \\ r = \frac{12}{N^2} \end{cases}$$

Compatible with data with  $N \simeq 60$  and  $\alpha \simeq (10^5 \, \ell_P)^2$ 

#### Planck 2015

$$\begin{cases} A_s = (2.196 \pm 0.05) \times 10^{-9} \\ n_s = 0.9603 \pm 0.0073 \\ r < 0.11 \quad \text{at } k_* = (20 \,\text{Mpc})^{-1} \end{cases}$$

<u>Pre-inflationary initial conditions</u> in  $R + \alpha R^2$  gravity EB-Fernandez-Satz, *in progress* 

Initial condition in post-Planckian phase

 $R_{\mu\nu\rho\sigma}R^{\mu\nu\rho\sigma} = 1/\ell_0^4 \qquad \text{with} \qquad \ell_P < \ell_0 \ll 10^5 \,\ell_P$ 

Background: select initial conditions s.t.  $N \ge 60$  e-folds of slow-roll

Quantum perturbation: scalar and tensor modes of the geometry local vacuum has low entanglement entropy at the curvature scale

Pre-slowroll phase produces Minkowskian correlations + squeezing

Expected enhancement of power in tensor modes at low  $\ell$ 



EB-Fernandez (2016)



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