# Extragalactic sources and ultra-high energy cosmic rays

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# Outline

- Cosmic-ray spectrum characteristics
- Sources (non- & relativistic)
- Shocks and jets
  - Properties
  - Particle acceleration mechanism
- Shock acceleration simulation studies overview
  - Numerical method
  - Individual and multiple relativistic shocks in AGN
  - Propagation and radiation
- Conclusion



# **Cosmic-rays**



- Cosmic-rays are subatomic particles & radiation of extra-terrestrial origin.
- First discovered in 1912 by Victor Hess, measuring radiation levels aboard a balloon up to 5300m
- Hess found increased radiation levels at higher altitudes: named it Cosmic Radiation



# **Cosmic-ray spectrum**

Energies and rates of the cosmic-ray particles





# The high energy regime - 'knee(s)' & 'ankle'





# The ultra-high energy regime – the 'toe'



	TA	Auger
$\gamma_1$	$3.33\pm0.04$	$3.27\pm0.02$
$\gamma_2$	$2.68\pm0.04$	$2.68\pm0.01$
$\gamma_3$	$4.2\pm0.7$	$4.2 \pm 0.1$
$lg(E_1/eV)$	$18.69\pm0.03$	$18.61\pm0.01$
$lg(E_2/eV)$	$19.68\pm0.09$	$19.41\pm0.02$



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Sources: Non-relativistic Relativistic



Sources: Non-relativistic Relativistic



## Supernovae







SN 1987A



Sources: Non-relativistic Relativistic



![](_page_13_Figure_0.jpeg)

![](_page_14_Picture_0.jpeg)

# A 'hidden force' in extragalactic jets: Shocks

![](_page_15_Picture_0.jpeg)

# Individual or multiple shocks

![](_page_15_Figure_2.jpeg)

#### Supersonic/superalfvenic strong compression waves → change gas/plasma's v, d, p, T - Collisional shocks (ordinary fluid) - Colissionless astrophysical shocks: In diffuse regions, low densities, large bulk speeds

![](_page_15_Figure_4.jpeg)

![](_page_15_Picture_5.jpeg)

![](_page_15_Picture_6.jpeg)

![](_page_15_Picture_7.jpeg)

PKS 1510-089

![](_page_15_Figure_9.jpeg)

![](_page_15_Figure_10.jpeg)

![](_page_15_Picture_11.jpeg)

![](_page_16_Picture_0.jpeg)

# **Shock classification - magnetic field orientation**

![](_page_16_Figure_2.jpeg)

![](_page_17_Picture_0.jpeg)

# Shock jump-conditions (Rankine-Hugoniot relations)

 $\rightarrow$  HD jump conditions (planar shocks,  $[X]_1^2 \equiv X_2 - X_1$ )

![](_page_17_Figure_3.jpeg)

$$\begin{split} \left[\rho_m V_z\right]_1^2 &= 0 \qquad \text{mass conservation} \\ \left[\rho_m \vec{V} V_z + p \vec{e}_z\right]_1^2 &= 0 \qquad \text{momentum cons.} \\ \left[\frac{1}{2}\rho_m V^2 V_z + \gamma p V_z / (\gamma - 1)\right]_1^2 &= 0 \qquad \text{energy cons.} \end{split}$$

→ MHD 
$$[\rho_m V_z]_1^2 = 0 \left[ \rho_m \vec{V} V_z + \left( p + \frac{B^2}{2\mu_0} \right) \vec{e}_z - \frac{B_z \vec{B}}{\mu_0} \right]_1^2 = 0 \left[ \left( \frac{1}{2} \rho_m V^2 + \frac{\gamma p}{\gamma - 1} + \frac{B^2}{\mu_0} \right) V_z - \frac{B_z \vec{B} \cdot \vec{V}}{\mu_0} \right]_1^2 = 0 [B_z]_1^2 = 0 [V_z \vec{B}_t - B_z \vec{V}_t]_1^2 = 0$$

Rankine (1870), Hugoniot (1887) Parker (1965), Hudson (1965), Parks (1984)

![](_page_18_Picture_0.jpeg)

# Particle acceleration mechanism at shocks

No doubt collisionless *astrophysical* shocks accelerate particles

Convincing evidence (early 80s) for efficient acceleration in heliospheric shocks and in SNRs

![](_page_19_Picture_0.jpeg)

# The Fermi mechanism

Transfer of the macroscopic kinetic energy of moving magnetized plasma to individual charged particles  $\rightarrow$  non-thermal distribution

2nd order Fermi acceleration (Fermi '49,'54)

@magnetic plasma clouds

#### 1st order Fermi acceleration - diffusive acceleration

(Krymskii '77, Bell '78, Blandford & Ostriker '78, Axford et al. '78)

@plasma shocks

![](_page_20_Picture_0.jpeg)

## 1st order Fermi acceleration – diffusive acceleration of CRs

Test particle - diffusion - n acceleration shock cycles

 $E_n = (x+1)^n \cdot E_0$ 

Energy gain: fraction of initial energy

$$\Delta E = E - E_0 = x \cdot E_0$$

Average energy gain per collision:

 $<\Delta E/E>\cong (2V/c)$ 

Leading to a power-law energy behaviour

$$N(>E) = \sum_{i=n}^{\infty} (1 - P_{esc})^{n(E)} = ... \propto E^{-\sigma}$$

$$\sigma = (r+2)/(r-1), r = V_1/V_2 = (\gamma+1)/(\gamma-1)$$

for mono-atomic gas:  $\gamma = 5/3 \Rightarrow r = 4 \Rightarrow E^{-2}$ 

Important: <u>Non-relativistic shocks</u>: σ is *constant* (~ 2.2) independent of shock-B inclination (Drury, '83) Relativistic shocks: Different story...

![](_page_20_Figure_13.jpeg)

(e.g. Krymskii '77, Bell '78, Drury '83)

![](_page_21_Picture_0.jpeg)

#### Note: Facts for non-relativistic shock acceleration

- Particles are everywhere in <u>isotropy</u> and the diffusive approximation for solution of the transport equation can apply
- Spectral index (σ) <u>independent</u> of: scattering nature (κ), inclination (ψ) and strength of magnetic field (B)

Concepts are well understood and well studied - they work well as a comparison basis for *relativistic* studies

![](_page_22_Picture_0.jpeg)

### Acceleration time scale & diffusion

The acceleration rate wins in competition with the time scale of the energy *losses* and the *escape rate*, defining the limit for the possible highest energies to be achieved.

#### **Acceleration rate:**

 $\tau(E) = (E \cdot \tau_{cycle}) / \Delta E = [3/(V_1 - V_2)] (\kappa_1 / V_1 + \kappa_2 / V_2) \text{ (Drury '83)}$ 

Confinement distance

**One cycle:** 

$$\tau_{cycle}$$
 (E)= (4/c )( $\kappa_1/V_1 + \kappa_2/V_2$ )

**Diffusion coefficient:** 

 $\kappa = \kappa_{\parallel} \cos^2 \psi$   $\kappa_{\parallel} = (1/3)\lambda \upsilon$   $\lambda = 10r_l$  (Quenby & Meli '05)

*i.e.* Proton 10GeV:  $\kappa$  about 10<sup>22</sup> cm<sup>2</sup>/s  $\rightarrow \tau_{cvcle}$  about 10<sup>4</sup> sec

![](_page_23_Picture_0.jpeg)

# Simulations of relativistic shock acceleration

![](_page_24_Picture_0.jpeg)

# **Relativistic shock acceleration: Questions**

- Is spectral index (σ) <u>universal</u>? Flat or steep?
- o depends on: gamma shock speed, inclination and scattering modes (turbulence of the media) ?
- Efficient acceleration → UHECRs ?
  - see: Ellison et al. (1995), Meli & Quenby (2003a,b, 2005), Niemec & Ostrowski (2004), Ellison & Double (2004), Stecker et al. (2007), Meli et al. (2008)

![](_page_24_Figure_6.jpeg)

![](_page_25_Picture_0.jpeg)

# **Numerical approaches**

- Semi-analytic solutions to diffusion-convection equation
  (e.g. Eichler '84, Berezhko & Ellison '99, Blasi & Gabici '02-'05)
- Numerical solutions to diffusion-convection equation with flow hydrodynamics & momentum dependent diffusion (e.g. Berezhko, Voelk et al. '96, Kang & Jones '91-'05, Malkov '97-'01)
- Monte Carlo simulations ('test-particle' approach)
  (e.g. Ellison et al. '02-'12, Baring '03-'13, Meli et al. '03-'14)

#### Particle-in-cell (PIC) simulations

(e.g. Dieckmann, Meli, et al. '08-'10, Nishikawa et al. (Meli), '13,'14)

![](_page_26_Picture_0.jpeg)

# Monte Carlo 'test-particle' approach principles

- Notion of 'test-particles' very efficient & very fast in describing particle random walks - large number of particles
- Random number generation → simulation of the random nature of a physical process (Cashwell & Everett '59)
- Powerful tool  $\rightarrow$  large dynamic ranges in spatial and momentum scales
- Scattering can be treated via large angle and pitch angle diffusion approach (e.g. Kennel & Petscheck '66, Forman et al. '74, Jokipii '87, Quenby & Meli '05, Meli & Biermann '06)

$$\kappa = \kappa_{\parallel} \cos^2 \psi + \kappa_{\perp} \sin^2 \psi \qquad \kappa_{\perp} = \kappa_{\parallel} \cdot (1 + (\lambda/r_l)^2)^{-1} \qquad \kappa_{\parallel} \gg \kappa_{\perp}$$

- Fully relativistic Lorentzian transformations
- Pesc (probability of escape)

![](_page_27_Picture_0.jpeg)

Relativistic jets and UHECRs: Individual shocks Multiple shocks

![](_page_28_Picture_0.jpeg)

# Relativistic jets and UHECRs: Individual shocks Multiple shocks

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![](_page_29_Picture_0.jpeg)

# Sub-luminal (oblique) shocks - spectra

![](_page_29_Figure_2.jpeg)

Scattering :  $1/\Gamma < \theta < 10/\Gamma$ 

![](_page_30_Picture_0.jpeg)

# Super-luminal (perpendicular) shocks - spectra

![](_page_30_Figure_2.jpeg)

#### Superluminal shocks $\rightarrow$ Not efficient accelerators $\rightarrow$ Irregular spectra

![](_page_31_Picture_0.jpeg)

# **Application to extragalactic astronomy**

![](_page_32_Picture_0.jpeg)

# **Contribution to the diffuse UHE cosmic-ray signal ?**

![](_page_32_Figure_2.jpeg)

#### Fitting between: 3EeV and 30EeV

<u>Black line</u>: assumed contribution of UHECR from GRBs with flat spectra,  $\sigma$ =1.5.

<u>**Red line</u>**: half-half contribution with  $\sigma$ =1.5 and  $\sigma$ =2.1 respectively.</u>

**Blue line**: only UHECRs from AGN with  $\sigma$ =2.1.

After averaging various spectra, we assume that a diffused proton spectrum measured at Earth is given by:

$$\frac{\mathrm{d}N_{\mathrm{p}}}{\mathrm{d}E_{\mathrm{p}}} = A_{\mathrm{p}} \int_{z_{\mathrm{min}}}^{z_{\mathrm{max}}} \left( x \cdot \frac{\mathrm{d}\Phi_{2.1}}{\mathrm{d}E_{\mathrm{p}}} (E_{\mathrm{p}}(z)) + (1-x) \cdot \frac{\mathrm{d}\Phi_{1.5}}{\mathrm{d}E_{\mathrm{p}}} (E_{\mathrm{p}}(z)) \right)$$
$$\times (1+z)^{-1} \cdot \exp\left(-\frac{E_{\mathrm{p}}(z)}{E_{\mathrm{cut}}(z)}\right) \cdot g(z)\mathrm{d}z$$

<u>Condition 1</u>: UHECRs produced in subluminal relativistic shocks with spectra of mean  $\sigma = -2.1$ , contribute a fraction  $0 \le x \le 1$ , and UHECRs with flat spectra of  $\sigma = -1.5$  contribute 1-*x*. (0.001<*z*<7)

> <u>Condition 2</u>: We take into account particle propagation, adiabatic energy losses, source evolution g(z), absorption at the highest energies and normalized the flux using observations above the ankle.

#### Meli and Ciarcelluti (2014)

![](_page_33_Picture_0.jpeg)

Relativistic jets: Individual shocks Multiple shocks

![](_page_34_Picture_0.jpeg)

# Model: Multiple shock patterns and cosmic-ray acceleration in extragalactic jets with a single particle-injection (Meli and Biermann, 2013)

![](_page_34_Figure_2.jpeg)

![](_page_34_Picture_3.jpeg)

![](_page_34_Picture_4.jpeg)

Repeated multiple shocks with opening angles *a*, *b*, *c*, *d*, in an AGN jet, e.g., PKS 1510-086, CenA, M87, NGC6251, etc

![](_page_35_Picture_0.jpeg)

![](_page_35_Figure_1.jpeg)

![](_page_36_Picture_0.jpeg)

# **Proton spectra at the source (in a shock sequence)**

![](_page_36_Figure_2.jpeg)

Meli & Biermann (2013)

![](_page_37_Picture_0.jpeg)

# **Cosmic-rays** $\iff$ gamma-ray and neutrino astronomy

![](_page_38_Picture_0.jpeg)

# **Radiation by cosmic-rays**

![](_page_38_Figure_2.jpeg)

![](_page_39_Picture_0.jpeg)

# Leptonic continuum emission from AGN

![](_page_39_Figure_2.jpeg)

![](_page_40_Picture_0.jpeg)

### **Hadronic interactions**

![](_page_40_Figure_2.jpeg)

![](_page_41_Picture_0.jpeg)

# Simulations of extragalactic *propagation* of *hadronic* cosmic-rays

![](_page_41_Figure_2.jpeg)

by W.Wagner

![](_page_42_Picture_0.jpeg)

![](_page_42_Figure_1.jpeg)

![](_page_43_Picture_0.jpeg)

# **Questions answered:**

- UHECRs seem to originate from extragalactic sources such as AGN and GRB jets
- Relativistic shocks *individual or multiple* can accelerate UHECRs with a variety of spectral features
- Spectral index of the *primary* spectrum ( $\sigma$ ) is <u>not universal</u>: observations  $\rightarrow$  *gamma-rays, neutrinos*
- σ depends on: shock speed, inclination and scattering modes (turbulence of the media)
  - Faster shocks generate flatter distributions
  - Subluminal (quasi-parallel) shocks efficient accelerators  $\rightarrow \sim 10^{21} \, \text{eV}$  (!)
  - Superluminal (quasi-perpendicular) shocks not efficient → ~10<sup>15</sup> eV

![](_page_44_Picture_0.jpeg)

# Take home lesson (Monte Carlo CR studies)

Relativistic *individual or multiple shocks* in extragalactic jets are powerful engines, producing very high energy CR via the Fermi acceleration mechanism with *distinctive spectral features* and *consequent radiation* 

Immediate applications to extragalactic observational astronomy:

✓ Hadronic radiation models

Gamma-ray & neutrino astronomy

Multi-messenger approach

![](_page_45_Figure_0.jpeg)