New aspects of the QCD phase transition in proto-neutron stars and core-collapse supernovae

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Fonds national suisse Schweizerischer Nationalfonds Fondo nazionale svizzero Swiss National Science Foundation



Motivation: core-collapse supernovae

- how do massive stars explode?
- which progenitors end as black holes, which as neutron stars?
- what is their nucleosynthesis contribution, galactical chemical evolution?



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NASA/Chandra

Supernova explosion mechanism

- presently favored: neutrino-driven mechanism (Colgate & White 1966)
- explosion triggered by energy deposition of neutrinos in the infalling matter
- requires multi-dimensional fluid instabilities: convection, turbulence, "standing accretion shock instability", in general no explosions in 1D simulations
- •e.g. increases time of matter in gain region and thereby the neutrino heating



Example of a 3D supernova simulation

- simulation by Kuo-Chuan Pan, Basel (arXiv:1505.02513)
- 15 M_{sun} progenitor (Woosley et al. 2002)
- HS(DD2) EOS, relativistic-mean field, nucleons, nuclei, electrons (MH and Schaffner-Bielich 2010)
- hydrodynamics: FLASH (Lee 2003)
- neutrino transport: IDSA (Liebendörfer et al. 2009)





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Open questions in core-collapse supernova theory

- typically low explosion energies, dannot explain all supernovae
- differences between 2D and 31



Role of the EOS in core-collapse supernovae

- different nuclear interactions/hadronic EOSs have only a moderate impact (Janka 2012, MH et al. 2012, Suwa 2013, Kuo-Chuan Pan et al. 2015, ...)
- no 2D or 3D simulations with non-nucleonic degrees of freedom yet (!)

 \rightarrow what is the role of quark matter?

Core-collapse supernova explosions triggered by the QCD phase transition

Quark-hadron hybrid EOS for supernovae

- 2009/2011: Sagert, Pagliara, Schaffner-Bielich, MH
- hybrid EOSs available as tables for various temperatures and asymmetries, suitable for core-collapse supernova simulations
- hadronic phase: "STOS", Shen, Toki, Oyamatsu and Sumiyoshi 1998, 2011

 $-n,p,\alpha,A,e$

- non-linear relativistic mean-field interactions (TM1)
- Thomas-Fermi approximation for finite nuclei
- quark phase: bag model
 - -u,d,s (m_s=100 MeV)
 - first-order corrections for strong interactions, α_S (Farhi and Jaffe 1984)

$$p_i(m_i, T, \mu_i, \alpha_s) = p_i(m_i, T, \mu_i, 0) - \left[\frac{7}{60}T^4\pi^2\frac{50\alpha_s}{21\pi} + \frac{2\alpha_s}{\pi}\left(\frac{1}{2}T^2\mu_i^2 + \frac{\mu_i^4}{4\pi^2}\right)\right]$$

phase transition:

global charge neutrality (Gibbs PT/non-congruent PT)

CCSN explosions by the QCD phase transition



- phase transition induces collapse of the proto-neutron star
- once pure quark matter is reached, collapse halts
- formation of a second shock
- higher temperatures, increased neutrino heating \rightarrow positive velocities
- shock merges with standing accretion shock
- explosion

Neutrino signal



- colored lines with phase transition, black without
- second neutrino burst due to quark matter
- peak and height determine density and strength of the phase transition
- measurable with present day neutrino detectors [DasGupta et al. 2009]

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Mass-radius relation of hybrid EOS and SN explosions



explosions in spherical symmetry (T. Fischer et al. ApJS 2011)

- no explosions for sufficiently high maximum mass
- weak phase transition
- quark matter
 behaves similarly as
 hadronic matter
 "masquerade"
- cf.: Fischer,
 Blaschke, et al.
 2012: PNJL hybrid
 EOS

Densities reached in the supernova

Why does B145 not explode?



 \rightarrow critical density too high

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Neutrino signal for B139 hybrid EOS

[Fischer, et al. Acta Phys. Polon. Suppl. 7 (2014)]



- no second collapse, no explosion, no neutrino burst
- only slight reconfiguration of proto-neutron star
- moderate changes in neutrino signal
- no smoking gun for quark matter in supernovae

only few models tested, mechanism still possible for others?

Thermal properties of the hybrid EOS

QCD phase diagrams

- fundamental question: phase diagram of strongly interacting matter
- typically shown in T-μ, sometime also in T-ρ



Non-congruence of the nuclear liquid-gas and QCD phase [MH, V. Dexheimer, S. Schramm, I. Iosilevskiy, PRC 88 (2013)]

- main topic: differences of phase transitions with single or multiple conserved charges
- effects of isospin symmetry on phase diagrams
- "congruent" and "non-congruent" phase transitions, commonly known as "Maxwell" and "Gibbs"



Setup

- liquid-gas phase transition of nuclear matter: non-linear relativistic mean-field model FSUgold [Todd-Rutel and Piekarewicz, PRL (2005)]
- QCD phase transition: Chiral SU(3) model, includes quarks and hadrons as a chemical mixture of quasi-particle degrees of freedom [Dexheimer and Schramm, PRC81 (2010)]
- neglect of all Coulomb interactions, "Coulomb-less" approximation (cf. works by Gulminelli, Raduta, Typel, ...)
- solve for thermal, mechanical, and chemical equilibrium
- in the following: symmetric nuclear matter, zero strangeness locally, no leptons

Phase diagram of symmetric baryonic matter — $T-\mu_B$



- different scales
- similar shape, both phase transitions terminate in critical point

Is the QCD PT of liquid-gas type?



Phase diagram of symmetric baryonic matter — T-P



[losilevskiy, arXiv:1403.8053]

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see also:

[Satarov, Dmitriev, Mishustin, PAN72 (2009)] [Bombaci et al., PLB680 (2009)] [Steinheimer, Randrup, Koch, PRC89 (2014)]

The entropic QCD PT (dP/dT|_{PT}<0)

Clausius-Clapeyron equation

$$\left. \frac{dP}{dT} \right|_{\rm PT} = \frac{S^I - S^{II}}{1/n_B^I - 1/n_B^{II}}$$

• Steiner et al. PLB 468 (2000):

$$S = T\pi^2 \frac{\sum_{i} p_{F_i} \sqrt{p_{F_i}^2 + (m_i^*)^2}}{\sum_{i} p_{F_i}}$$

- more degrees of freedom (color, strangeness) in the quark phase, and more relativistic
- leads to high specific heat capacity and low temperatures
- \rightarrow QCD PT always entropic?
- what about color-superconducting phases? (cf. Rüster et al. PRD73 (2006))

General properties of *entropic* PTs (dP/dT|_{PT}<0)

- note: Clausius-Clapeyron equation only valid for a *congruent* (aka Maxwell)
 PT, i.e. where one has only one conserved charge
- then one also has:

$$\left. \frac{dP}{dT} \right|_{\rm PT} = \left. \frac{\partial P}{\partial T} \right|_{n_B}$$

• using general thermodynamic relations: unusual sign of 2nd cross derivatives, "abnormal thermodynamics", e.g.: [losilevskiy, arXiv:1403.8053]

$$\left. \frac{\partial P}{\partial T} \right|_{n_B} < 0 \Leftrightarrow \left. \frac{\partial T}{\partial n_B} \right|_S < 0$$

- dT/dn_B|_S<0 observed by many authors, also well-known in HIC
 - Steiner et al. PLB 2000
 - Nakazato et al. APJ 2010
 - Fischer et al. APJS 2011
 - Yudin et al. Astron. L 2013



Inverted convection in proto-hybrid stars

[A.V. Yudin, MH, D.K. Nadyozhiny, T.L. Razinkova, arXiv:1507.04598, accepted by MNRAS]

Matter and convection in proto-neutron stars

- entropy per baryon of 0-5
- trapped neutrinos in the early stage of the supernova, characterized by lepton fraction Y_L~0.4
- after one minute: neutrinofree, Y_v=0, beta-equilibrium
- negative entropy gradients leading to convection in the supernova and protoneutron star
- effect of convection: outward transport of hot matter, enhanced neutrino luminosities



[Roberts et al. PRL108 (2012)]

grayed regions: convectively unstable

Convection in proto-hybrid stars



 unusual thermal properties (abnormal thermodynamics) due to entropic PT: positive entropy gradients are convectively unstable (!)

[Yudin, MH, et al. arXiv:1507.04598]

Convection criteria

• relativistic Ledoux criterion for convection:

$$\frac{\partial \epsilon}{\partial S}\Big|_{P,Y_L} \frac{dS}{dr} + \frac{\partial \epsilon}{\partial Y_L}\Big|_{P,S} \frac{dY_L}{dr} > 0$$

- let's ignore composition changes, keep $Y_L=0.4=const. \rightarrow d\epsilon/dS$ determines convection
- dε/dS usually negative → negative entropy gradients are convectively unstable

$$\frac{\partial \epsilon}{\partial S}\Big|_{P,Y_L} = n_B T \left\{ 1 - \frac{\epsilon + P}{T \left. \frac{\partial P}{\partial T} \right|_{S,Y_L}} \right\}$$

• first term: small, relativistic correction • $\rightarrow dP/dT|_{s}>0 \Leftrightarrow d\epsilon/dS|_{P}<0$

Convection criteria for B165



 positive values: convectively unstable for positive entropy gradients

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Temperature for isentropes of B165



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Realizability and possible consequences

- realizability depends on
 - EOS model of hadronic and quark phase
 - description of the phase transition (surface tension)
 - structure of the proto-hybrid star
- possible consequences:
 - starting from quark core with low temperature
 - low entropy bubbles moving outwards, high entropy inwards
 - inward heat flux, increasing temperature of quark core, decreasing temperature of hadronic mantle, until negative entropy gradient is achieved
 - impact on supernova dynamics ???
 - imprint of quark matter on neutrino signal ???



CCSN explosions and the QCD PT

[MH, O. Heinimann A. Yudin, I. Iosilevskiy, M. Liebendörfer, F.-K. Thielemann, arXiv:1511.06551]

A third family of proto-compact stars



- third family feature ("twins") arises for high entropies
- result of the thermal properties of the EOS
- transition from second to third family releases gravitational energy of 10⁵⁰ to 10⁵³ erg

- explains the supernova explosions of Sagert and Fischer et al:
- proto-neutron star first on the second branch
- accretion until maximum reached
- collapse to third family, energy release, formation of 2nd shock, explosion

A third family of proto-compact stars — neutrino free



 for B139: third family arises only for very high entropies, much less pronounced

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A third family of proto-compact stars — trapped neutrinos



- neutrinos tend to suppress the third family feature
- less gravitational binding energy release, if at all

Matthias Hempel Frankfurt, 24.11.2015 Unusual thermal properties and stability of compact stars

- for convection: dε/dP|_S
- for stability: P(ε,S)
- to characterize thermal effects: $dP/dS|_{\epsilon}$
- dP/dS| $_{\epsilon}$ >0: stiffening, dP/dS| $_{\epsilon}$ <0: softening for increasing entropy
- using general thermodynamic relationships:

$$\left. \frac{\partial P}{\partial S} \right|_{\epsilon} = -T n_B \left(\frac{c_s}{c} \right)^2 + \frac{T}{C_V} \left. \frac{\partial P}{\partial T} \right|_{n_B}$$

- first term small, relativistic correction
- \rightarrow abnormal thermodynamics/entropic PT induces a softening of the EOS with increasing temperature/entropy (!)

Pressure-energy density relation



- hadronic and quark matter stiffens when it is heated
- in the phase coexistence region it softens (!)
- note: effect occurs only in parts of the phase coexistence (noncongruent PT)

→ the unusual thermal properties of the entropic PT are responsible for the supernova explosions





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Summary and conclusions

- phase diagram in P-T can provide interesting information
- is the QCD PT *entropic* (dP/dT|_{PT}<0)?
- entropic PTs lead to unusual thermal properties of the EOS, "abnormal thermodyamics"
- possible consequences in astrophysics:
 - inverted convection in proto-neutron stars
 - third family of proto-compact stars which exists only at finite entropy
 - core-collapse supernova explosions
- is it possible to achieve explosions by the QCD PT and have a maximum mass above 2 M_{sun}?
 - difficult to answer, requires new EOSs and new simulations
 - but: the maximum mass is determined at T=0, for the supernova the thermal properties are crucial

Comparing B139 and B165



- almost no temperature decrease for B139
- extremely extended phase coexistence
- "masquerade"

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Comparing B139 and B165



- interactions stiffen the quark phase
- softening with entropy very weak, if at all

Unusual thermal properties in other EOS



 relativistic mean-field EOS including hyperons and deltas





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NJL EOS

transition to