

# Viscosity in General Relativity

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Mathematically, the theory is very rich from both its analytic and geometric points of view. Over the past few decades, the subject of mathematical general relativity has matured into an active and exciting field of research among mathematicians.

# Important features of the Minkowski metric

In special relativity, fields live in **Minkowski space**, which is  $\mathbb{R}^4 = \mathbb{R} \times \mathbb{R}^3$  endowed with the **Minkowski metric**

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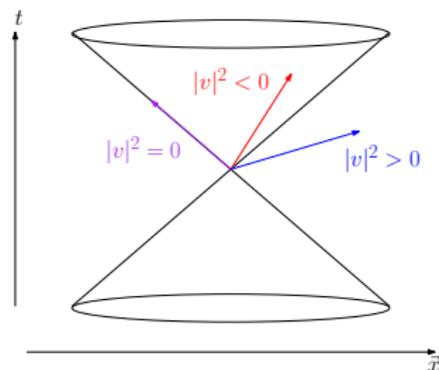
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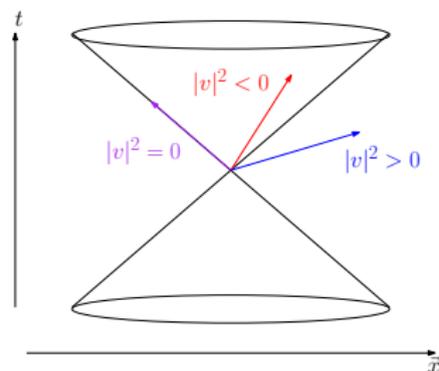
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**Causality:** The four-velocity of any physical entity satisfies  $|v|^2 = \eta_{\alpha\beta} v^\alpha v^\beta \leq 0$ . “Nothing propagates faster than the speed of light.”

# From special to general relativity

In general relativity the metric  $\eta$  is no longer fixed but changes due to the presence of matter/energy:  $\eta_{\alpha\beta} \rightarrow g_{\alpha\beta}(x)$ , where  $x = (x^0, x^1, x^2, x^3)$  are space-time coordinates.

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The values of  $g_{\alpha\beta}(x)$  depend on the matter/energy near  $x$ . Thus distances and lengths vary according to the distribution of matter and energy on space-time. This distribution, in turn, depends on the geometry of the space-time, i.e., it depends on  $g_{\alpha\beta}(x)$ .

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# The Ricci and scalar curvature

The Ricci curvature of  $g$  is

$$R_{\alpha\beta} = g^{\mu\nu} \left( \frac{\partial^2 g_{\alpha\beta}}{\partial x^\mu \partial x^\nu} + \frac{\partial^2 g_{\mu\nu}}{\partial x^\alpha \partial x^\beta} - \frac{\partial^2 g_{\alpha\nu}}{\partial x^\mu \partial x^\beta} - \frac{\partial^2 g_{\mu\beta}}{\partial x^\alpha \partial x^\nu} \right) + F_{\alpha\beta}(g, \partial g).$$

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$R_{\alpha\beta}$  and  $g_{\alpha\beta}$  are symmetric two-tensors ( $4 \times 4$  “matrix”), and  $R$  is a scalar.

Thus, Einstein’s equations are a system of second order partial differential equations for  $g_{\alpha\beta}$  (and whatever other fields come from  $T_{\alpha\beta}$ ).

# Coupling gravity and matter

Consider Einstein's equations

$$(\star) \begin{cases} R_{\alpha\beta} - \frac{1}{2}R g_{\alpha\beta} = T_{\alpha\beta}, \\ \nabla^\alpha T_{\alpha\beta} = 0, \end{cases}$$

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Matter fields = everything that is not gravity.

To couple Einstein's equations to any matter field, **all we need is  $T_{\alpha\beta}$** .

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Here,  $u$  is a (time-like) **unit** (i.e.,  $|u|^2 = g_{\alpha\beta}u^\alpha u^\beta = -1$ ) vector field representing the four-velocity of the fluid particles;  $p$  and  $\varrho$  are real valued functions describing the pressure and energy density of the fluid.

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The system is closed by an **equation of state**:  $p = p(\varrho)$ .

# Causality in general relativity

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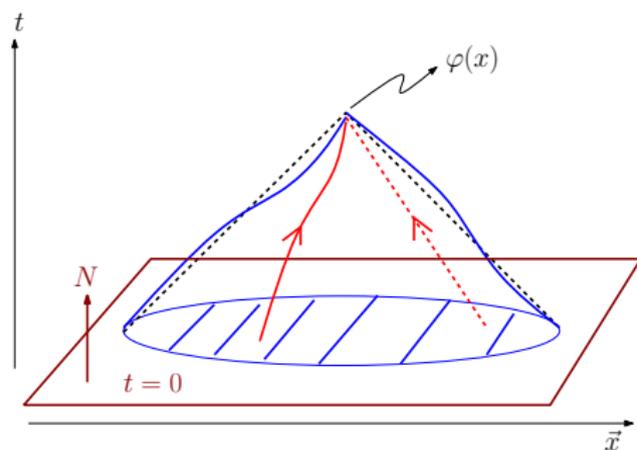
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Note that the causal structure is far more complicated than in Minkowski space since  $g_{\alpha\beta} = g_{\alpha\beta}(x)$ . One can better formulate causality in terms of the **domain of dependence** of solutions to Einstein's equations:



A theory is **causal** if for any field  $\varphi$  its value at  $x$  depends only on the “past domain of dependence of  $x$ .”

# What about fluids with viscosity?

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- ▶ Real fluids have viscosity.
- ▶ Cosmology. Perfect fluids exhibit no dissipation. Maartens ('95):  
“The conventional theory of the evolution of the universe includes a number of dissipative processes, as it must if the current large value of the entropy per baryon is to be accounted for. (...) important to develop a robust model of dissipative cosmological processes in general, so that one can analyze the overall dynamics of dissipation without getting lost in the details of particular complex processes.”

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- ▶ Astrophysics. Viscosity can have important effects on the stability of neutron stars (Duez et al., '04); source of anisotropies in highly dense objects (Herrera et al., '14).
- ▶ The treatment of viscous fluids in the context of special relativity is also of interest in heavy-ion collisions (Rezzolla and Zanotti, '13).

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All we need then is  $T_{\alpha\beta}^{NS}$  ( $T_{\alpha\beta}$  for Navier-Stokes).

$T_{\alpha\beta}$  is determined by the variational formulation/action principle of the matter fields.

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The stress-energy tensor is given by

$$T_{\alpha\beta} = \frac{1}{\sqrt{-\det(g)}} \frac{\delta \mathcal{L}}{\delta g^{\alpha\beta}}.$$

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Remark: stress-energy for the Navier-Stokes equations in non-relativistic physics is constructed "by hand."

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Eckart ('40) proposed the following stress-energy tensor for a relativistic viscous fluid

$$T_{\alpha\beta}^E = (\rho + p)u_\alpha u_\beta + p g_{\alpha\beta} - \left(\zeta - \frac{2}{3}\vartheta\right)\pi_{\alpha\beta} \nabla_\mu u^\mu - \vartheta \pi_\alpha^\mu \pi_\beta^\nu (\nabla_\mu u_\nu + \nabla_\nu u_\mu) - \kappa (q_\alpha u_\beta + q_\beta u_\alpha),$$

where  $\pi_{\alpha\beta} = g_{\alpha\beta} + u_\alpha u_\beta$ ,  $\zeta$  and  $\vartheta$  are the coefficients of bulk and shear viscosity, respectively,  $\kappa$  is the coefficient of heat conduction, and  $q_\alpha$  is the heat flux.

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$$T_{\alpha\beta}^E = (\rho + \varrho)u_\alpha u_\beta + \rho g_{\alpha\beta} - \left(\zeta - \frac{2}{3}\vartheta\right)\pi_{\alpha\beta}\nabla_\mu u^\mu \\ - \vartheta\pi_\alpha^\mu\pi_\beta^\nu(\nabla_\mu u_\nu + \nabla_\nu u_\mu) - \kappa(q_\alpha u_\beta + q_\beta u_\alpha),$$

where  $\pi_{\alpha\beta} = g_{\alpha\beta} + u_\alpha u_\beta$ ,  $\zeta$  and  $\vartheta$  are the coefficients of bulk and shear viscosity, respectively,  $\kappa$  is the coefficient of heat conduction, and  $q_\alpha$  is the heat flux.

$T_{\alpha\beta}^E$  reduces to the stress-energy tensor for a perfect fluid when  $\zeta = \vartheta = \kappa = 0$ , it is a covariant generalization of the non-relativistic stress-energy tensor for Navier-Stokes, and satisfies basic thermodynamic properties.

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Despite the results of Hiscock and Lindblom,  $T_{\alpha\beta}^E$  is still used in applications (particularly in cosmology) for the construction of phenomenological models.

# Entropy production

Define the **entropy current** as

$$S^\alpha = snu^\alpha + \kappa \frac{q^\alpha}{T},$$

where  $s$  is the specific entropy,  $n$  is the rest mass density, and  $T$  is the temperature.

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Equation (1) cannot be assumed. Rather, it has to be verified as a consequence of the equations of motion. This is one of the main constraints for the construction of relativistic theories of viscosity.

## Second order theories: the Mueller-Israel-Stewart theory

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In the MIS theory, the quantities  $\Pi$ ,  $\Pi_{\alpha\beta}$ , and  $Q_\alpha$  are treated as **new variables** on the same footing as  $\varrho$ ,  $u^\alpha$ , etc.

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This gives [equations for  \$\Pi\$ ,  \$\Pi\_{\alpha\beta}\$ , and  \$Q\_\alpha\$](#)  that are appended to Einstein's equations.

# Summary of results for second order theories

For the MIS and other second order theories:

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- ▶ Coupling to Einstein's equations? (Existence of solutions?)

## Back to first order theories

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Lichnerowicz ('55) proposed the following stress-energy tensor for a relativistic viscous fluid:

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Lichnerowicz's stress-energy tensor had been mostly ignored for many years, but recently it has been showed as potentially viable candidate for relativistic viscosity.

## Some results

Using Lichnerowicz's stress-energy tensor, it is possible to show (D-, '14; D- and Czubak '16; D-, Kephart, and Scherrer, '15):

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- ▶ Applications to cosmology lead to different models, in particular big-rip scenarios.
- ▶ None of these results consider all dissipative variables (e.g. shear viscosity but no bulk viscosity, etc).

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Geroch and Lindblom ('90) developed a general framework for theories of relativistic viscosity that leads to causal dynamics under many circumstances. One then has to show that a particular theory (e.g. MIS) fits in the formalism **under the conditions that give rise to causality.**

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– Thank you for your attention –