

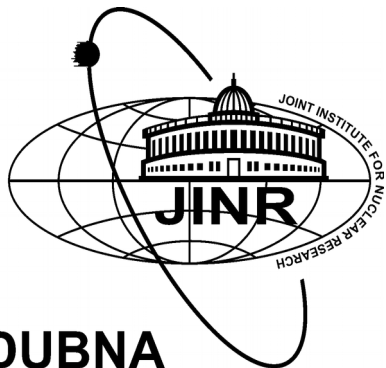
Recent news on hybrid stars & mass twins

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1. Introduction: “alternative facts” = fake news vs. reality check
2. Third family, twins, ABCD, HMT, triples, fifth family & all that:
 - Constant Speed of Sound (CSS) model
 - Multi-Polytrope (MP) model
 - Relativistic Density Functional (RDF) approach
3. Outlook: Full thermal hybrid EoS for Supernovae & Mergers

Astro-Coffee, FIAS, Goethe-University Frankfurt, 19.12.2017

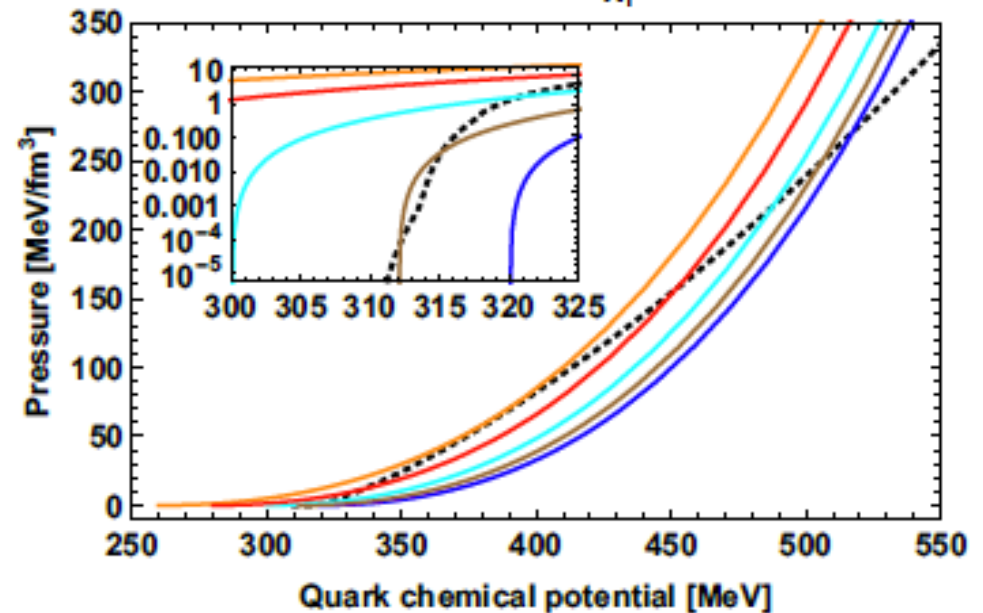
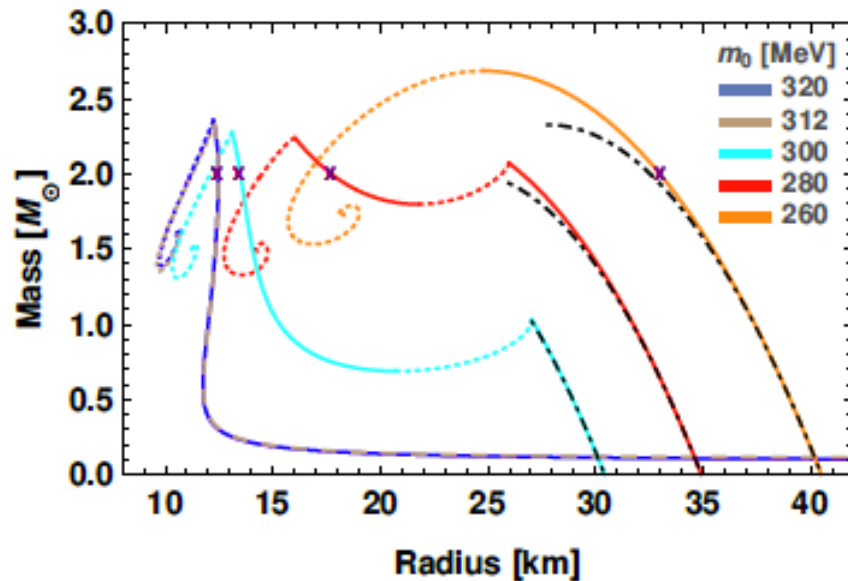
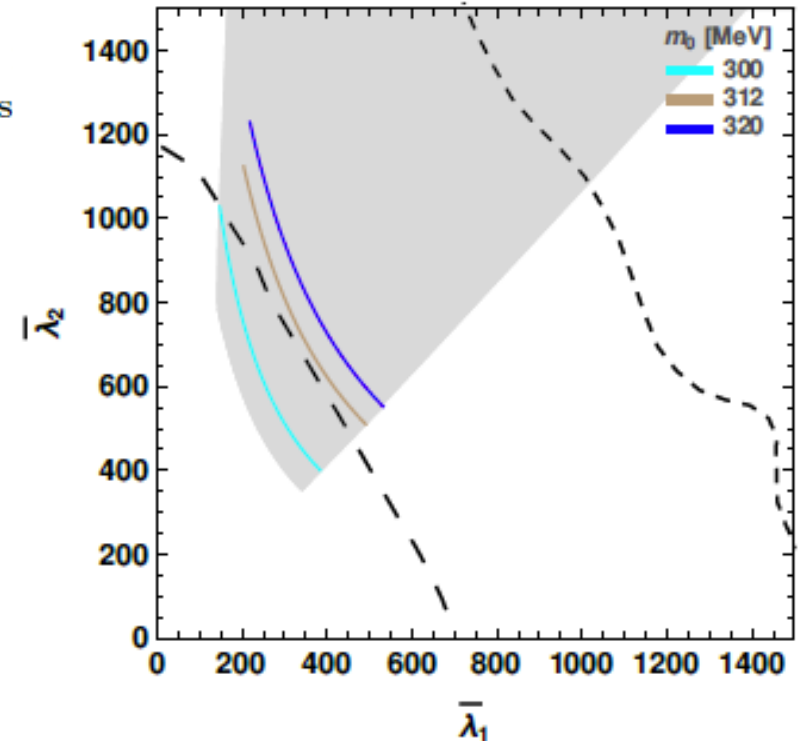
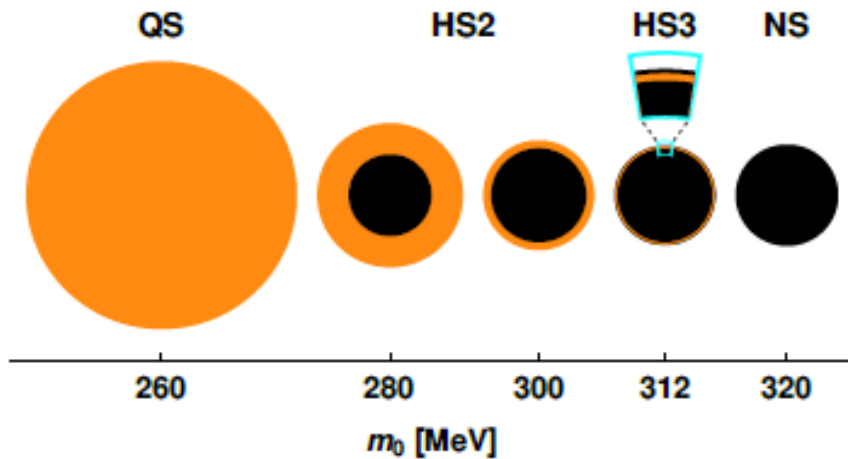


Alternative facts of the day: New hybrid star solutions!

arxiv:1711.06244v1, 1611.2017

Holographic compact stars meet gravitational wave constraints

Eemeli Annala,^{1,*} Christian Ecker,^{2,†} Carlos Hoyos,^{3,‡} Niko Jokela,^{1,§}
David Rodríguez Fernández,^{3,4,¶} and Alekski Vuorinen^{1,**}

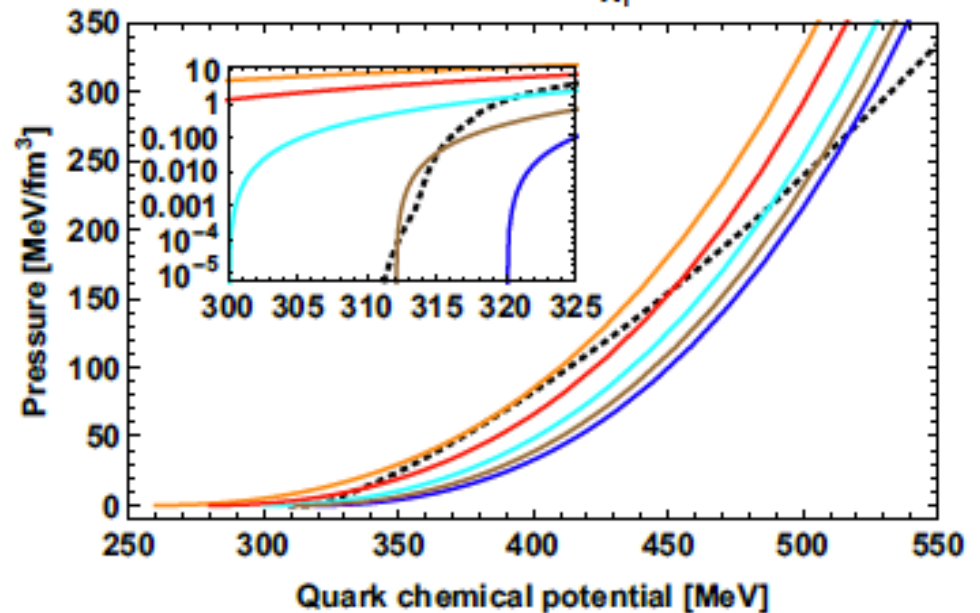
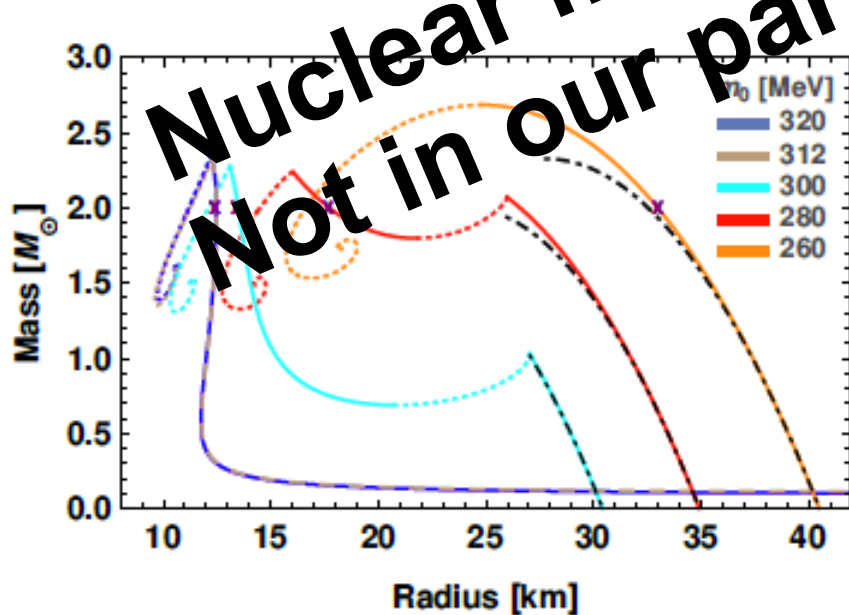
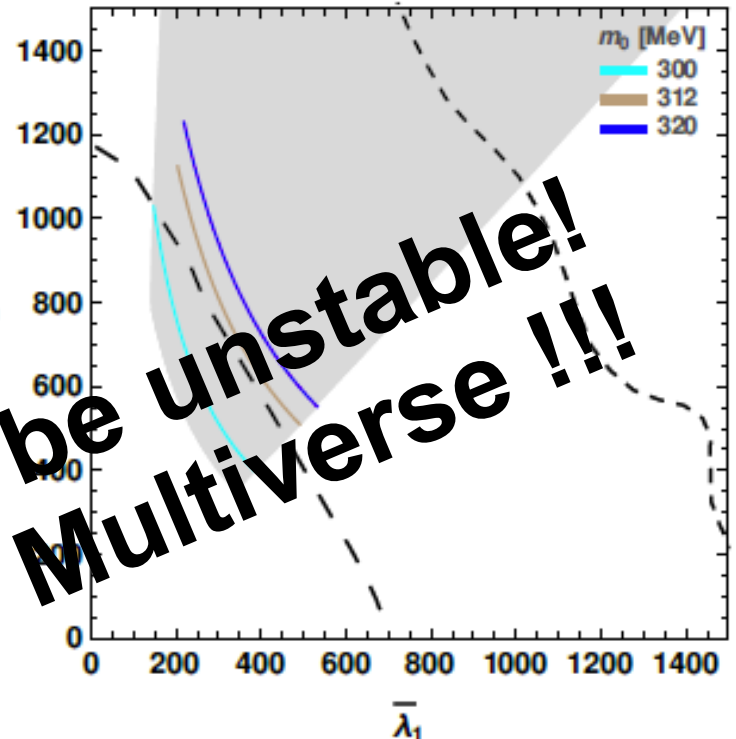
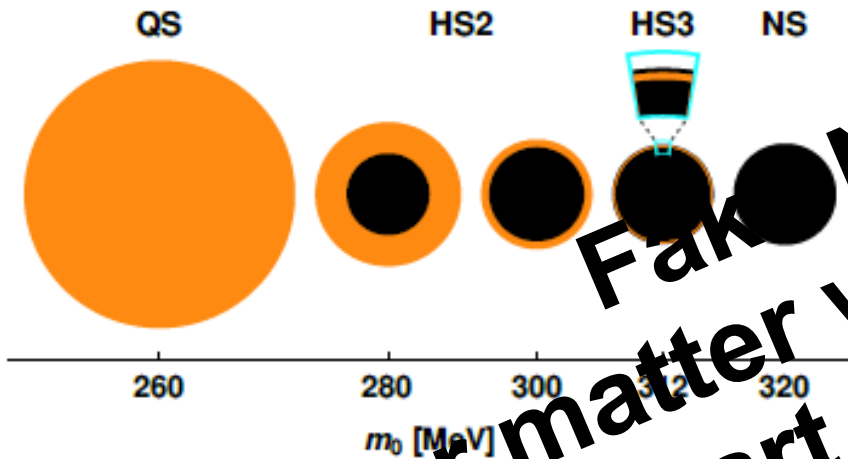


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Fake News!
Nuclear matter would be unstable!
Not in our part of the Multiverse !!!

History: Third family & Nonidentical Twins

PHYSICAL REVIEW

VOLUME 172, NUMBER 5

25 AUGUST 1968

Equation of State at Supranuclear Densities and the Existence of a Third Family of Superdense Stars*†

ULRICH H. GERLACH‡§

Palmer Physical Laboratory, Princeton University, Princeton, New Jersey

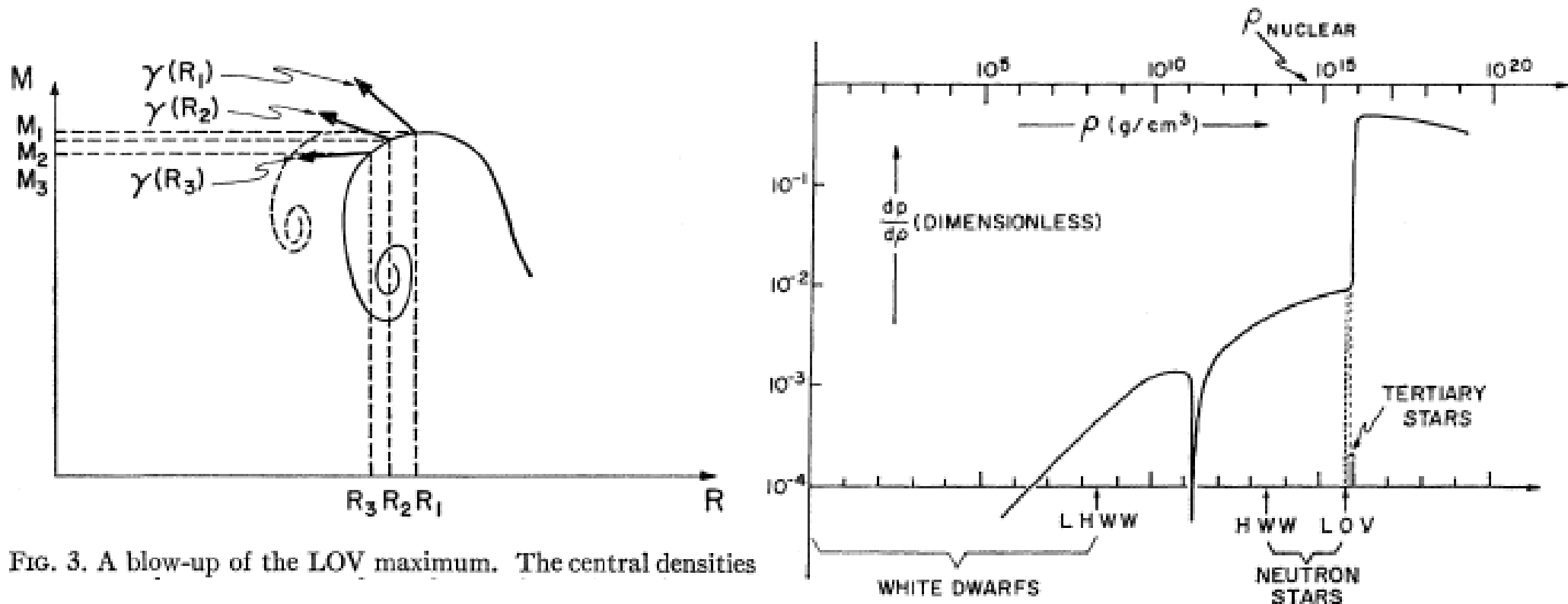


FIG. 3. A blow-up of the LOV maximum. The central densities

History: Third family & Nonidentical Twins

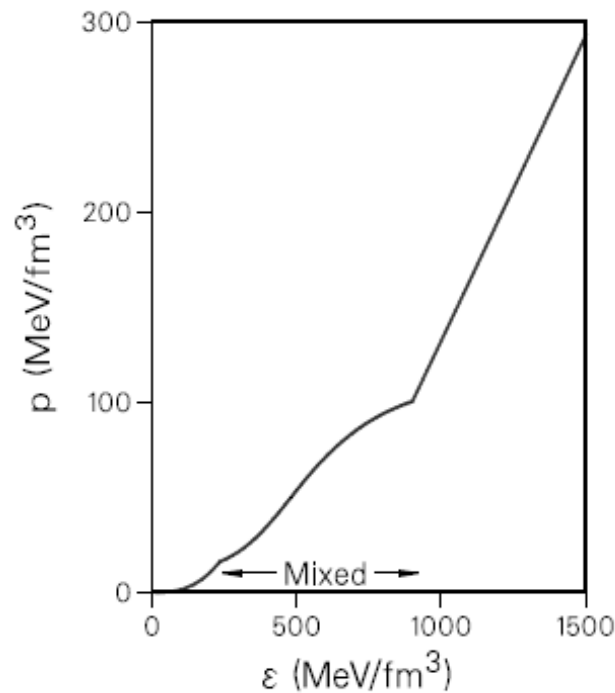
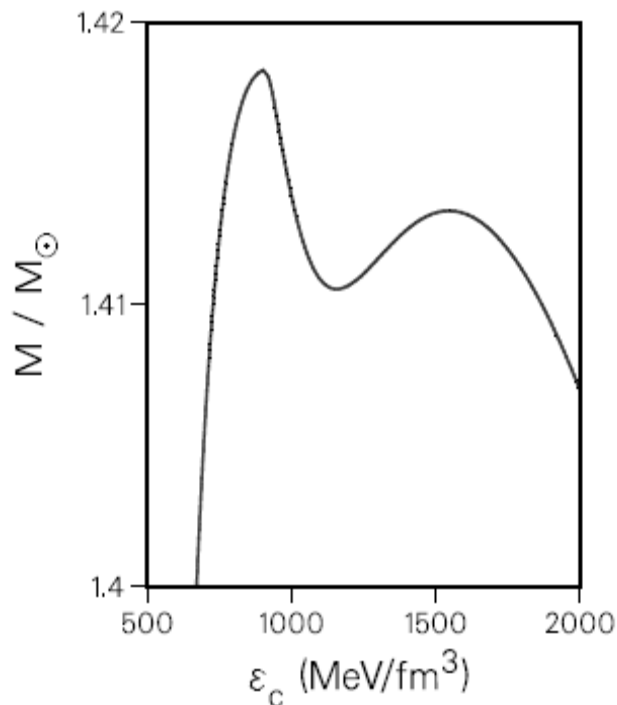
Non-Identical Neutron Star Twins

Norman K. Glendenning

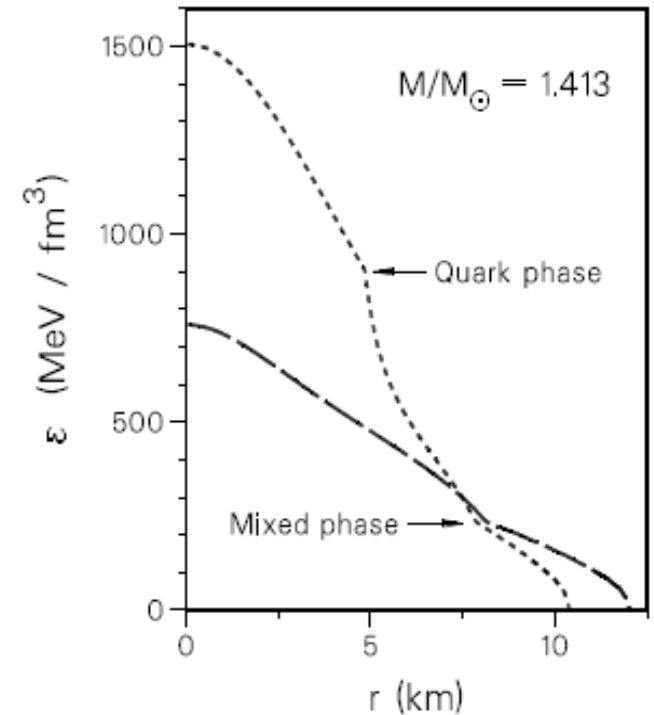
*Nuclear Science Division, Lawrence Berkeley National Laboratory,
University of California, Berkeley, CA 94720, USA*

Christiane Kettner

*Institut fuer theoretische Physik I, Universitaet Augsburg
Memmingerstr. 6, 86135 Augsburg
(June 17, 1998)*

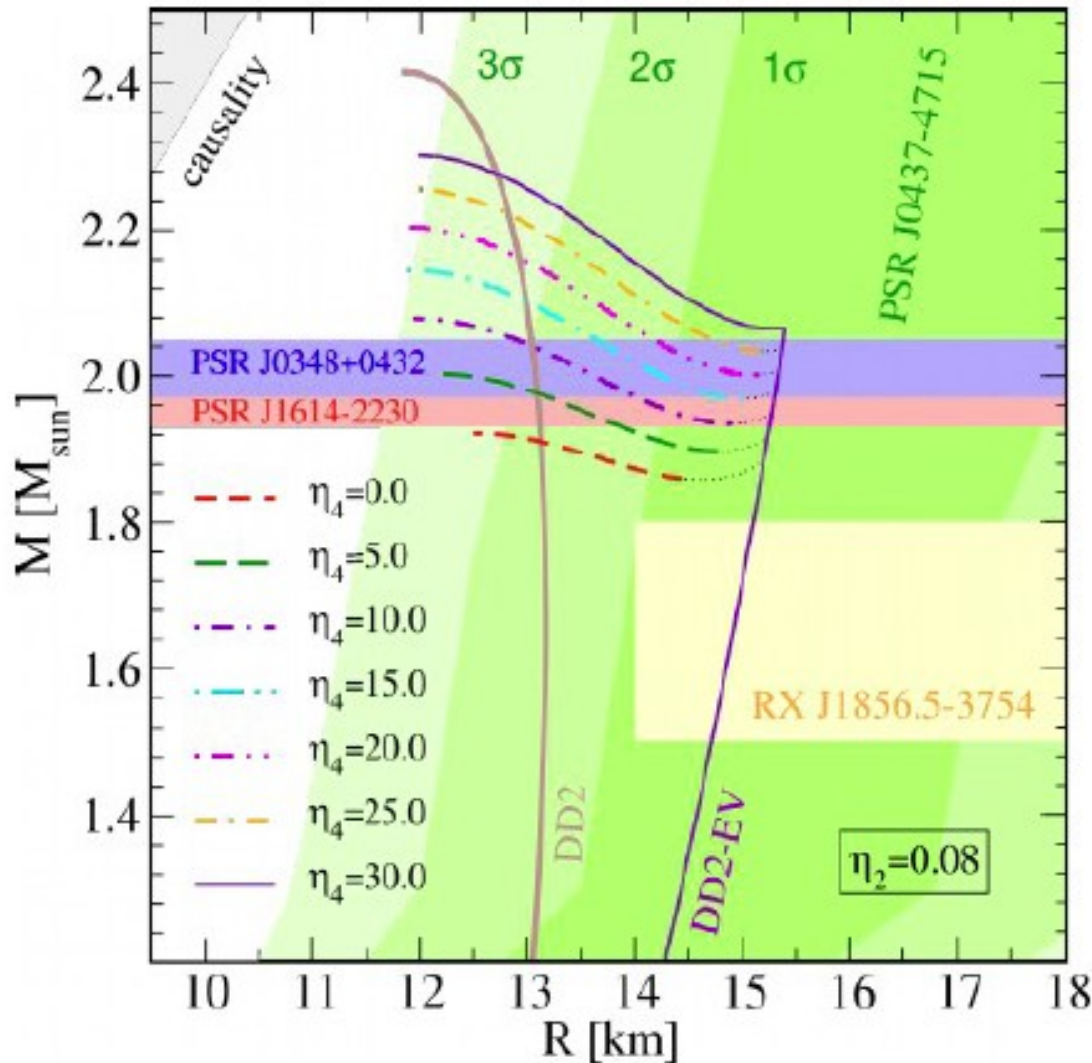


astro-ph/9807155; A&A (2000) L9

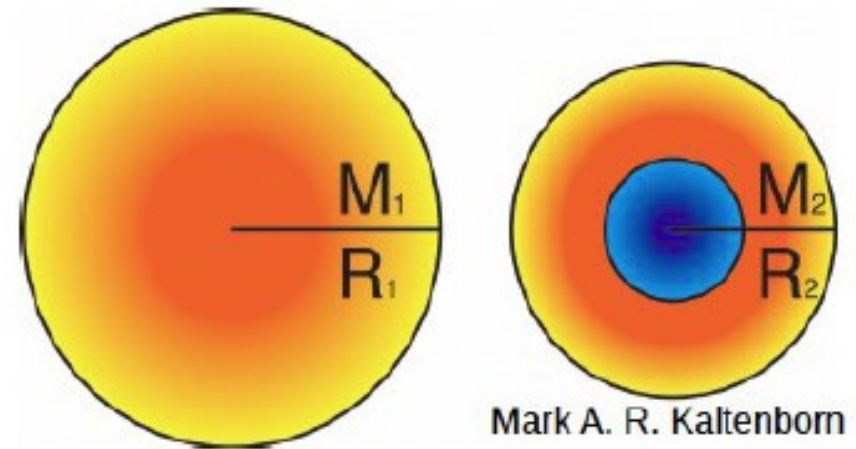


The original Twin paper uses
Glendenning construction, not
Maxwell one -
Surface tension zero vs. infity!
Pasta phases in-between ...

Neutron Star Interiors: Strong Phase Transition?



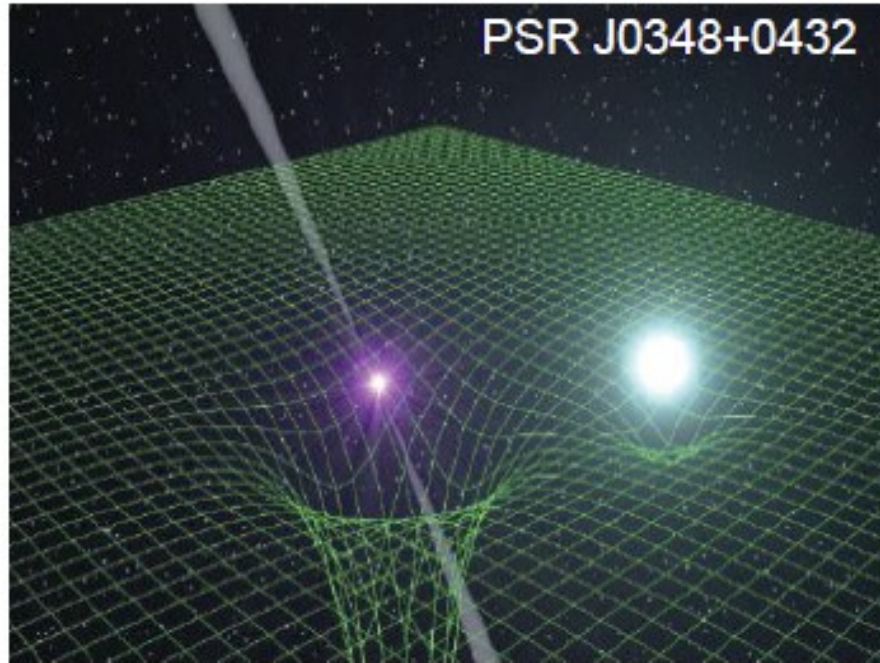
- Star configurations with same masses, but different radii



- **New class of EOS, that features high mass twins**
- NASA NICER mission: radii measurements ~ 0.5 km
- Existence of twins implies 1st order phase-transition and hence a critical point

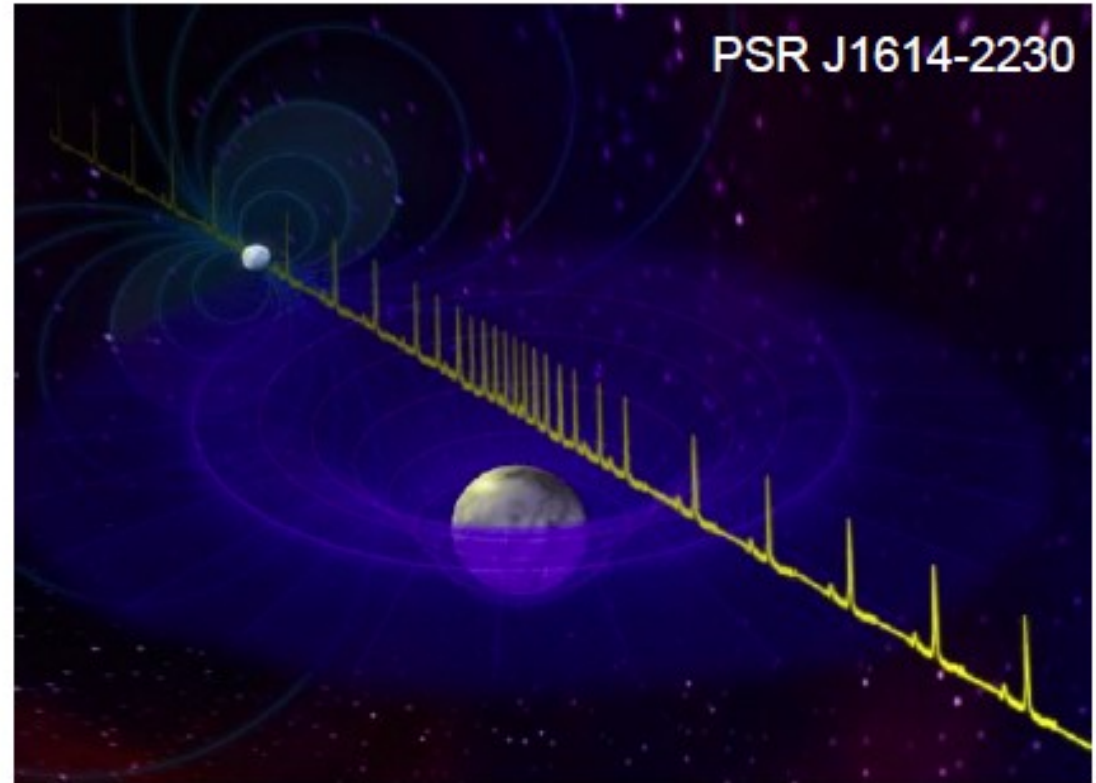
Neutron Star Interiors: Strong Phase Transition?

$M=2.01 \pm 0.04 M_{\text{sun}}$



Antoniadis et al., Science 340 (2013) 448
Demorest et al., Nature 467 (2010) 1081
Fonseca et al., arxiv:1603.00545

$M=1.928 \pm 0.017 M_{\text{sun}}$

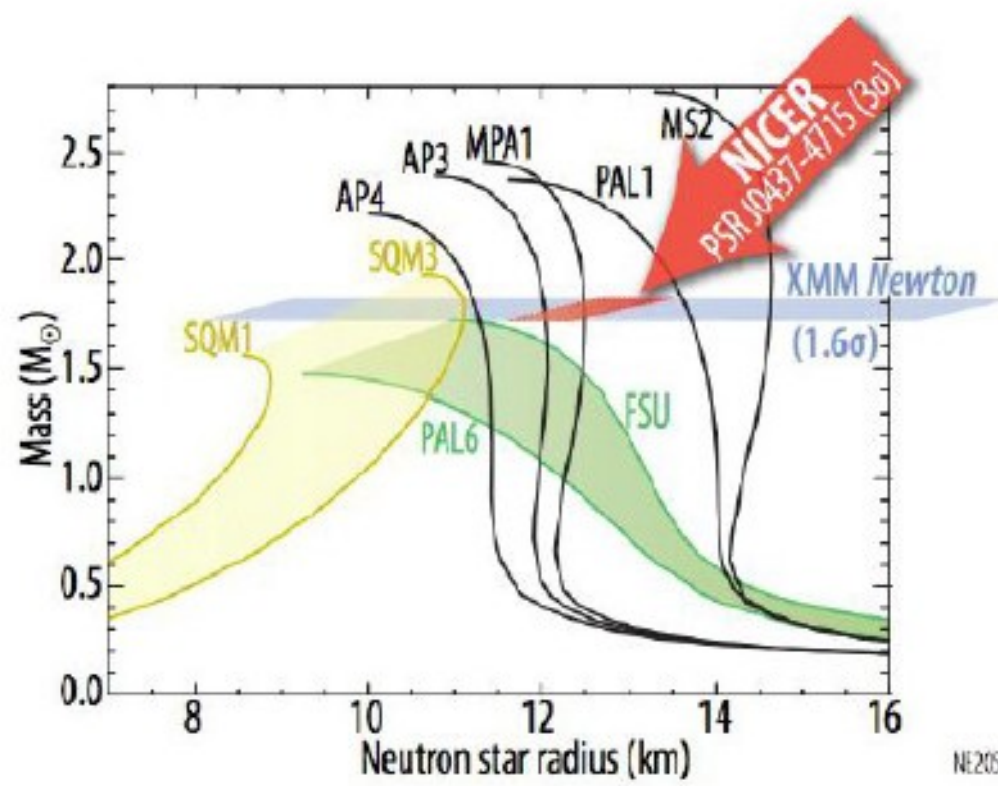
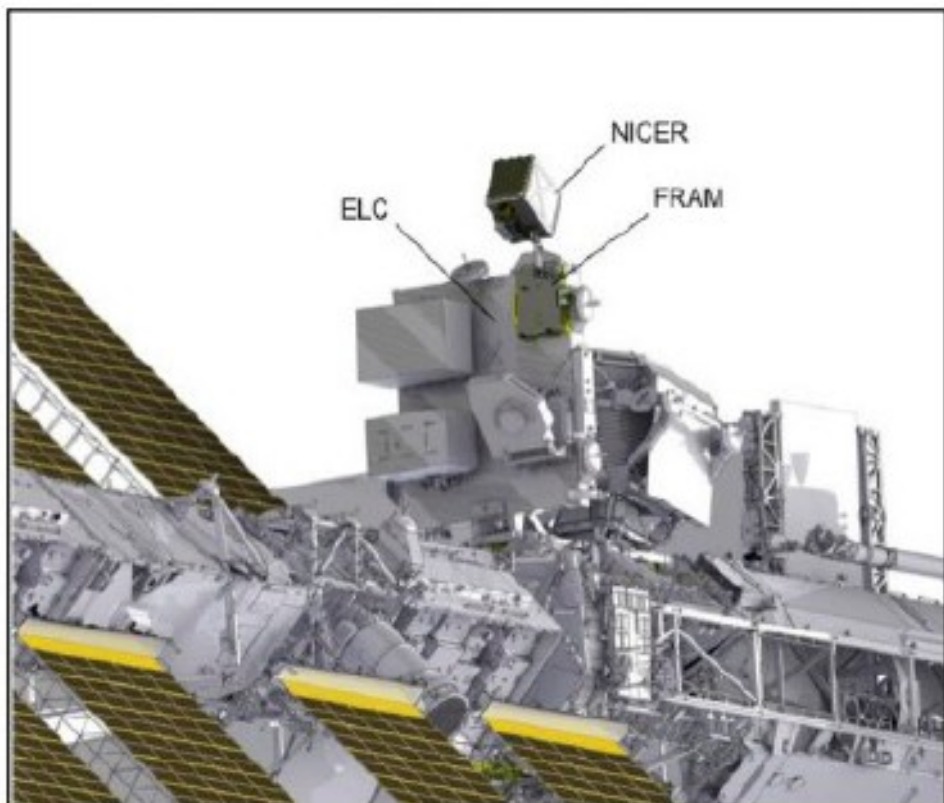
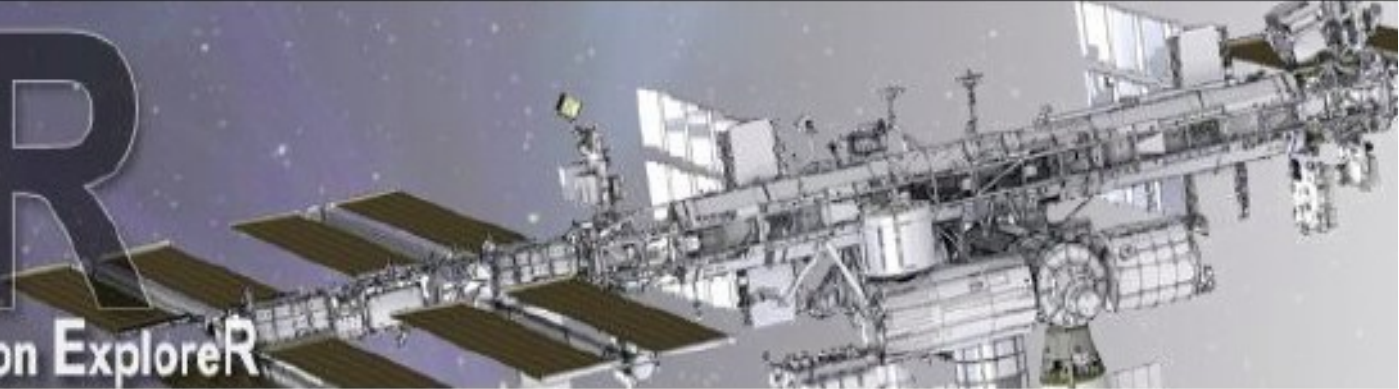


What if they were high-mass twin stars?

→ radius measurement required ! → NICER (2017)

NICER

Neutron star Interior Composition Explorer

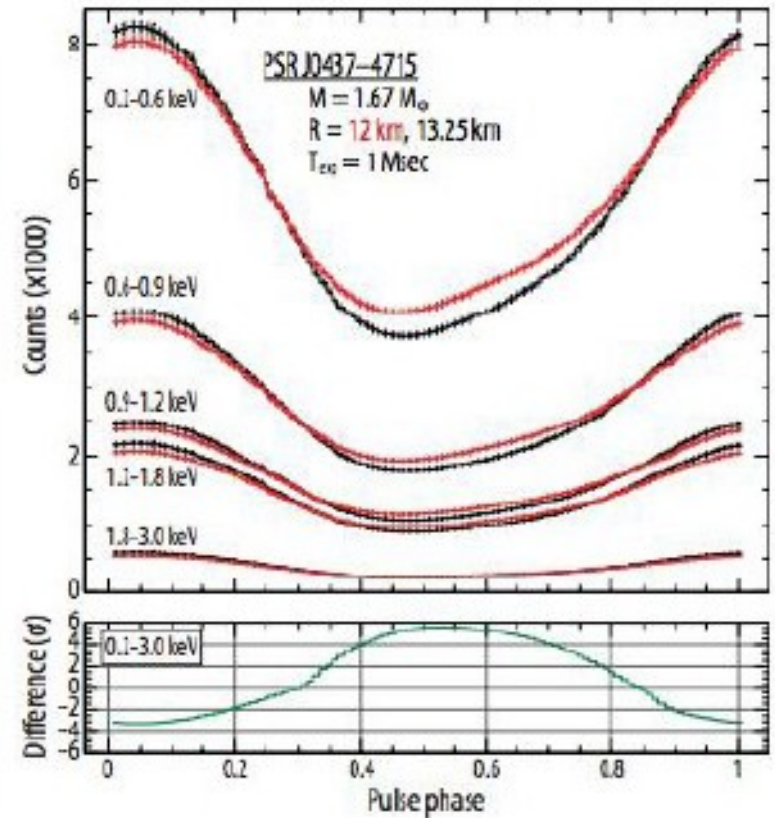
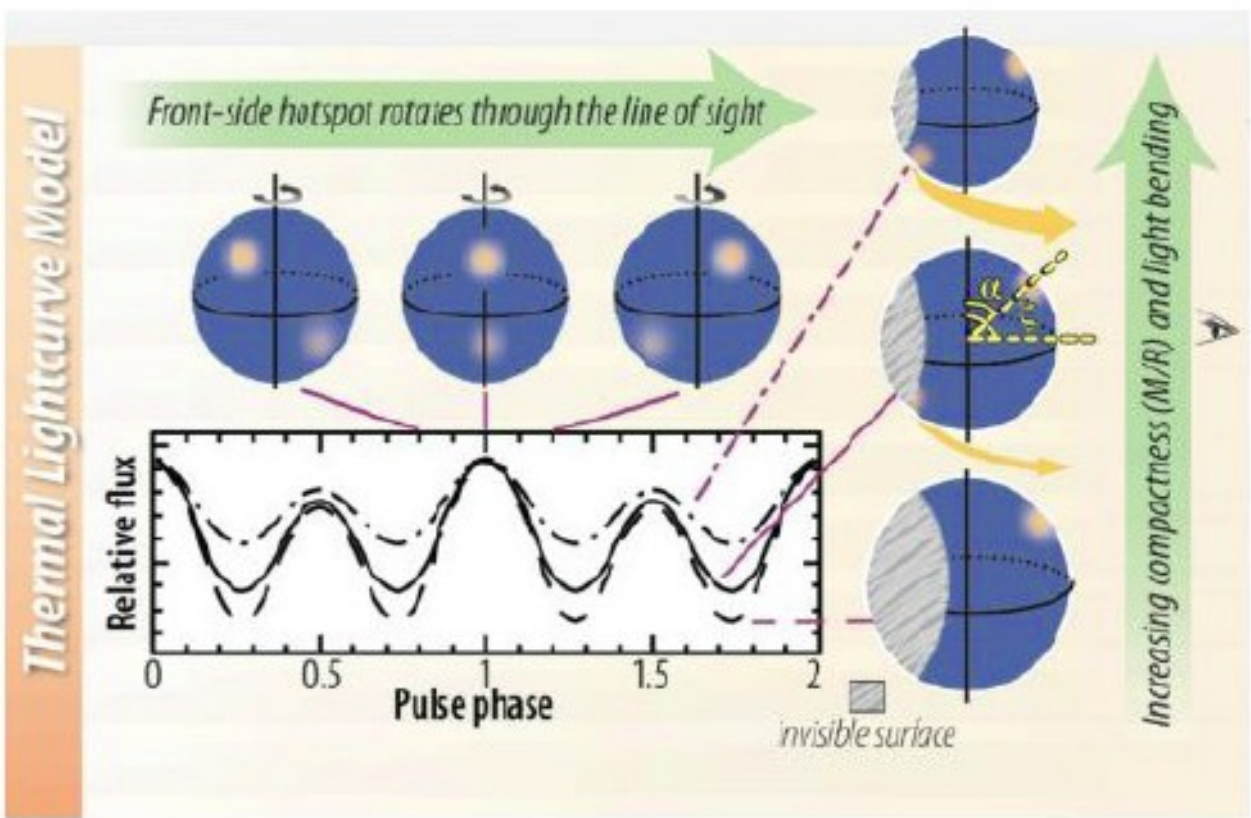
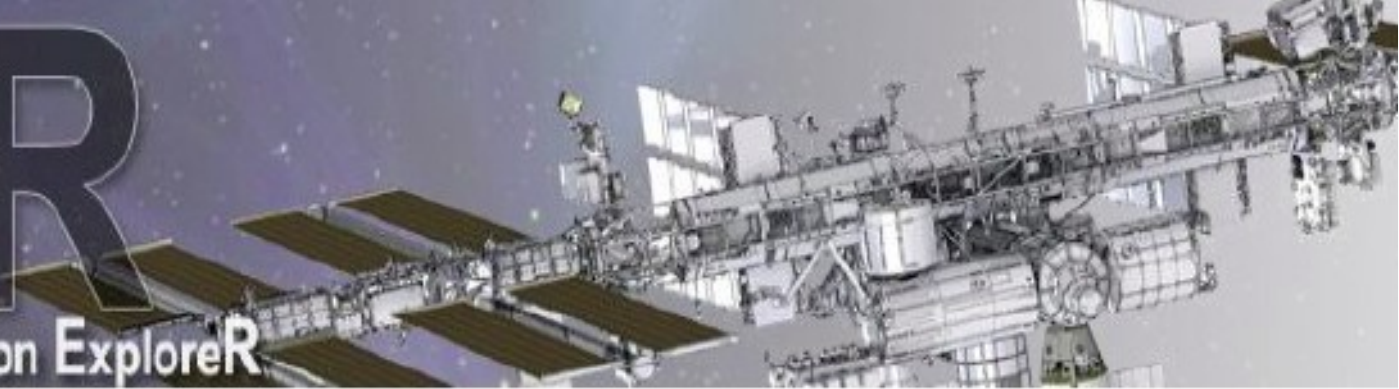


NICER 2017

Gendreau, K. C., Arzoumanian, Z., & Okajima, T. 2012, Proc. SPIE, 8443, 844313

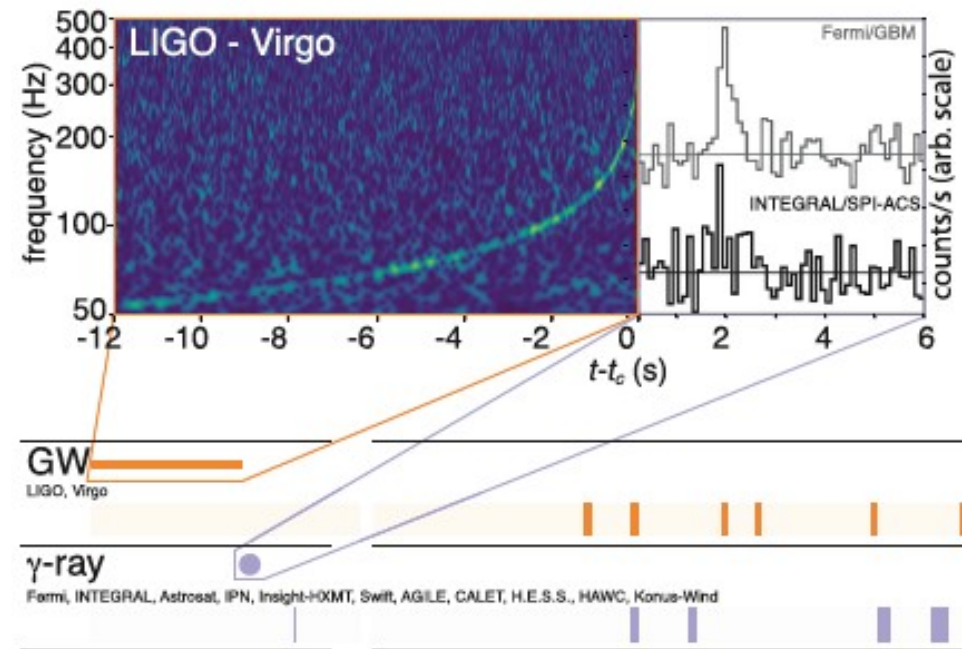
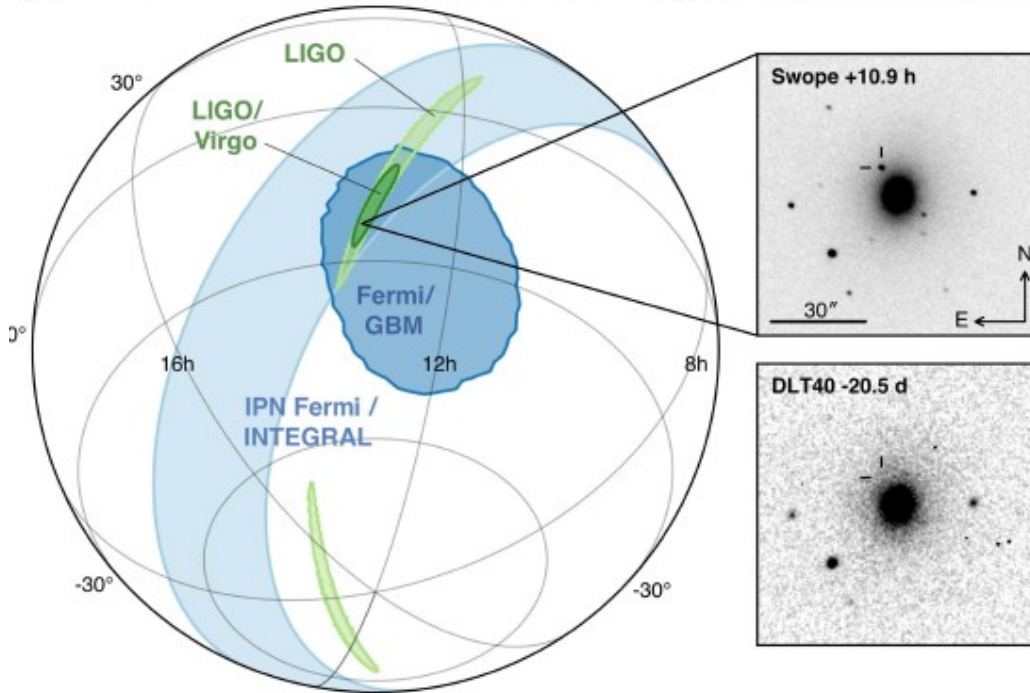
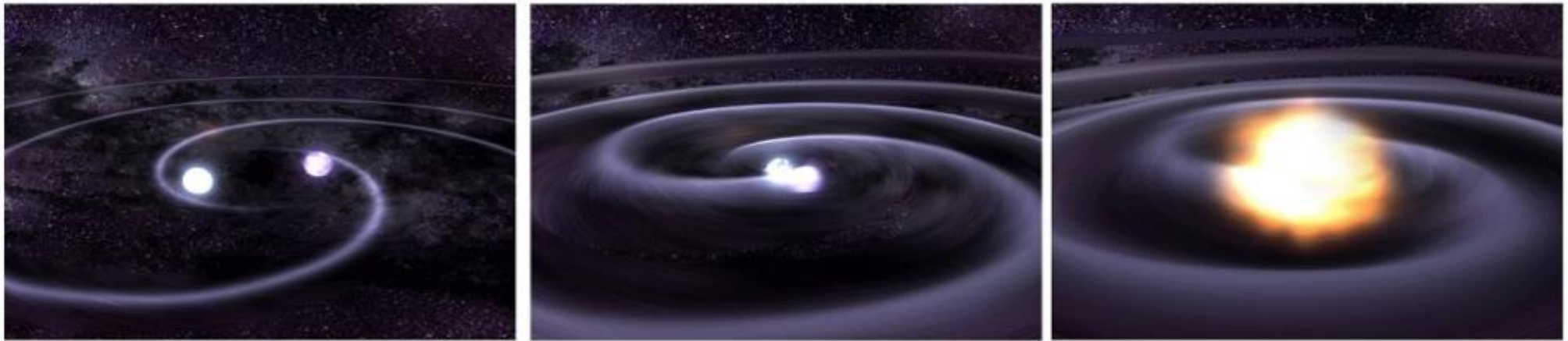
NICER

Neutron star Interior Composition Explorer



Hot Spots

Mass-Radius Constraints: GW170817 & NICER



GW170817, announced on 16.10.2017

B.P. Abbott et al. [LIGO/Virgo Collab.], PRL 119, 161101 (2017); ApJLett 848, L12 (2017)

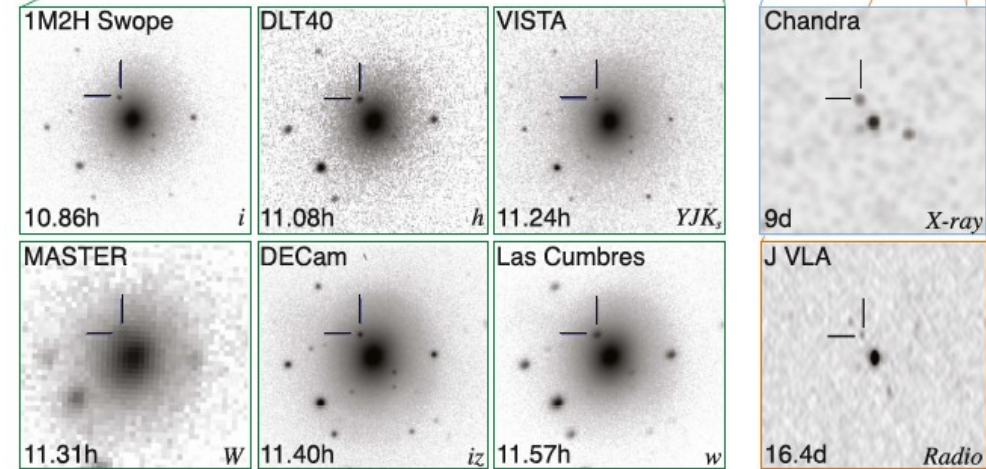
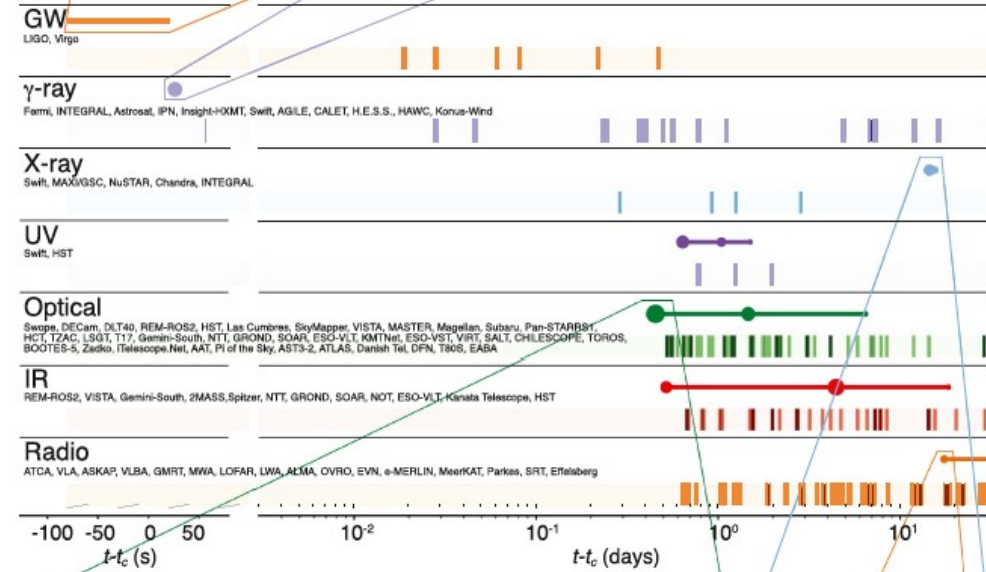
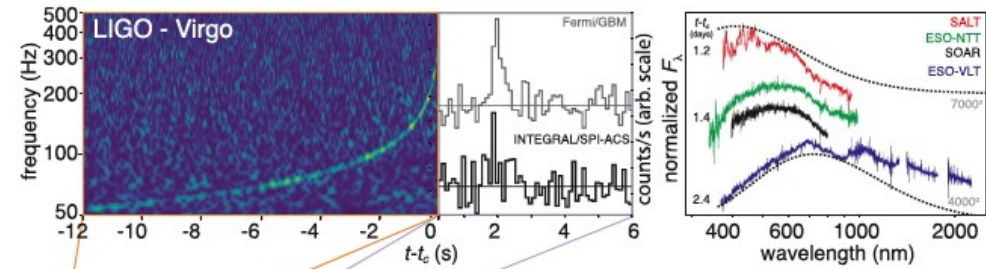
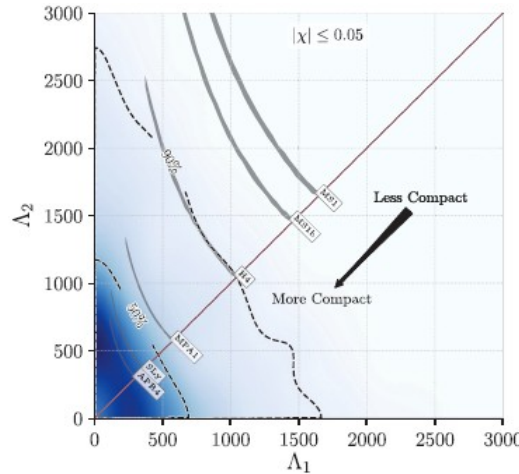
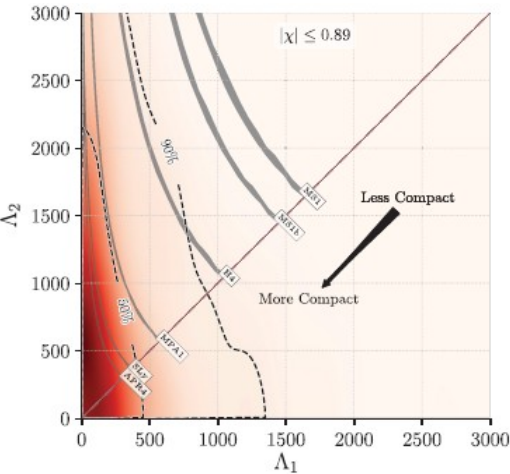
GW170817: NS-NS Merger

Multi-Messenger Astrophysics !!

$M < 2.17 M_{\text{sun}}$ (arxiv:1710.05938)

Low-spin priors ($|\chi| \leq 0.05$)

Primary mass m_1	$1.36\text{--}1.60 M_{\odot}$
Secondary mass m_2	$1.17\text{--}1.36 M_{\odot}$
Chirp mass \mathcal{M}	$1.188^{+0.004}_{-0.002} M_{\odot}$
Mass ratio m_2/m_1	$0.7\text{--}1.0$
Total mass m_{tot}	$2.74^{+0.04}_{-0.01} M_{\odot}$
Radiated energy E_{rad}	$> 0.025 M_{\odot} c^2$
Luminosity distance D_L	40^{+8}_{-14} Mpc



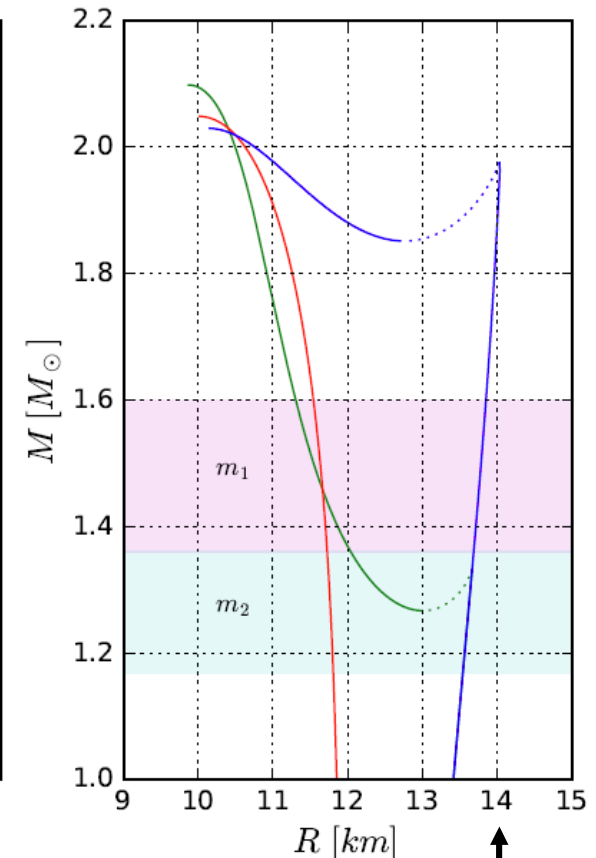
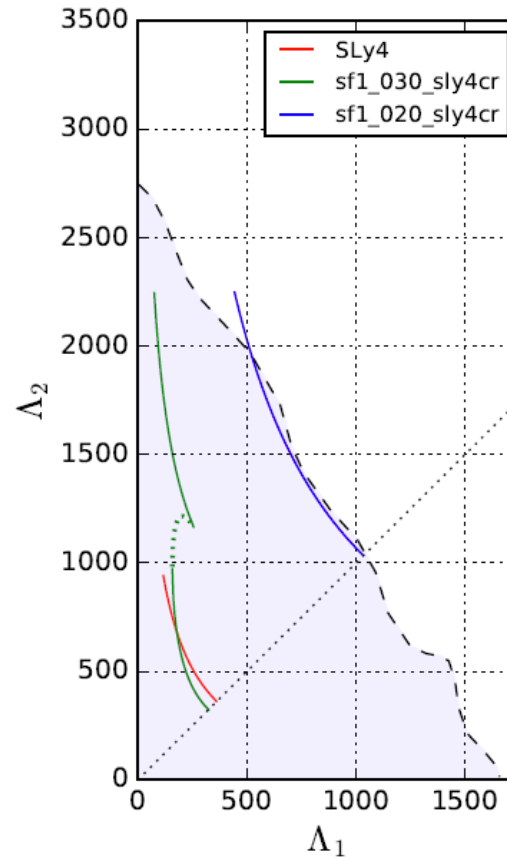
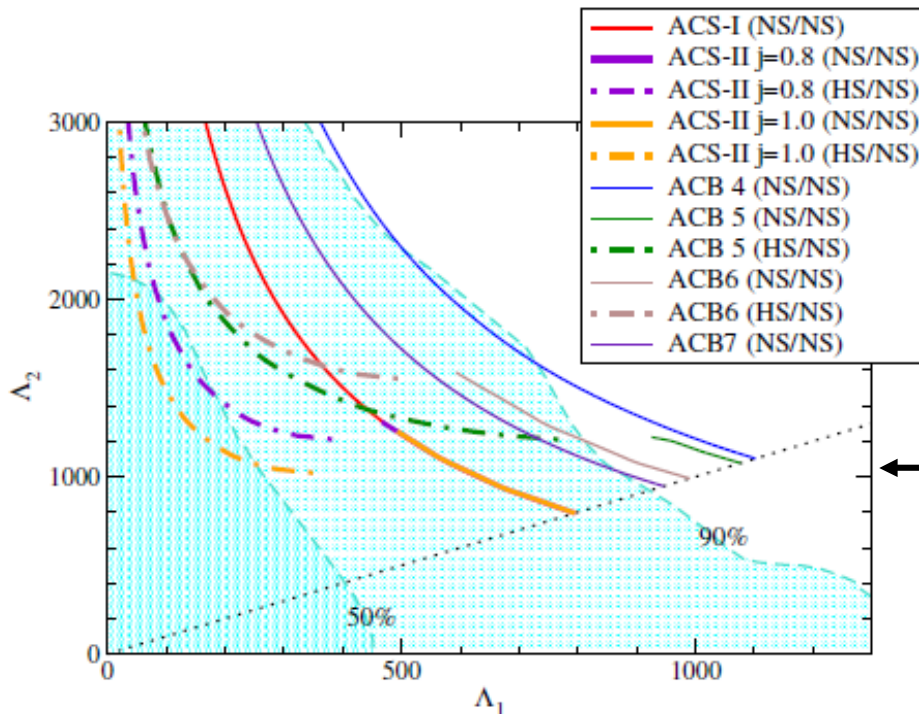
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GW170817: NS-NS Merger – Equation of State Constraints

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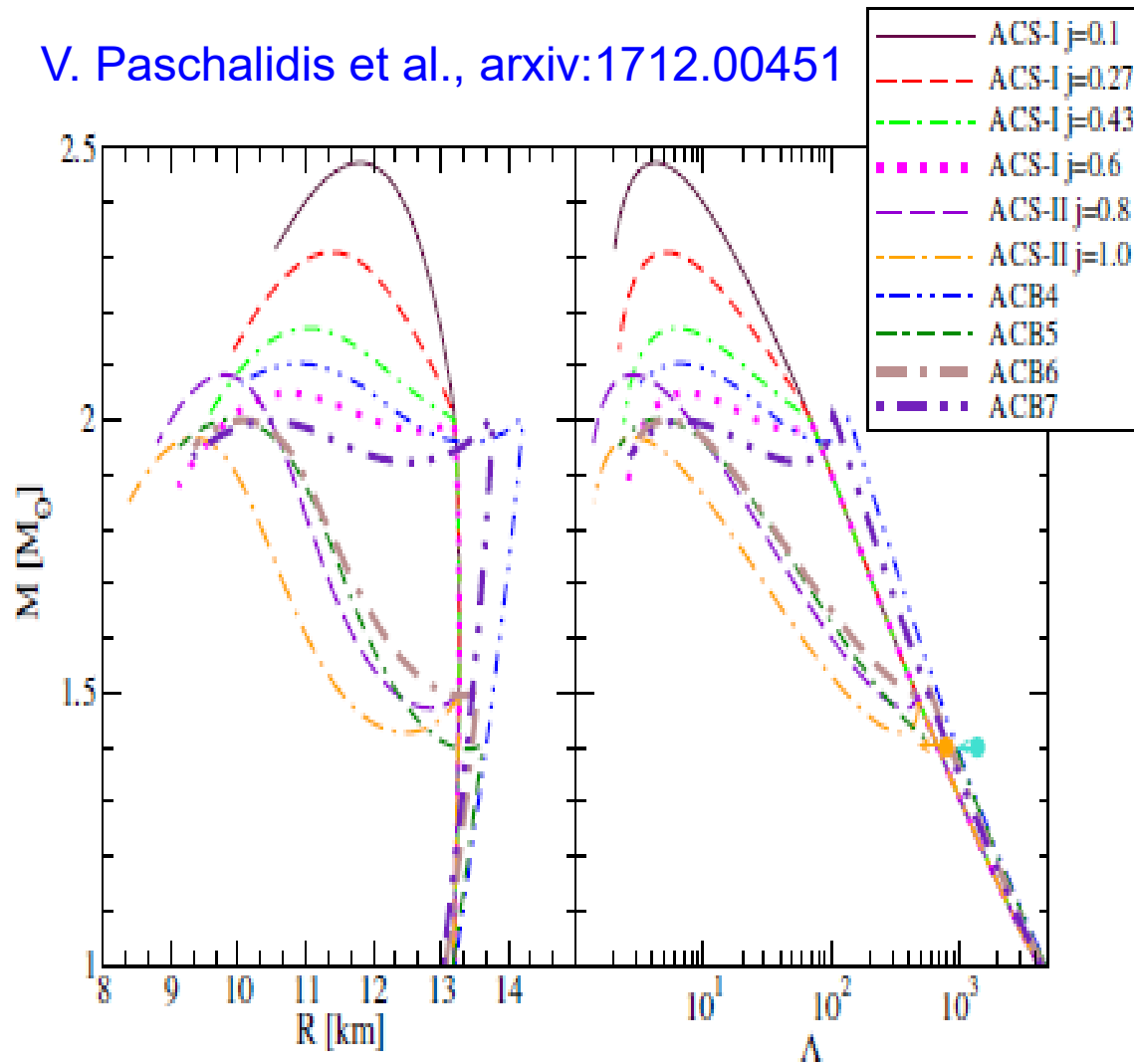
M. Bejger, D.B., et al., in preparation (2017)

V. Paschalidis, K. Yagi, D. Alvarez-Castillo, D.B., A. Sedrakian, arxiv:1712.00451 (2017)

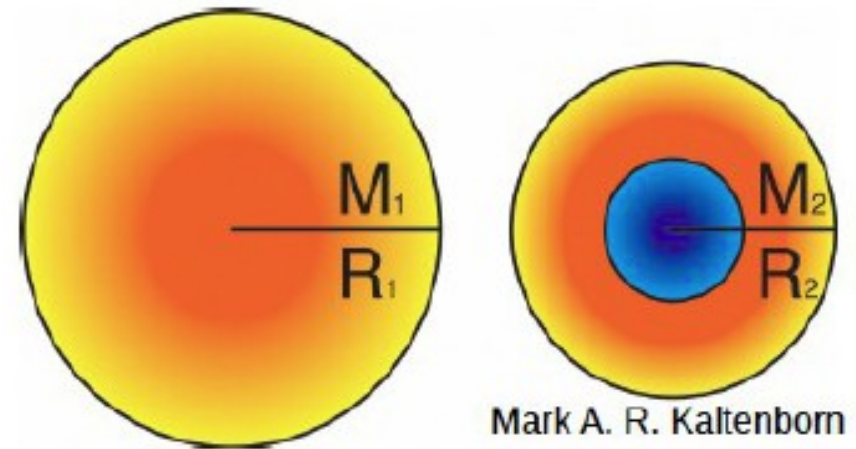
Suggestion: The heavier NS be a hybrid star (HS) with a quark core, evtl. member of a “third family”!

Neutron Star Interiors: Strong Phase Transition? M-R Relation!

V. Paschalidis et al., arxiv:1712.00451



- Star configurations with same masses, but different radii



- **New class of EOS, that features high mass twins**
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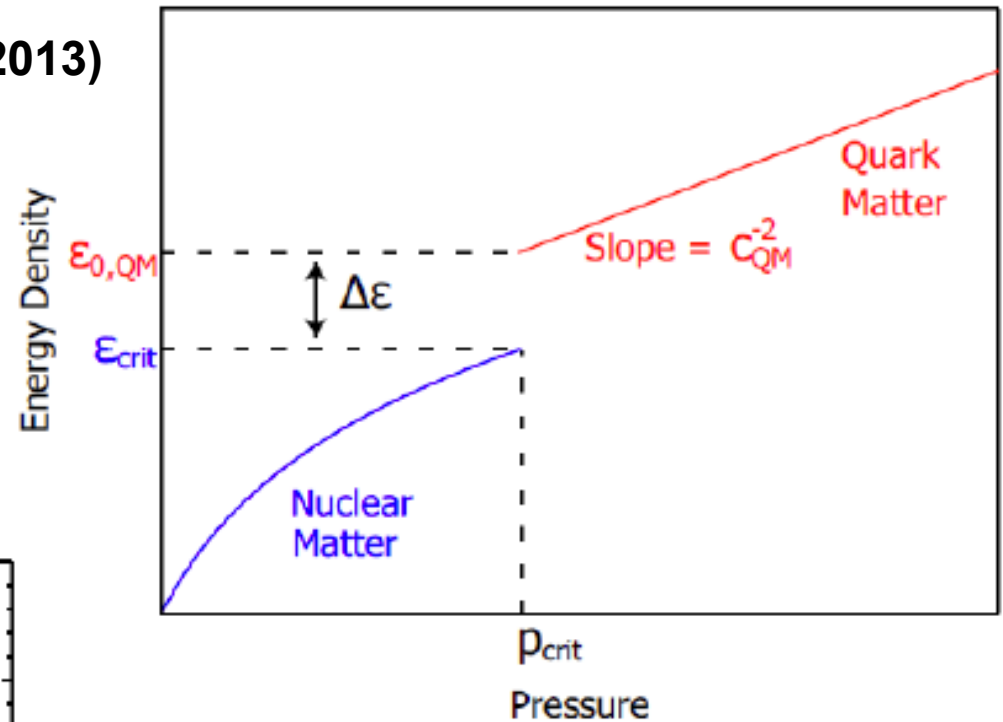
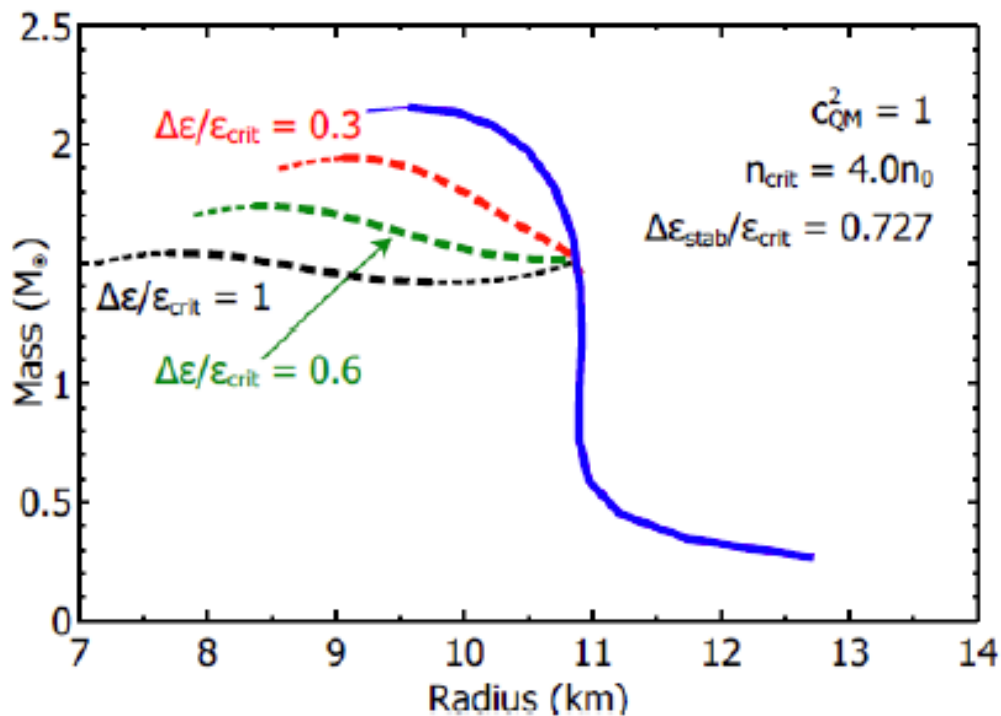
High-mass twins (HMT) or typical-mass twins (TMT) ?

For a classification see: J.-E. Christian, A. Zacchi, J. Schaffner-Bielich, arxiv:1707.07524

2.1 Constant Speed of Sound (CSS) Model

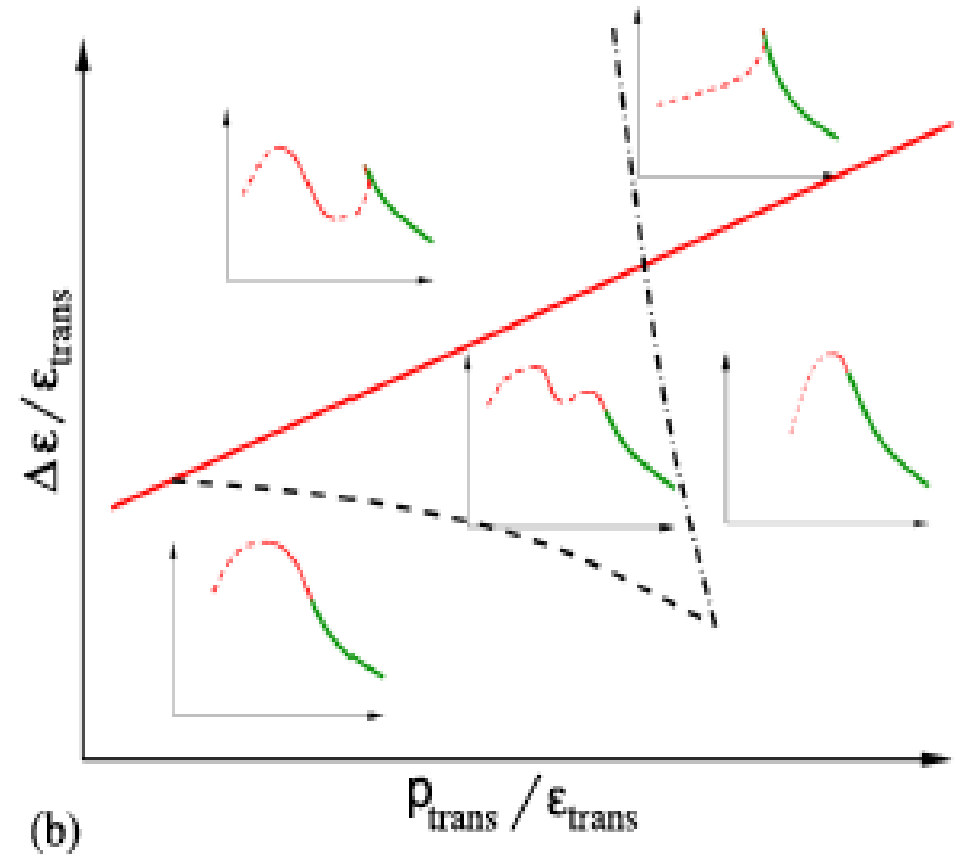
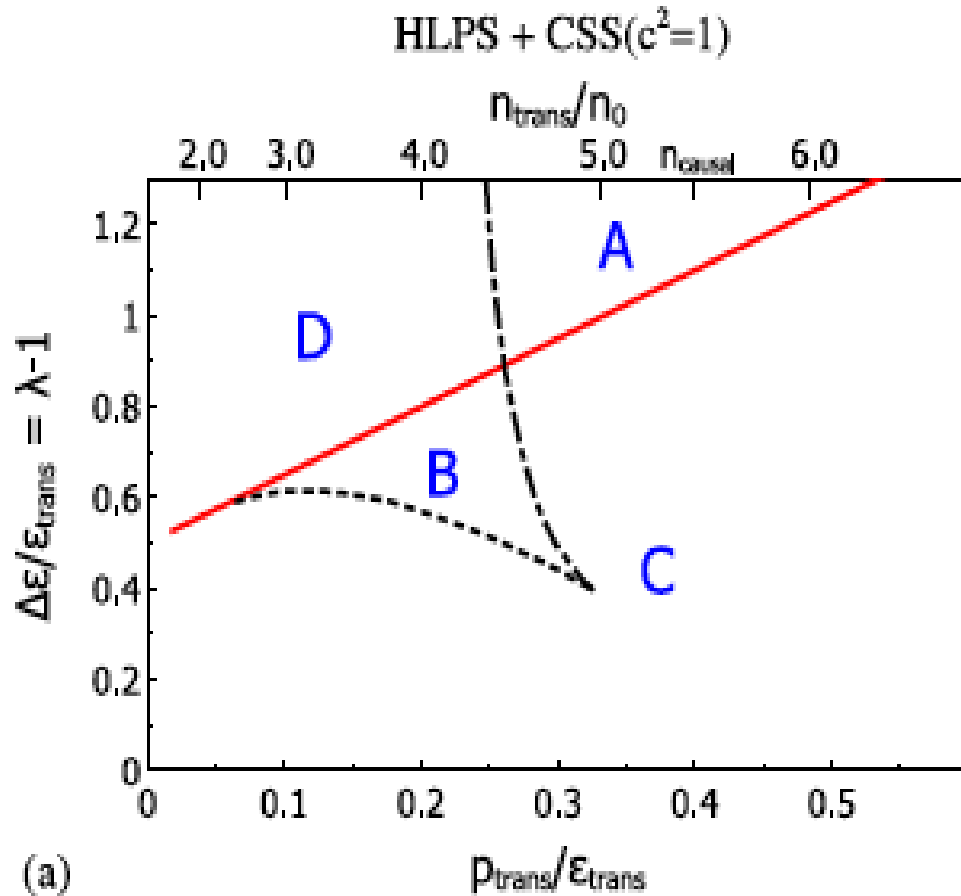
Alford, Han, Prakash, PRD88, 013083 (2013)

First order PT can lead to a stable branch of hybrid stars with quark matter cores which, depending on the size of the “latent heat” (jump in energy density), can even be disconnected from the hadronic one by an unstable branch → “third family of CS”.



Measuring two disconnected populations of compact stars in the M-R diagram would be the detection of a first order phase transition in compact star matter and thus the indirect proof for the existence of a critical endpoint (CEP) in the QCD phase diagram!

Key fact: Mass “twins” \leftrightarrow 1st order PT



Systematic Classification [Alford, Han, Prakash: PRD88, 083013 (2013)]

EoS $P(\epsilon) \leftrightarrow$ Compact star phenomenology $M(R)$

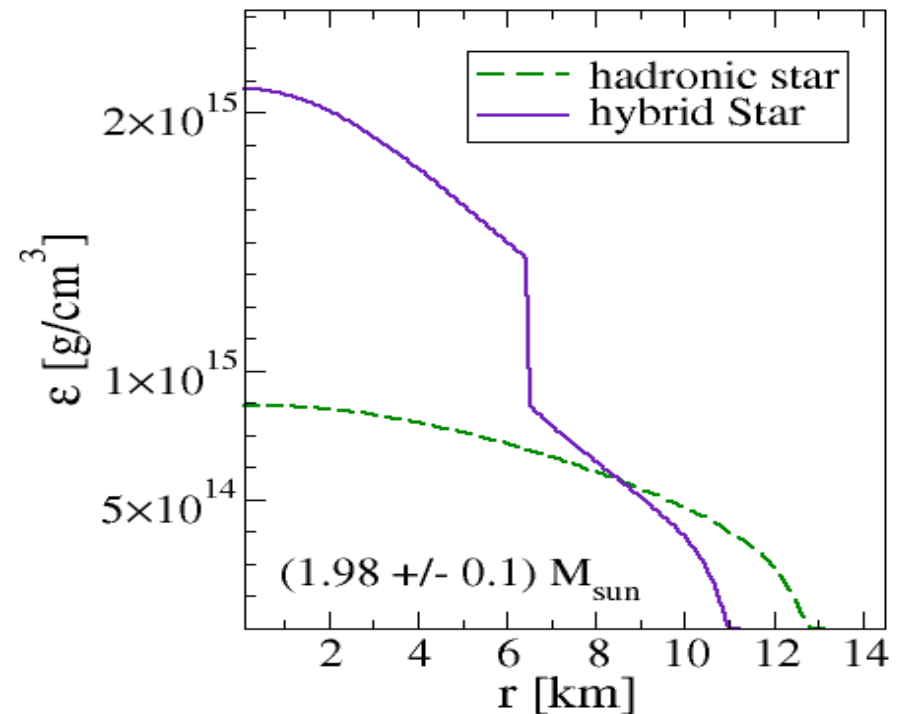
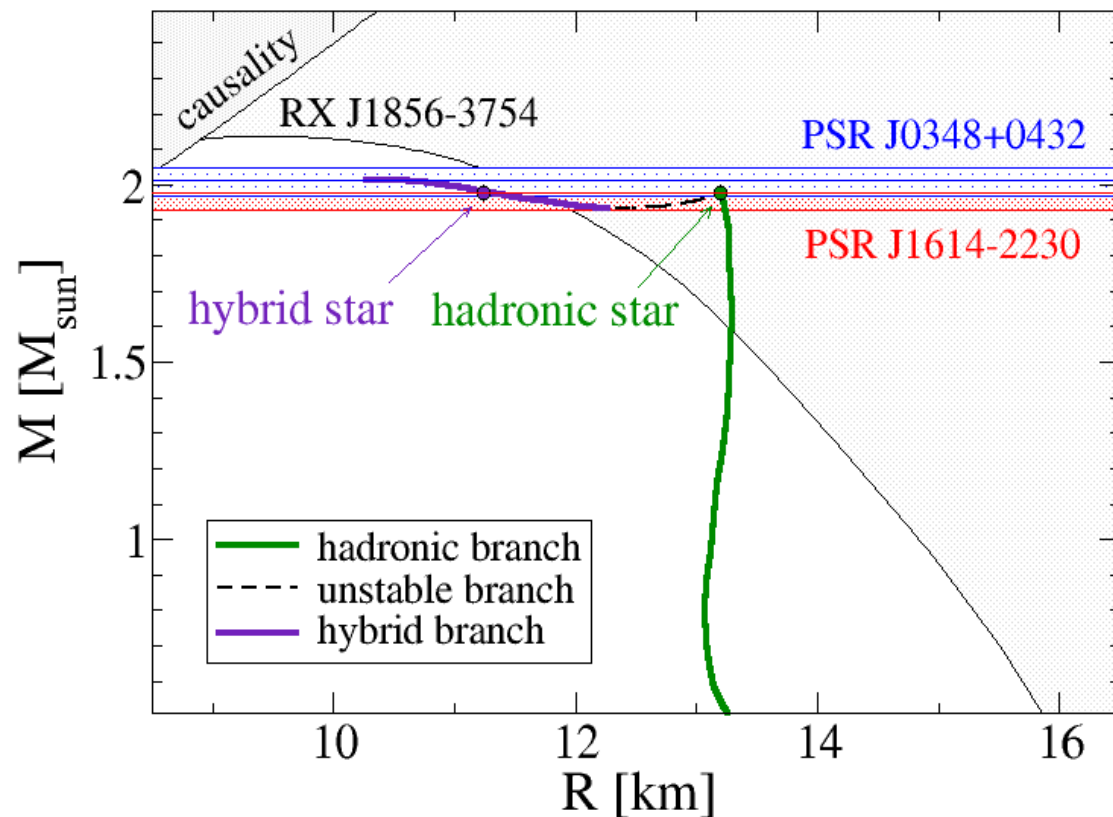
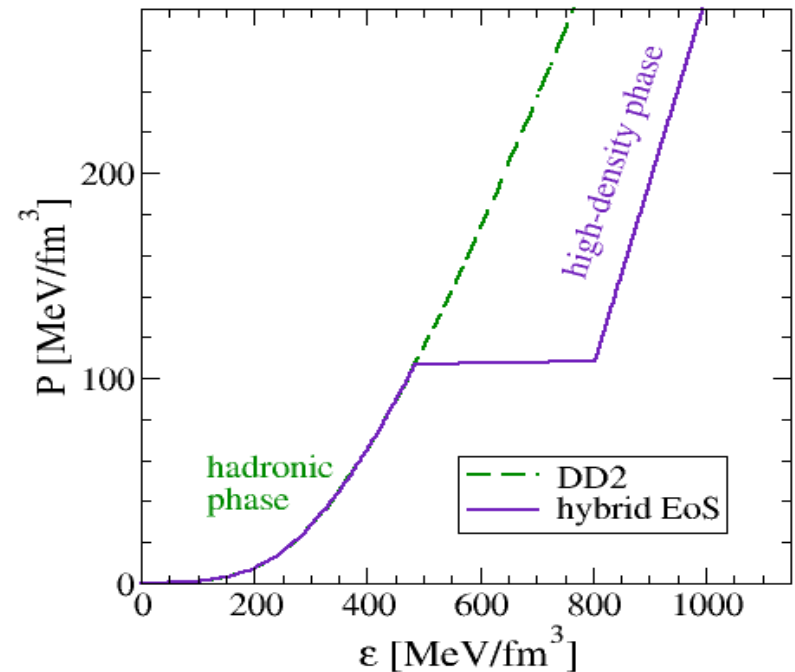
Most interesting and clear-cut cases: (D)isconnected and (B)oth – high-mass twins!

“Holy Grail” - High-Mass Twin Stars

Twins prove existence of **disconnected populations** (third family) in the M-R diagram

Consequence of a **first order phase transition**

Question: Do twins prove the 1st order phase trans.?

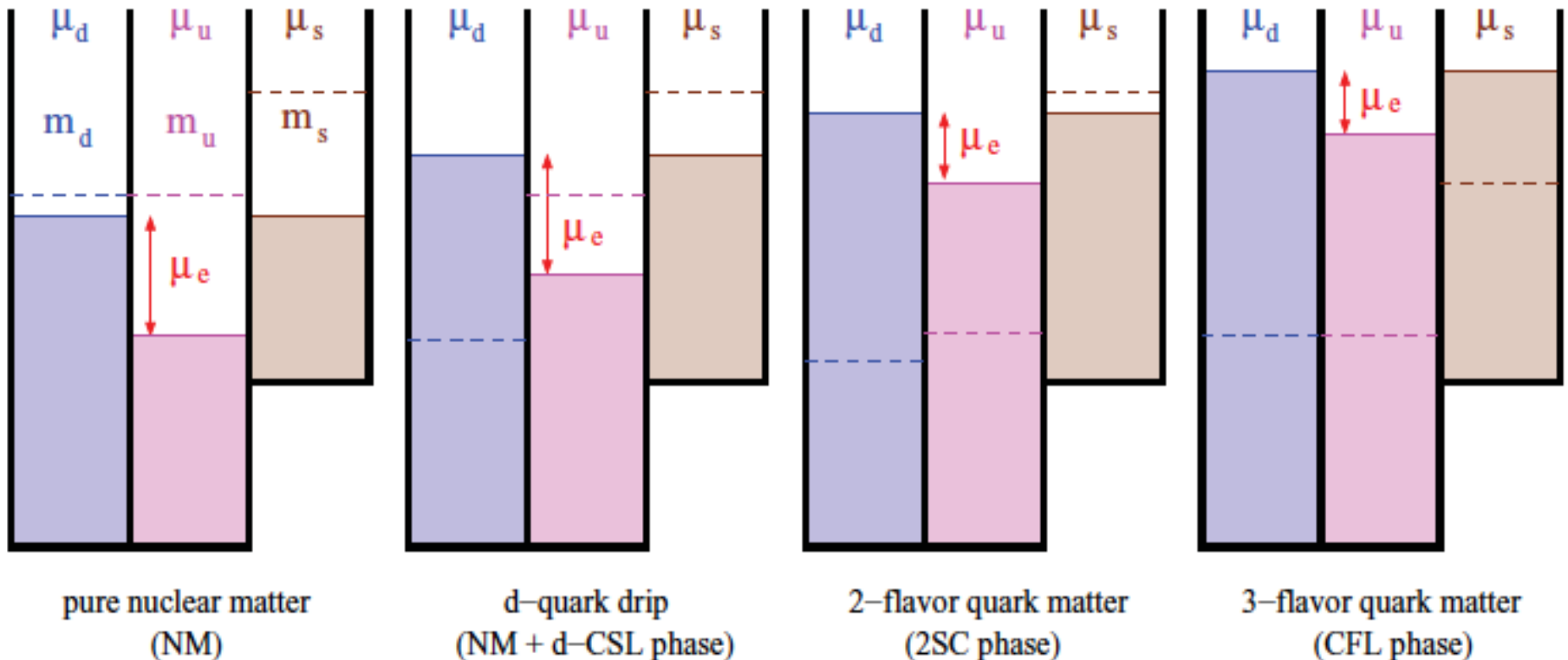


Neutron Star Interiors: Sequential Phase Transitions?

How likely is it that s-quarks (and no s-bar) exist and survive in neutron stars in a QGP or in hyperons. How large is then the ratio $s/(u+d)$ in neutron stars and in the Universe?

There could also be single flavor quark matter, mixed with nuclear matter (d-quark dripline)

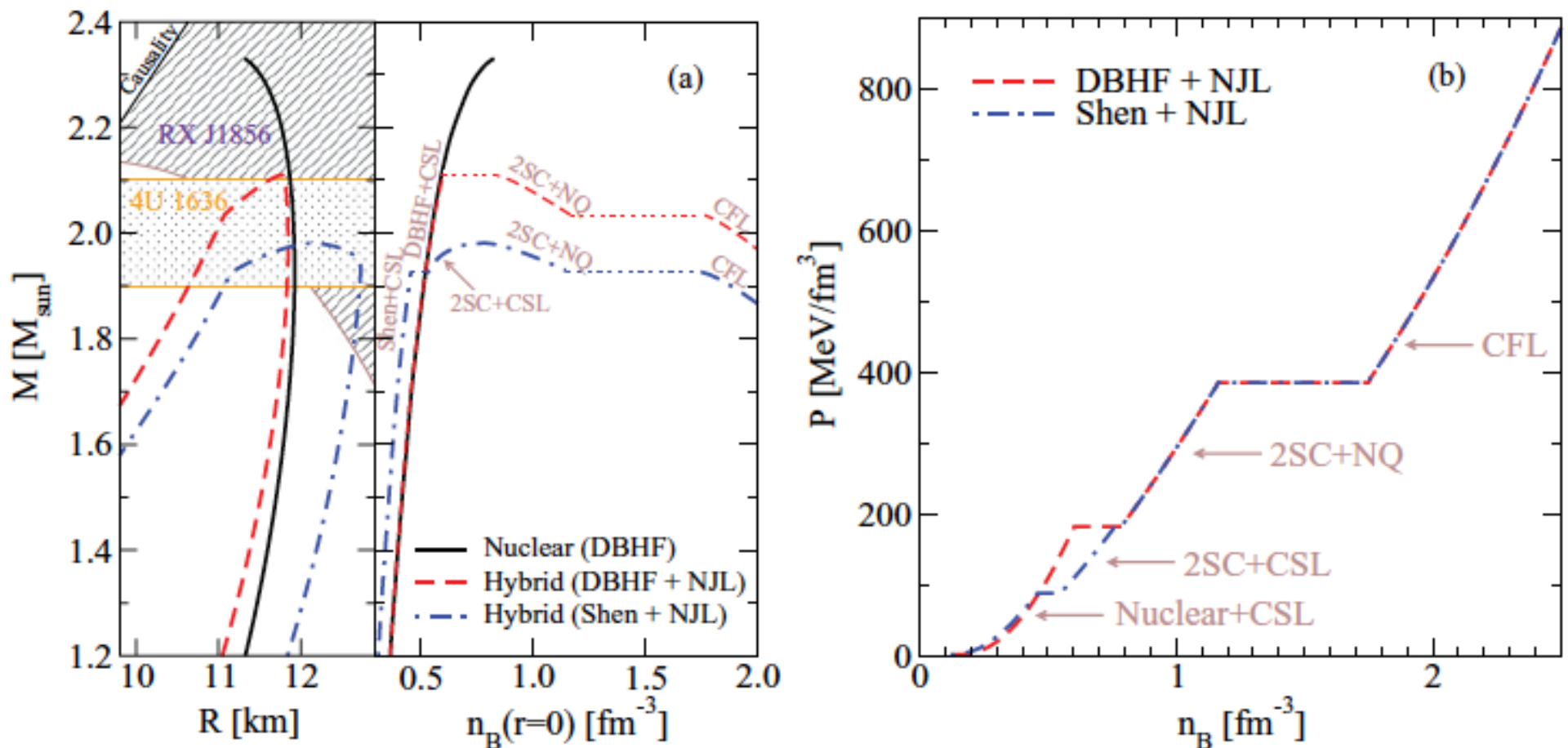
Increasing density



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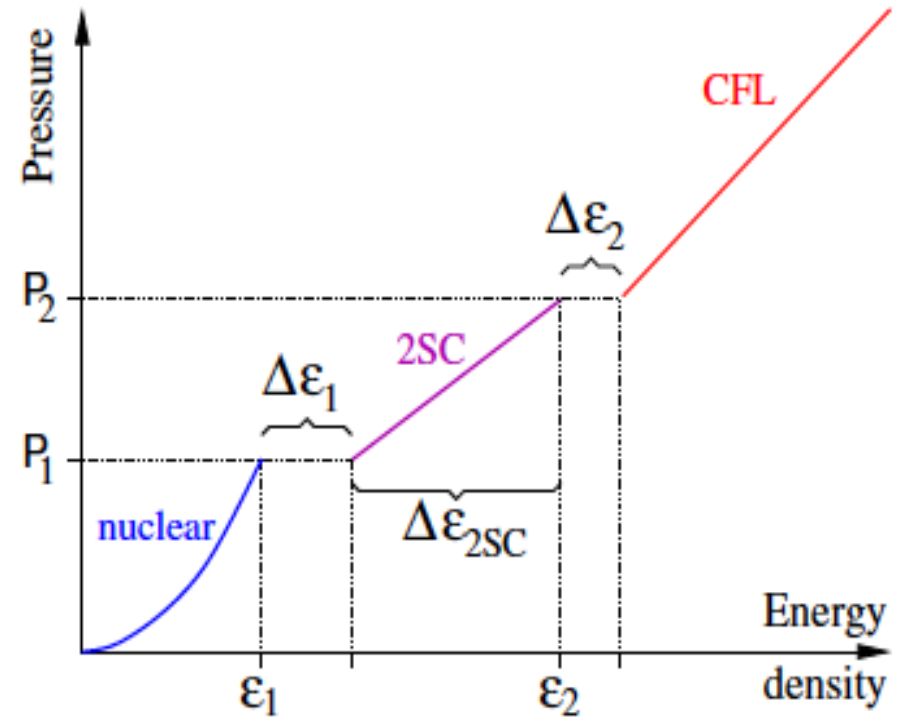
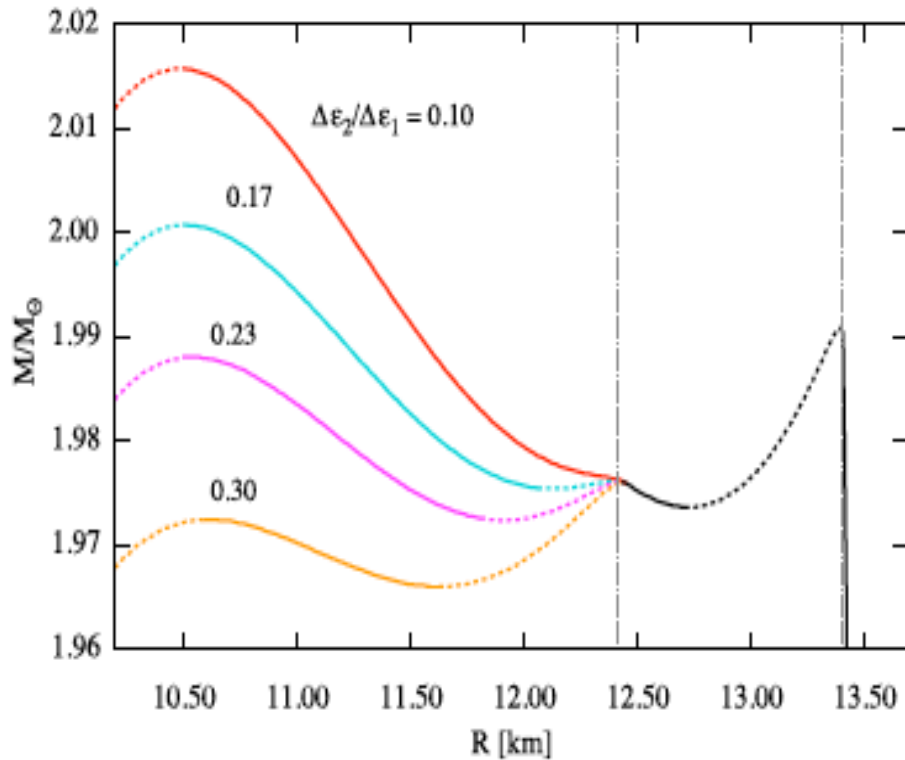


Neutron Star Interiors: Sequential Phase Transitions?

Measuring Mass vs. Radius



Equation of state



High-mass twins:

D. Blaschke et al., PoS CPOD 2013
S. Benic et al., A&A 577 (2015) A50

High-mass triples and fourth family:

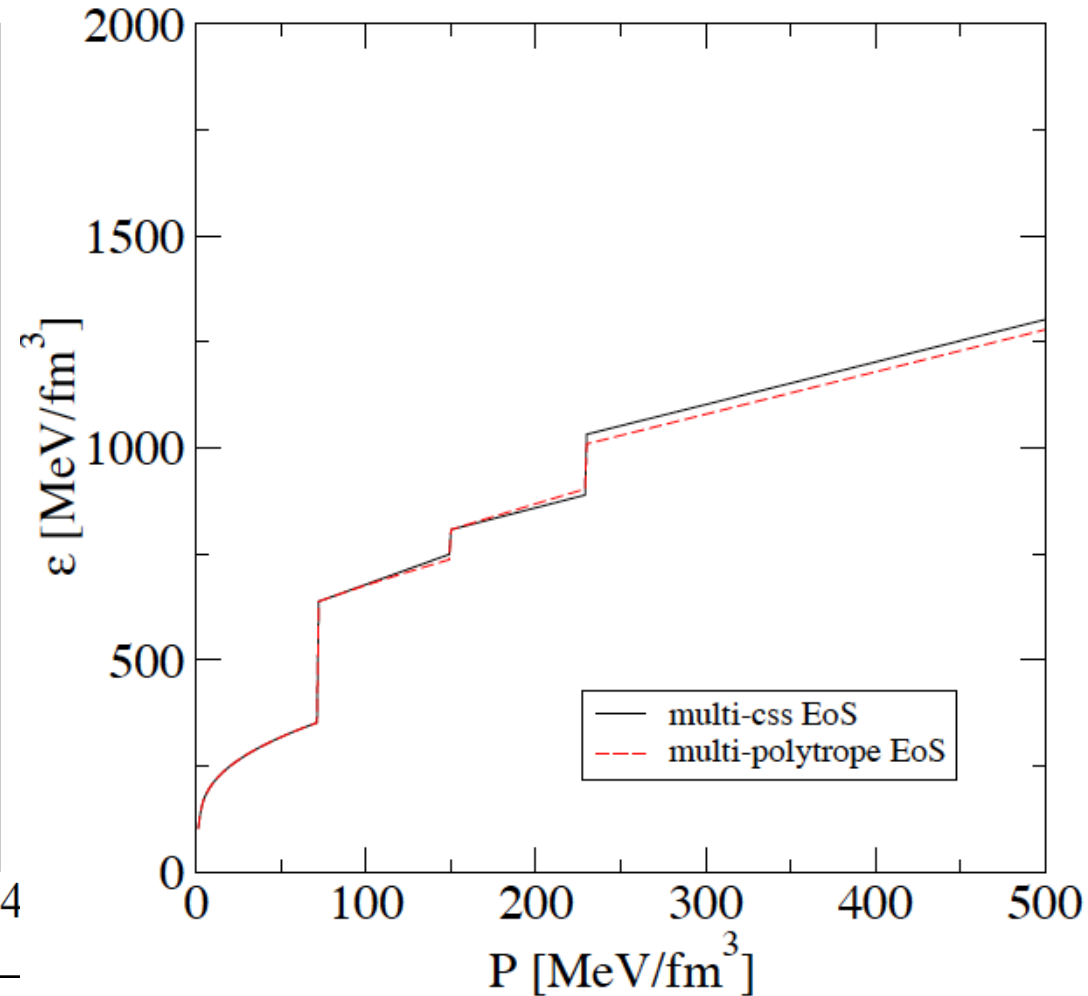
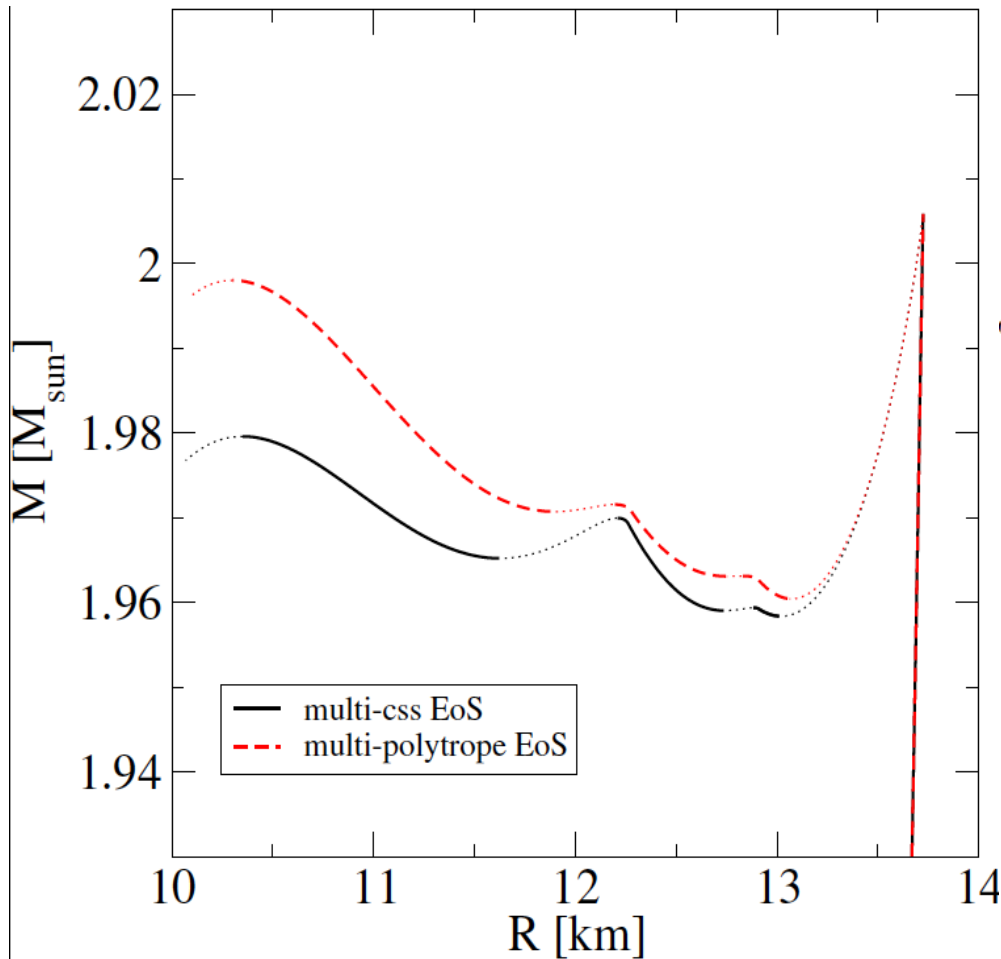
M. Alford and A. Sedrakian, arxiv:1706.01592
PRL 119 (2017)

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High-mass triples and fifth family:

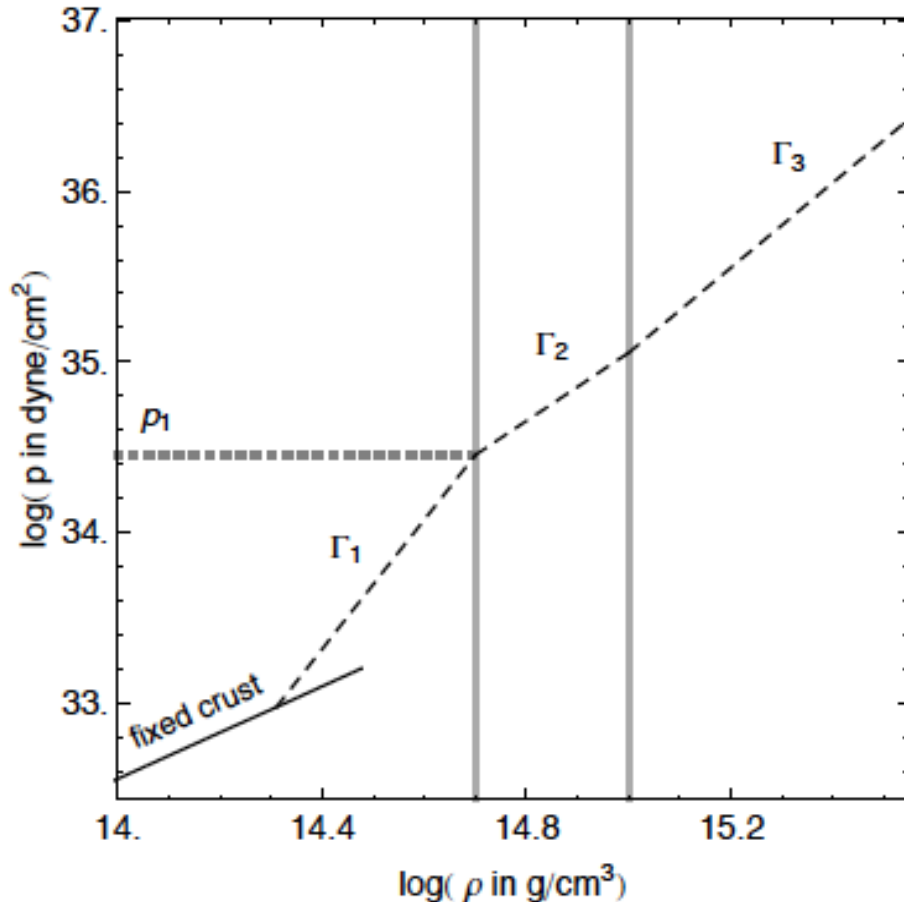
A. Ayriyan, D.B., H. Grigorian, in preparation (2017)

3. Piecewise polytrope EoS – high mass twins (HMT)?

J. Read et al., PRD 79, 124032 (2009)

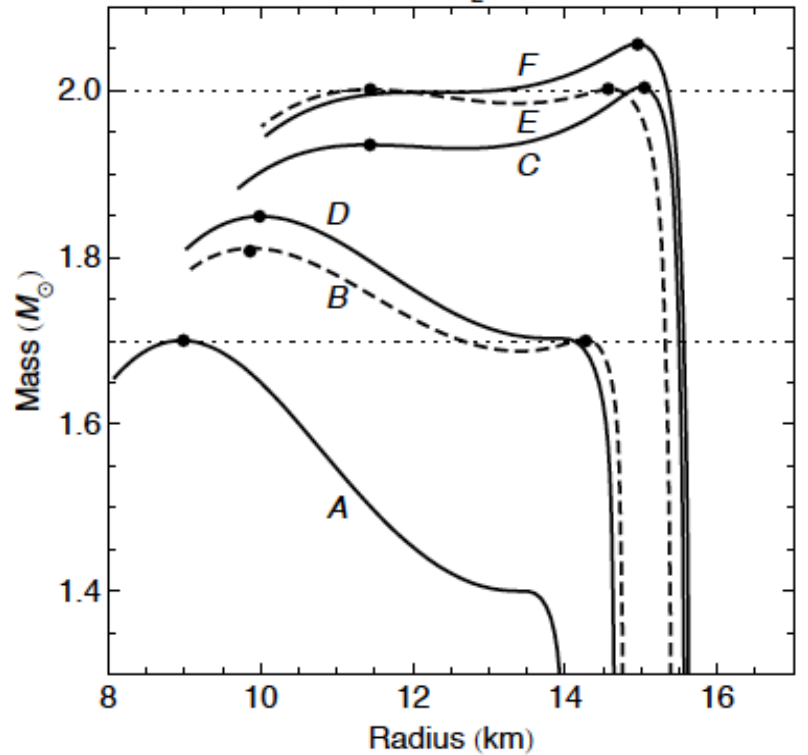
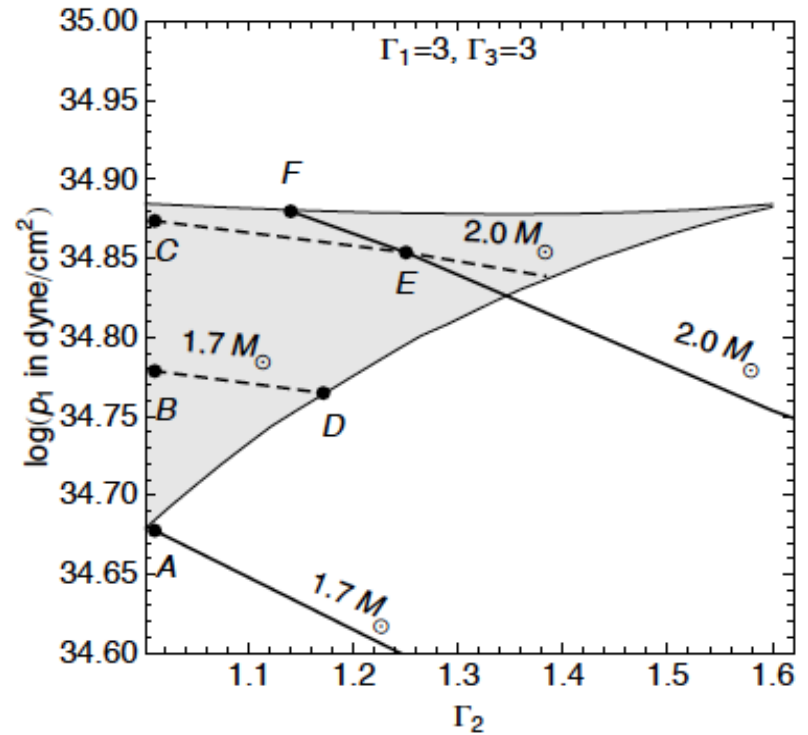
$$P_i(n) = \kappa_i n^{\Gamma_i}$$

- $i = 1 : n_1 \leq n \leq n_{12}$
- $i = 2 : n_{12} \leq n \leq n_{23}$
- $i = 3 : n \geq n_{23}$,



Case E:

**HMT @
2 M_{sun}**



3. Piecewise polytropic EoS – high mass twins?

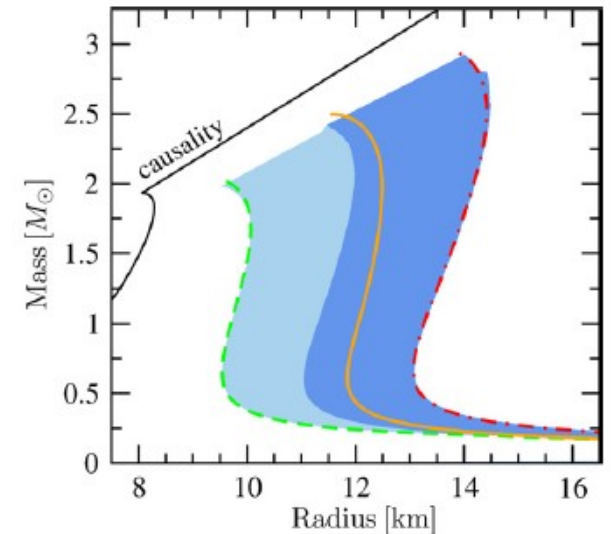
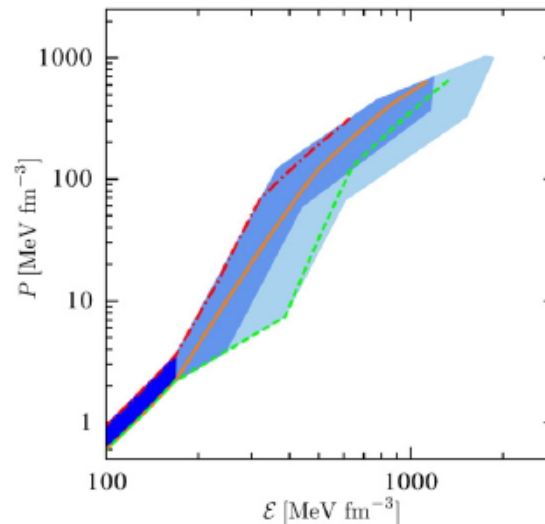
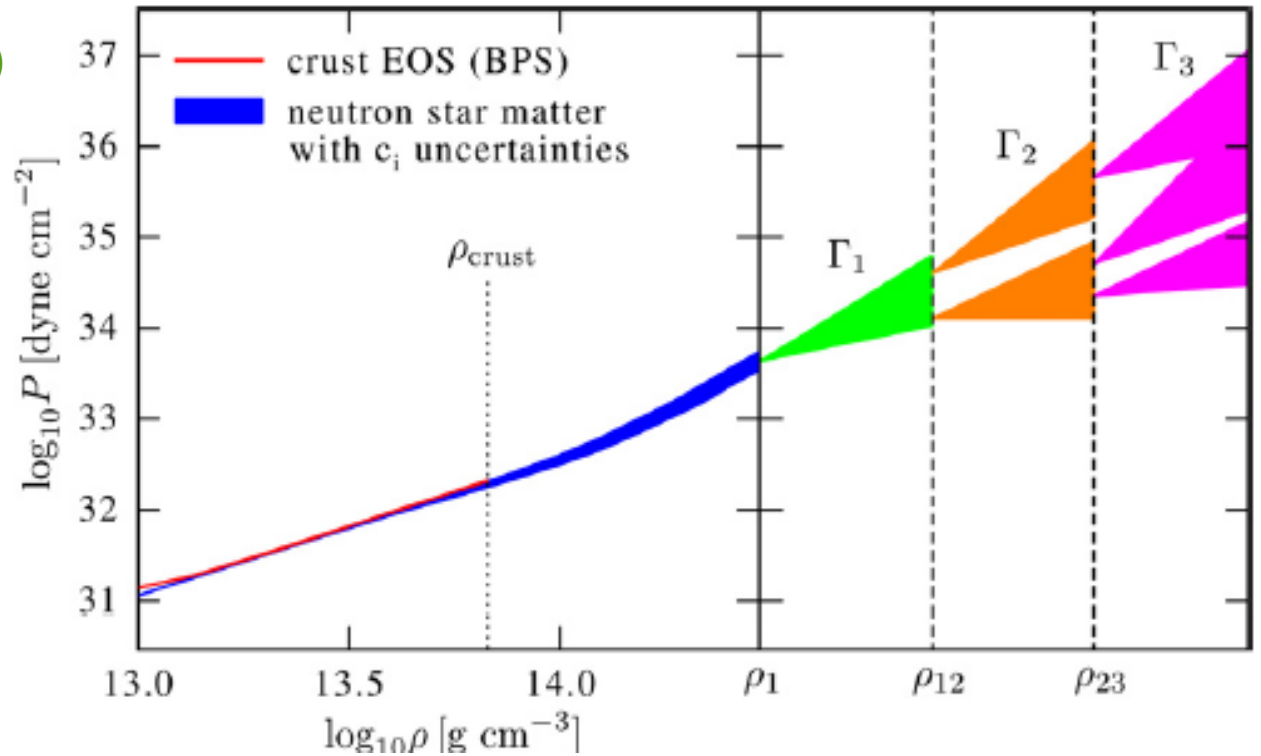
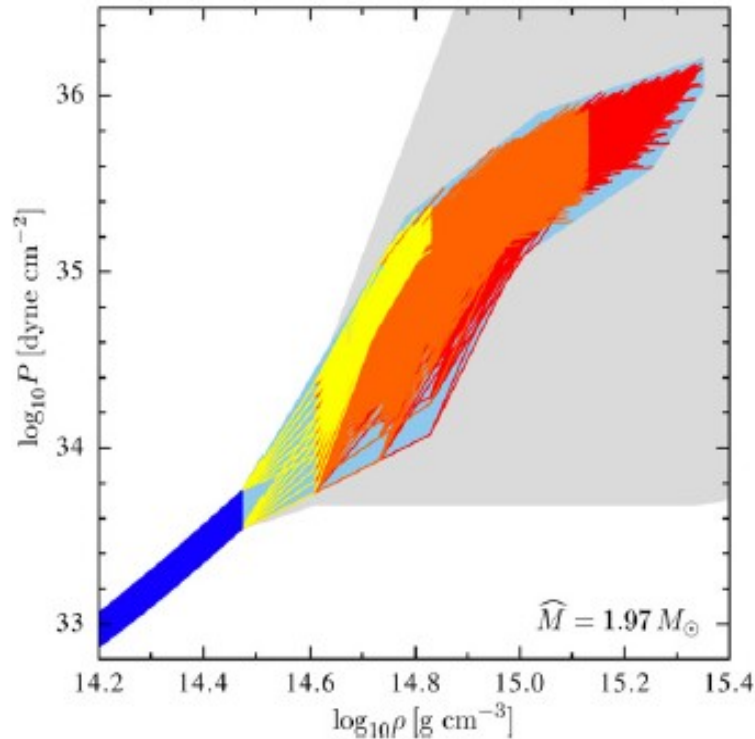
Hebeler et al., ApJ 773, 11 (2013)

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3. Piecewise polytrope EoS – high mass twins?

Hebeler et al., ApJ 773, 11 (2013)

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$$i = 1 : n_1 \leq n \leq n_{12}$$

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$$i = 3 : n \geq n_{23} ,$$

Here, 1st order PT in region 2:

$$\Gamma_2 = 0 \text{ and } P_2 = \kappa_2 = P_{\text{crit}}$$

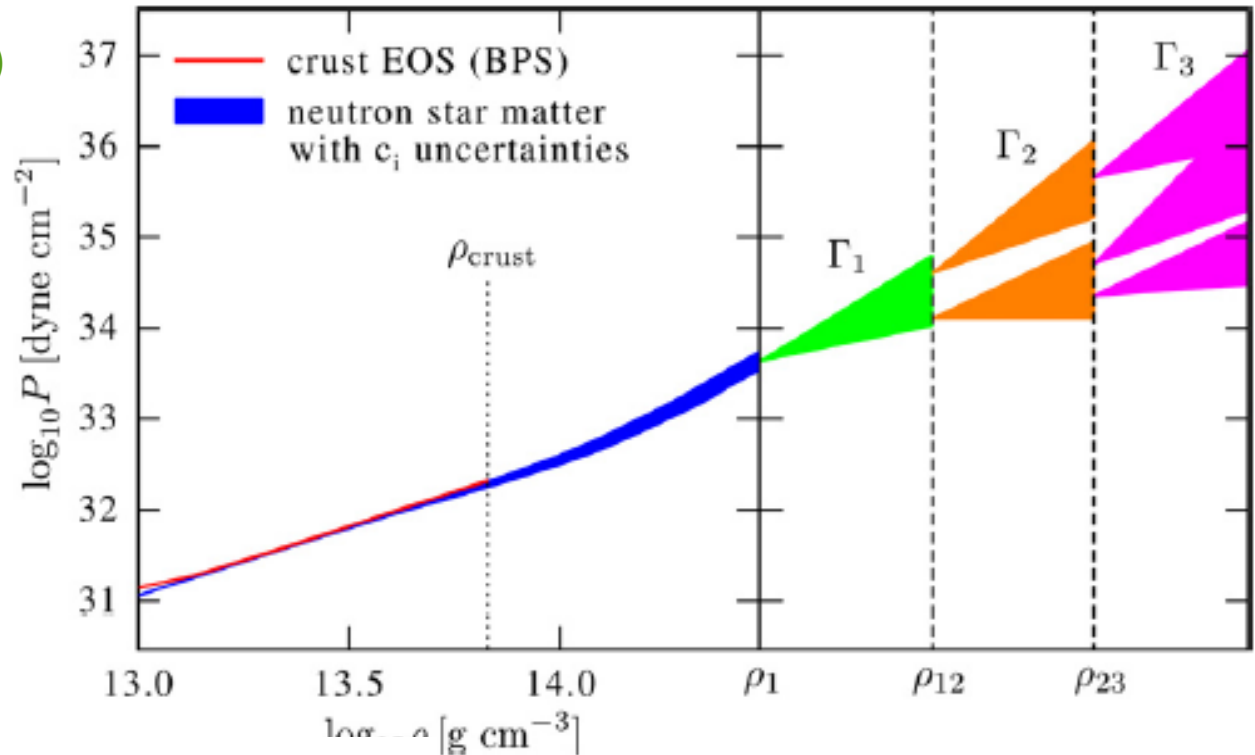
$$P(n) = n^2 \frac{d(\varepsilon(n)/n)}{dn},$$

$$\varepsilon(n)/n = \int dn \frac{P(n)}{n^2} = \int dn \kappa n^{\Gamma-2} = \frac{\kappa n^{\Gamma-1}}{\Gamma-1} + C,$$

$$\mu(n) = \frac{P(n) + \varepsilon(n)}{n} = \frac{\kappa \Gamma}{\Gamma-1} n^{\Gamma-1} + m_0,$$

Seidov criterion for instability:

$$\frac{\Delta\varepsilon}{\varepsilon_{\text{crit}}} \geq \frac{1}{2} + \frac{3}{3} \frac{P_{\text{crit}}}{\varepsilon_{\text{crit}}}$$



$$n(\mu) = \left[(\mu - m_0) \frac{\Gamma - 1}{\kappa \Gamma} \right]^{1/(\Gamma-1)}$$

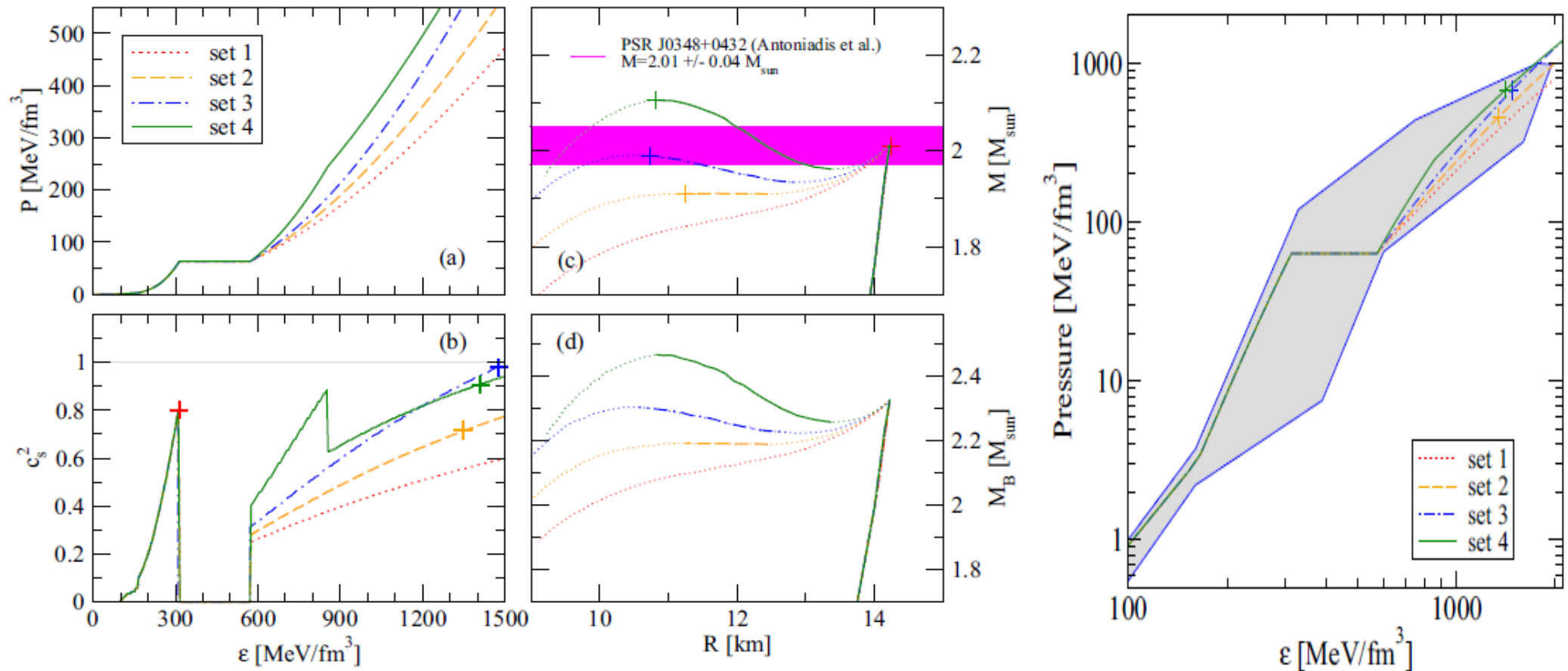
$$P(\mu) = \kappa \left[(\mu - m_0) \frac{\Gamma - 1}{\kappa \Gamma} \right]^{\Gamma/(\Gamma-1)}$$

Maxwell construction:

$$P_1(\mu_{\text{crit}}) = P_3(\mu_{\text{crit}}) = P_{\text{crit}}$$

$$\mu_{\text{crit}} = \mu_1(n_{12}) = \mu_3(n_{23})$$

3. Piecewise polytropic EoS – high mass twins?



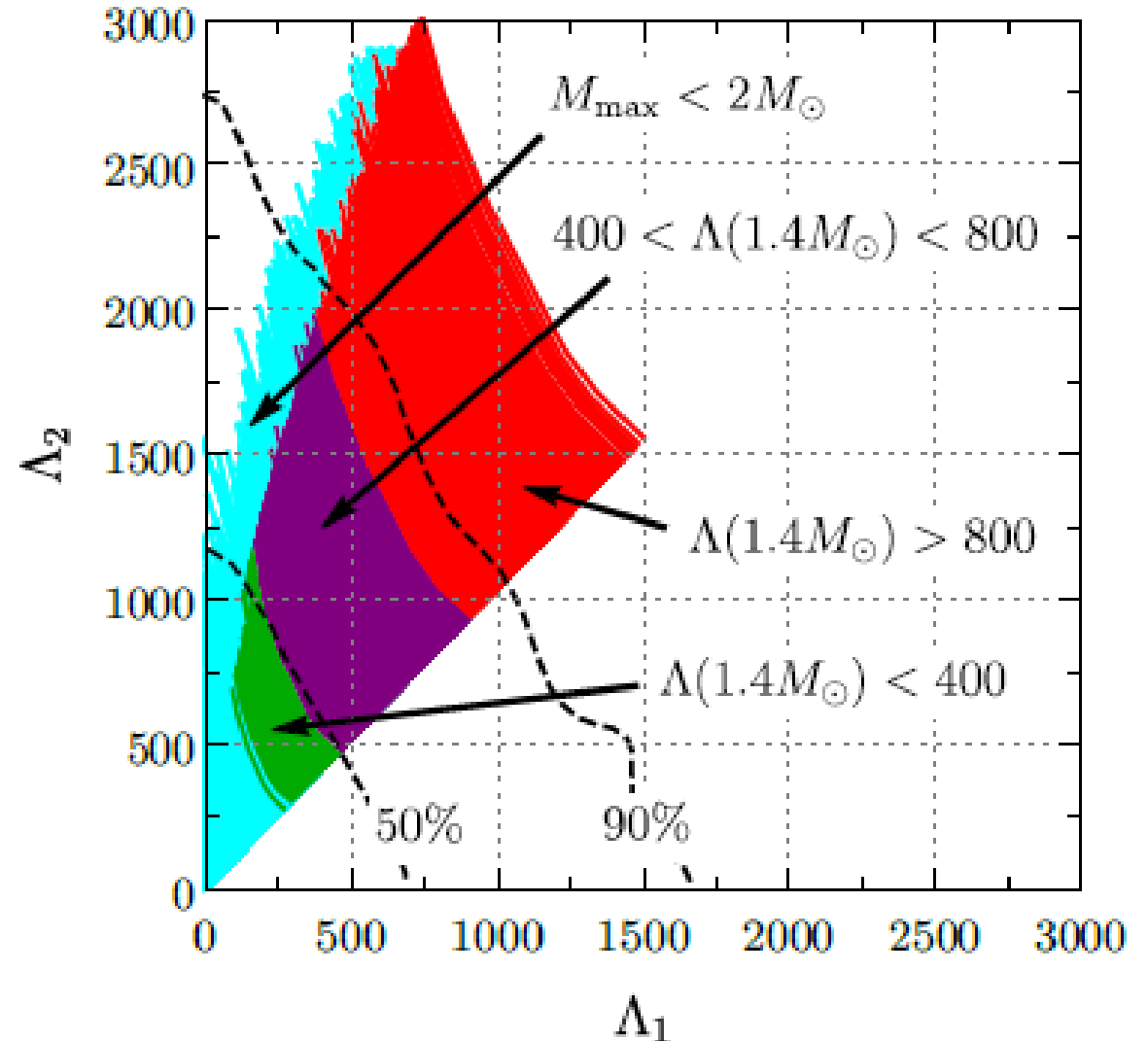
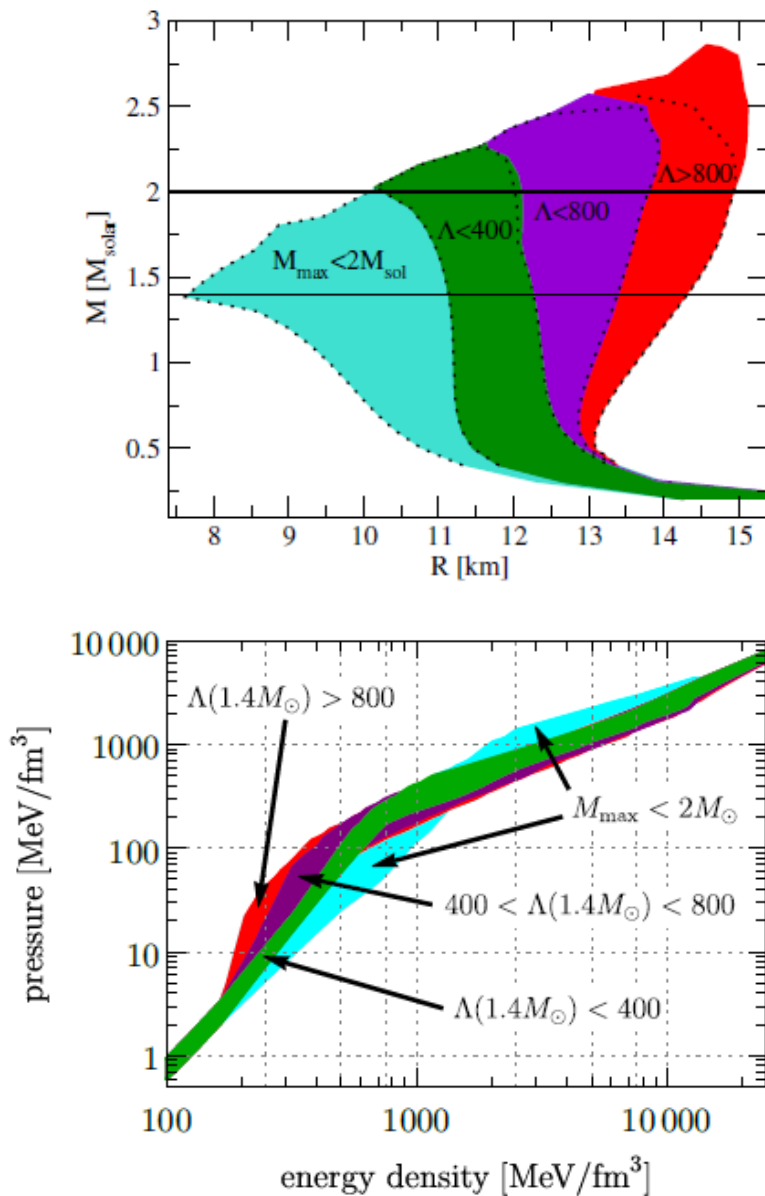
	Γ_3	κ_3 [MeV fm $^{3(\Gamma_3-1)}$]	$m_{0,3}$ [MeV]	M_{\max}^{NS} [M_{\odot}]	M_{\max}^{HS} [M_{\odot}]	M_{\min}^{HS} [M_{\odot}]
set 1	2.50	302.56	991.75	2.01	–	–
set 2	2.80	365.12	1004.88	2.01	1.910	1.909
set 3	3.12	447.16	1014.87	2.01	1.991	1.934
set 4a	4.00	774.375	1031.815			
set 4b	2.80	548.309	958.553	2.01	2.106	1.961

All sets with same onset of phase transition;
 $P_{\text{crit}} = 63.2 \text{ MeV/fm}^3$, $\epsilon_{\text{crit}} = 318.3 \text{ MeV/fm}^3$
 and same jump in energy density
 $\Delta\epsilon = 253.9 \text{ MeV/fm}^3$; varying Γ_3

Third family solutions found at 2 Msol (HMT),
 4-tropes favored; match with Hebeler et al.!
 [D. Alvarez & D.B. PRC 96 (2017) 045809]

Gravitational-wave constraints on the neutron-star-matter Equation of State

Eemeli Annala,¹ Tyler Gorda,¹ Aleksi Kurkela,² and Aleksi Vuorinen¹



Unfortunately, twins and third family forgotten !!!
 For this aim, 2- and 3-tropes not sufficient, 4-tropes!

Implications from GW170817 and I-Love-Q relations for relativistic hybrid stars

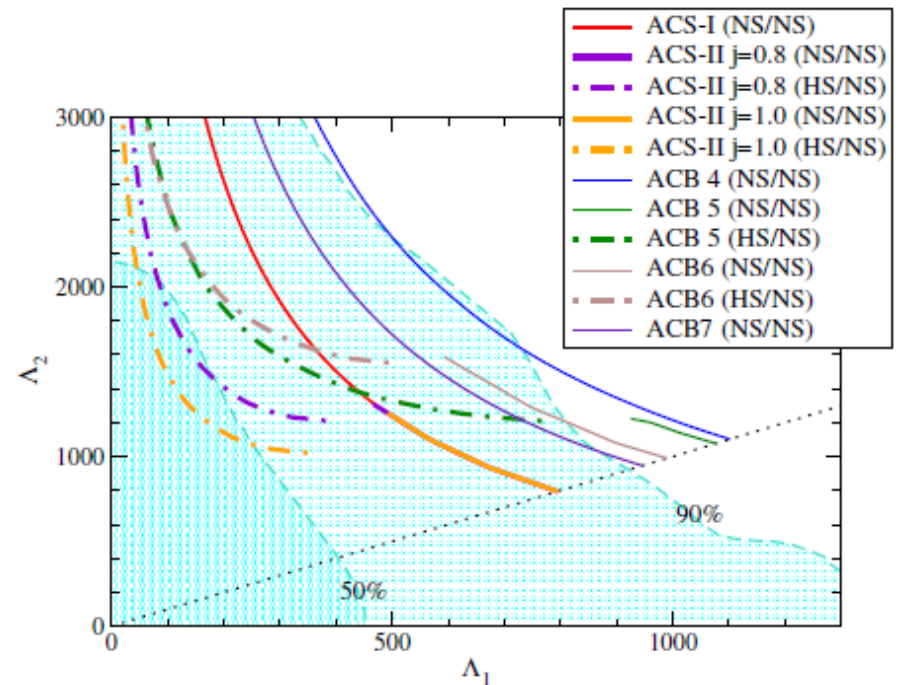
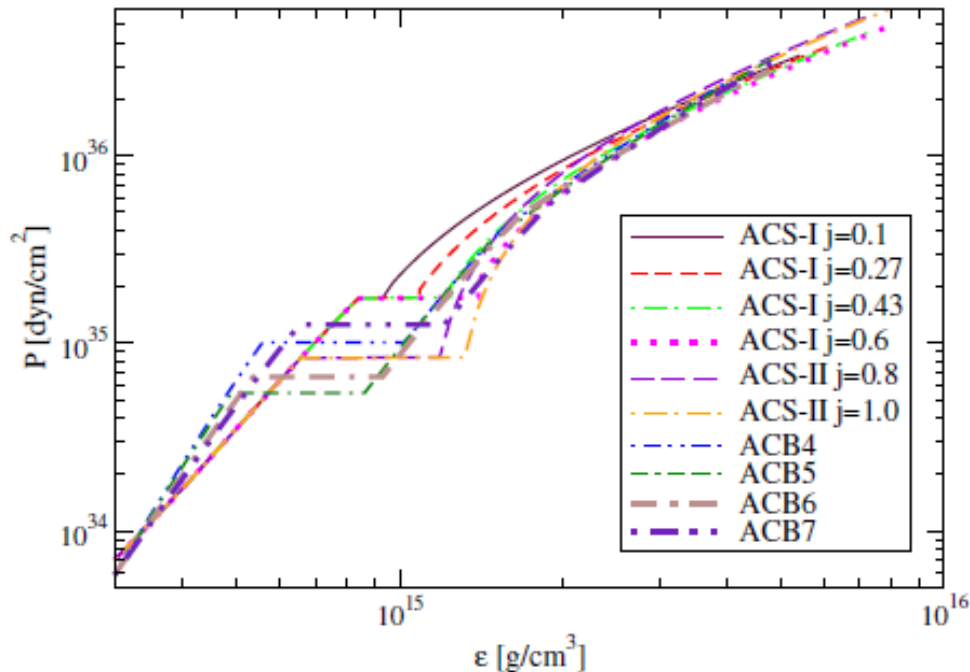
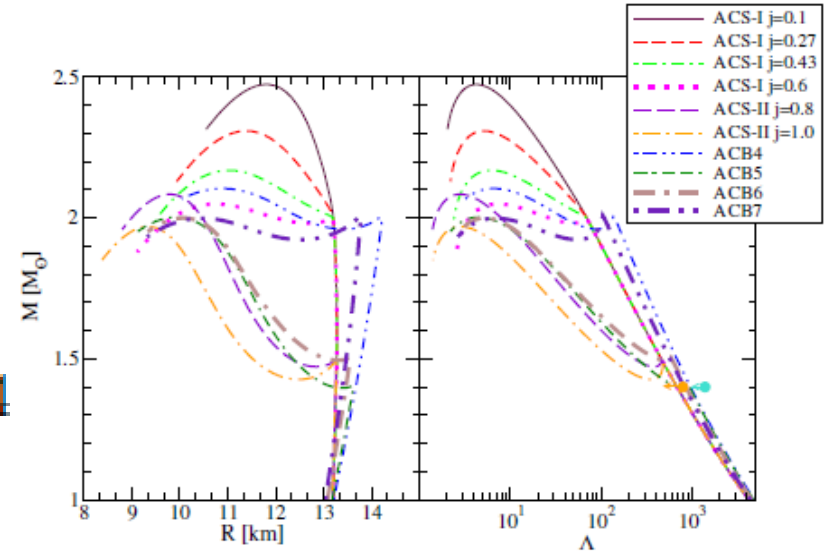
Vasileios Paschalidis,¹ Kent Yagi,² David Alvarez-Castillo,^{3,4} David B. Blaschke,^{3,5,6} and Armen Sedrakian⁷

CSS model (ACSX)

$$P(\varepsilon) = \begin{cases} P_{\text{tr}}, & \varepsilon_1 \leq \varepsilon \leq \varepsilon_2, \\ P_{\text{tr}} + c_s^2(\varepsilon - \varepsilon_2), & \varepsilon > \varepsilon_2, \end{cases}$$

MP model (ACBx)

$$P(n) = \kappa_i (n/n_0)^{\Gamma_i}, \quad n_i < n < n_{i+1}, \quad i = 1 \dots 4$$



Relativistic density functional approach to quark matter - string-flip model (SFM)



PHYSICAL REVIEW D

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1 DECEMBER 1986

Pauli quenching effects in a simple string model of quark/nuclear matter

G. Röpke and D. Blaschke

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H. Schulz

*Central Institute for Nuclear Research, Rossendorf, 8051 Dresden, German Democratic Republic
and The Niels Bohr Institute, 2100 Copenhagen, Denmark*

(Received 16 December 1985)

Relativistic density functional approach* (I)

$$\mathcal{Z} = \int \mathcal{D}\bar{q}\mathcal{D}q \exp \left\{ \int_0^\beta d\tau \int_V d^3x [\mathcal{L}_{\text{eff}} + \bar{q}\gamma_0\hat{\mu}q] \right\}, \quad q = \begin{pmatrix} q_u \\ q_d \end{pmatrix}, \quad \hat{\mu} = \text{diag}(\mu_u, \mu_d)$$

$$\mathcal{L}_{\text{eff}} = \mathcal{L}_{\text{free}} - U(\bar{q}q, \bar{q}\gamma_0q), \quad \mathcal{L}_{\text{free}} = \bar{q} \left(-\gamma_0 \frac{\partial}{\partial \tau} + i\vec{\gamma} \cdot \vec{\nabla} - \hat{m} \right) q, \quad \hat{m} = \text{diag}(m_u, m_d)$$

General nonlinear functional of quark density bilinears: scalar, vector, isovector, diquark ...
Expansion around the expectation values:

$$U(\bar{q}q, \bar{q}\gamma_0q) = U(n_s, n_v) + (\bar{q}q - n_s)\Sigma_s + (\bar{q}\gamma_0q - n_v)\Sigma_v + \dots,$$

$$\langle \bar{q}q \rangle = n_s = \sum_{f=u,d} n_{s,f} = - \sum_{f=u,d} \frac{T}{V} \frac{\partial}{\partial m_f} \ln \mathcal{Z}, \quad \Sigma_s = \left. \frac{\partial U(\bar{q}q, \bar{q}\gamma_0q)}{\partial (\bar{q}q)} \right|_{\bar{q}q=n_s} = \frac{\partial U(n_s, n_v)}{\partial n_s},$$

$$\langle \bar{q}\gamma_0q \rangle = n_v = \sum_{f=u,d} n_{v,f} = \sum_{f=u,d} \frac{T}{V} \frac{\partial}{\partial \mu_f} \ln \mathcal{Z}, \quad \Sigma_v = \left. \frac{\partial U(\bar{q}q, \bar{q}\gamma_0q)}{\partial (\bar{q}\gamma_0q)} \right|_{\bar{q}\gamma_0q=n_v} = \frac{\partial U(n_s, n_v)}{\partial n_v}$$

$$\mathcal{Z} = \int \mathcal{D}\bar{q}\mathcal{D}q \exp \{ \mathcal{S}_{\text{quasi}}[\bar{q}, q] - \beta V \Theta[n_s, n_v] \}, \quad \Theta[n_s, n_v] = U(n_s, n_v) - \Sigma_s n_s - \Sigma_v n_v$$

$$\mathcal{S}_{\text{quasi}}[\bar{q}, q] = \beta \sum_n \sum_{\vec{p}} \bar{q} G^{-1}(\omega_n, \vec{p}) q, \quad G^{-1}(\omega_n, \vec{p}) = \gamma_0(-i\omega_n + \hat{\mu}^*) - \vec{\gamma} \cdot \vec{p} - \hat{m}^*$$

*This work was inspired by the textbook on “Thermodynamics and statistical mechanics” of the “red” series on Theoretical Physics by Walter Greiner and Coworkers.

Relativistic density functional approach (II)

$$\mathcal{Z} = \int \mathcal{D}\bar{q}\mathcal{D}q \exp \{ \mathcal{S}_{\text{quasi}}[\bar{q}, q] - \beta V \Theta[n_s, n_v] \}, \quad \Theta[n_s, n_v] = U(n_s, n_v) - \Sigma_s n_s - \Sigma_v n_v$$

$$\mathcal{Z}_{\text{quasi}} = \int \mathcal{D}\bar{q}\mathcal{D}q \exp \{ \mathcal{S}_{\text{quasi}}[\bar{q}, q] \} = \det[\beta G^{-1}], \quad \ln \det A = \text{Tr} \ln A$$

$$P_{\text{quasi}} = \frac{T}{V} \ln \mathcal{Z}_{\text{quasi}} = \frac{T}{V} \text{Tr} \ln[\beta G^{-1}] \quad \text{“no sea” approximation ...}$$

$$= 2N_c \sum_{f=u,d} \int \frac{d^3p}{(2\pi)^3} \left\{ T \ln \left[1 + e^{-\beta(E_f^* - \mu_f^*)} \right] + T \ln \left[1 + e^{-\beta(E_f^* + \mu_f^*)} \right] \right\}$$

$$P_{\text{quasi}} = \sum_{f=u,d} \int \frac{dp}{\pi^2} \frac{p^4}{E_f^*} [f(E_f^* - \mu_f^*) + f(E_f^* + \mu_f^*)] \quad E_f^* = \sqrt{p^2 + m_f^{*2}}$$

$$f(E) = 1/[1 + \exp(\beta E)]$$

$$P = \sum_{f=u,d} \int_0^{p_{F,f}} \frac{dp}{\pi^2} \frac{p^4}{E_f^*} - \Theta[n_s, n_v], \quad p_{F,f} = \sqrt{\mu_f^{*2} - m_f^{*2}}$$

$$\hat{m}^* = \hat{m} + \Sigma_s$$

$$\hat{\mu}^* = \hat{\mu} - \Sigma_v$$

Selfconsistent densities

$$n_s = - \sum_{f=u,d} \frac{\partial P}{\partial m_f} = \frac{3}{\pi^2} \sum_{f=u,d} \int_0^{p_{F,f}} dp p^2 \frac{m_f^*}{E_f^*}, \quad n_v = \sum_{f=u,d} \frac{\partial P}{\partial \mu_f} = \frac{3}{\pi^2} \sum_{f=u,d} \int_0^{p_{F,f}} dp p^2 = \frac{p_{F,u}^3 + p_{F,d}^3}{\pi^2}.$$

Relativistic density functional approach (III)

Density functional for the SFM

$$U(n_s, n_v) = D(n_v)n_s^{2/3} + an_v^2 + \frac{bn_v^4}{1 + cn_v^2},$$

Quark selfenergies

$$\Sigma_s = \frac{2}{3}D(n_v)n_s^{-1/3}, \quad \text{Quark "confinement"}$$

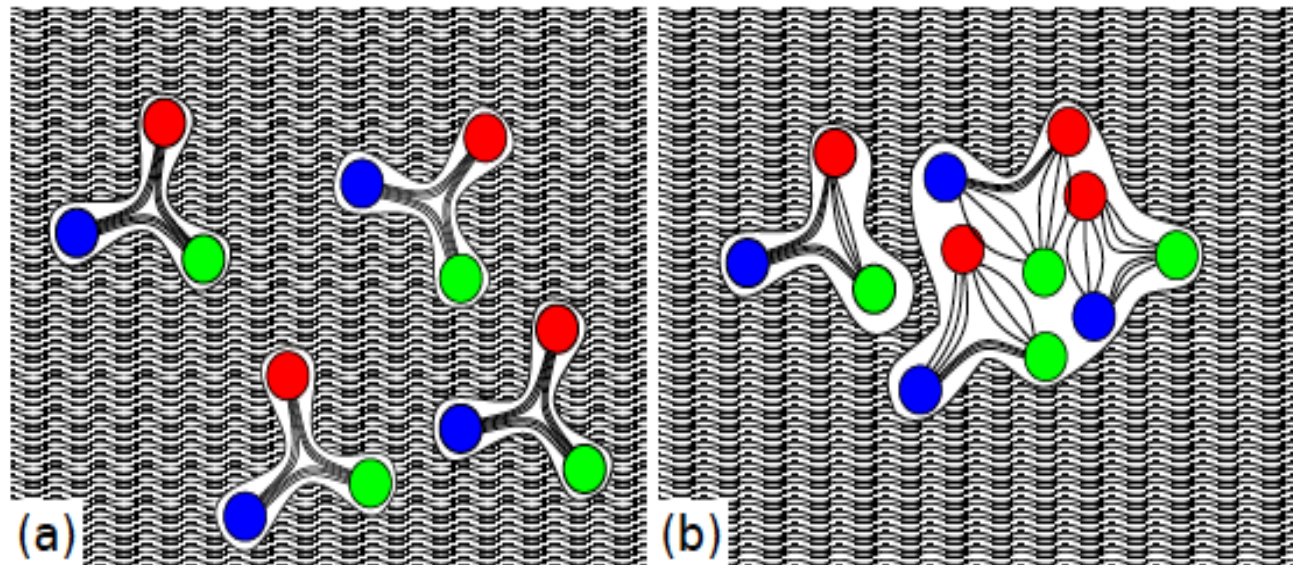
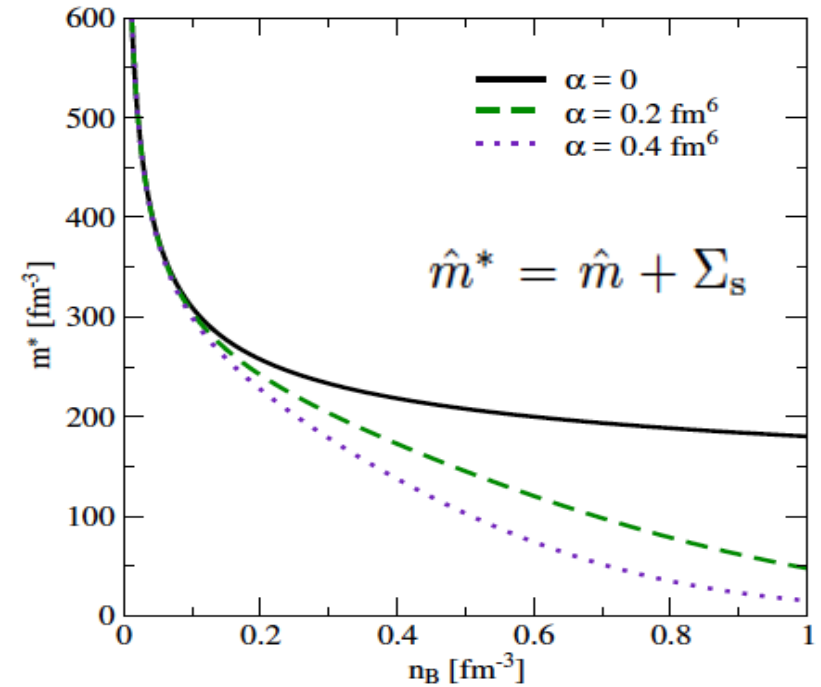
$$\Sigma_v = 2an_v + \frac{4bn_v^3}{1 + cn_v^2} - \frac{2bcn_v^5}{(1 + cn_v^2)^2} + \frac{\partial D(n_v)}{\partial n_v}n_s^{2/3}$$

String tension & confinement due to dual Meissner effect (dual superconductor model)

$$D(n_v) = D_0\Phi(n_v)$$

Effective screening of the string tension in dense matter by a reduction of the available volume $\alpha = v|v|/2$

$$\Phi(n_B) = \begin{cases} 1, & \text{if } n_B < n_0 \\ e^{-\alpha(n_B - n_0)^2}, & \text{if } n_B > n_0 \end{cases}$$



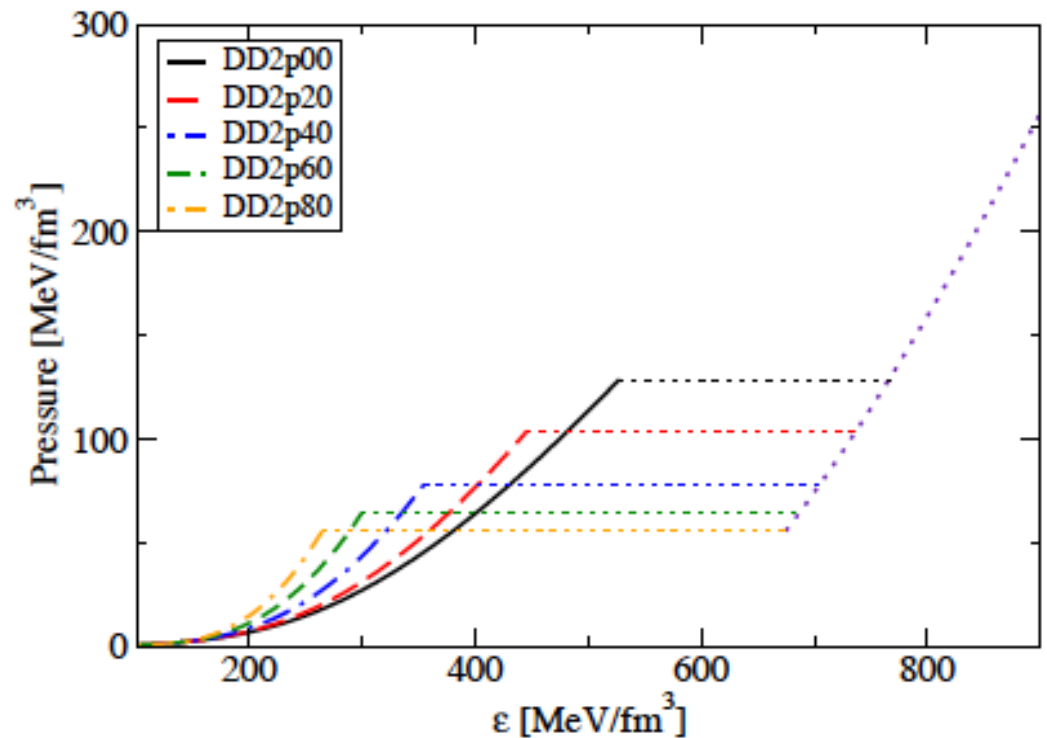
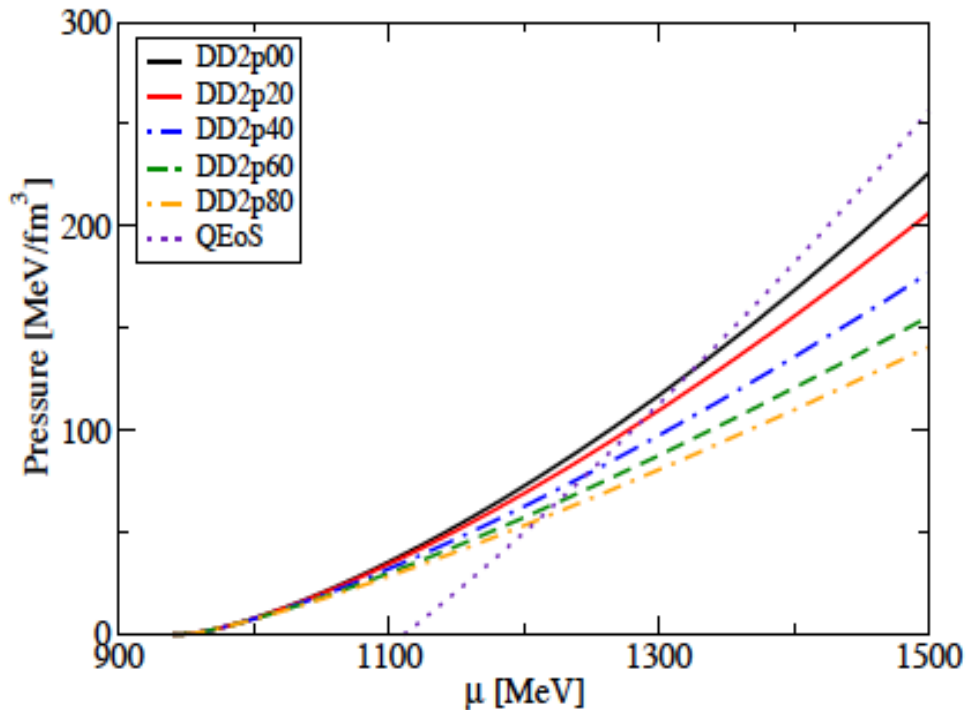
3. Phase transition to SFM quark matter

Hadronic matter: DD2 with excluded volume

[S. Typel, EPJA 52 (3) (2016)]

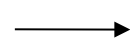
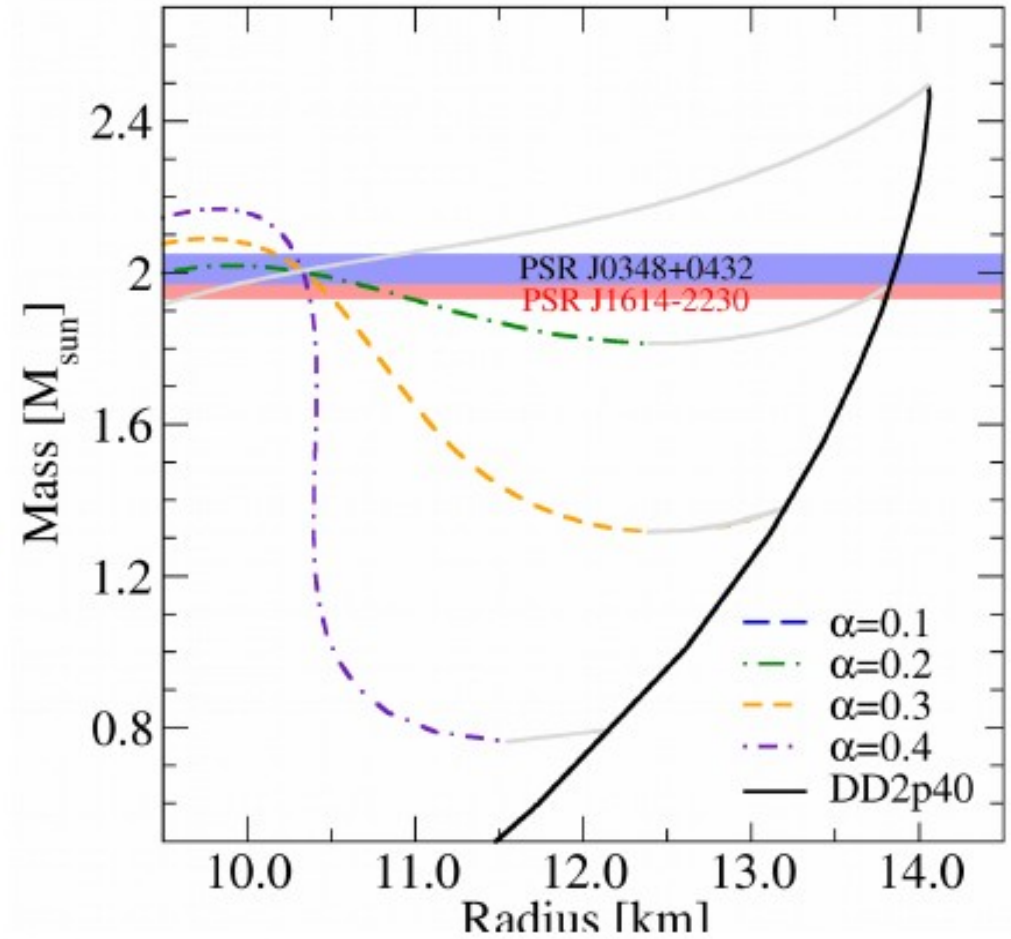
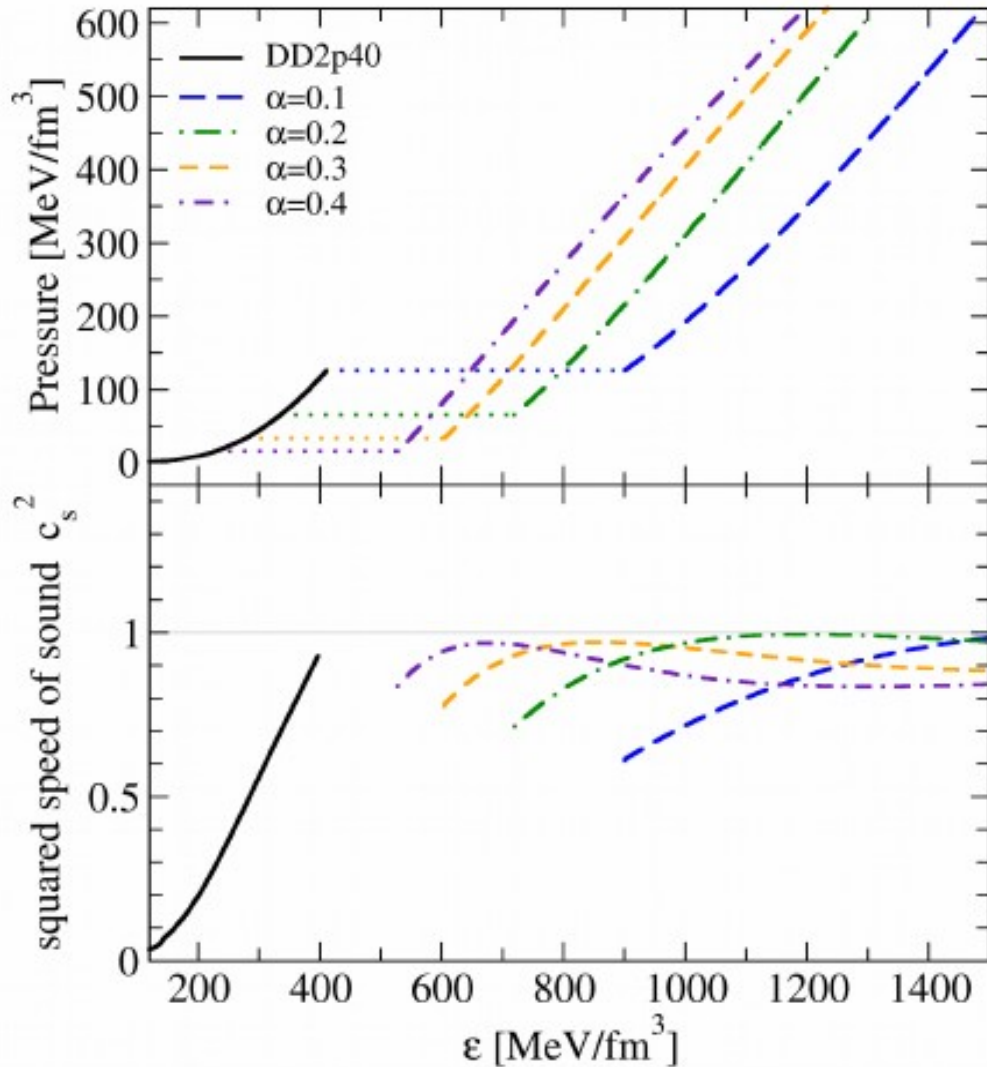
$$\Phi_n = \Phi_p = \begin{cases} 1, & \text{if } n_B < n_0 \\ e^{-\frac{v|v|}{2}(n_B - n_0)^2}, & \text{if } n_B > n_0 \end{cases}$$

Varying the hadronic excluded volume parameter, p00 \rightarrow v=0, ... , p80 \rightarrow v=8 fm³



Hybrid EOS - parameters

α, a, b



KVOR_cut2 vs. string-flip model (SFM)

Robustness of Twins against Pasta Phase Effects

arxiv:1711.03926v1, 10.11.2017

How robust is a third family of compact stars against pasta phase effects?

A. Ayriyan,^{1,*} N.-U. Bastian,^{2,†} D. Blaschke,^{2,3,4,‡} H. Grigorian,^{1,§} K. Maslov,^{3,4,¶} and D. N. Voskresensky^{3,4,**}

¹Laboratory for Information Technologies, Joint Institute for Nuclear Research, Joliot-Curie street 6, 141980 Dubna, Russia

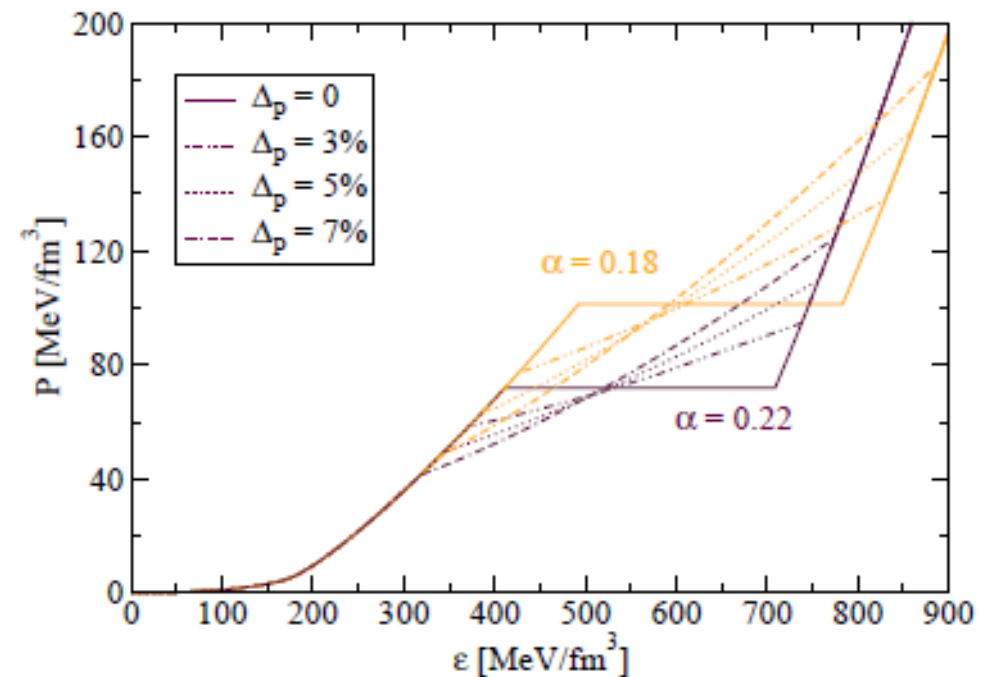
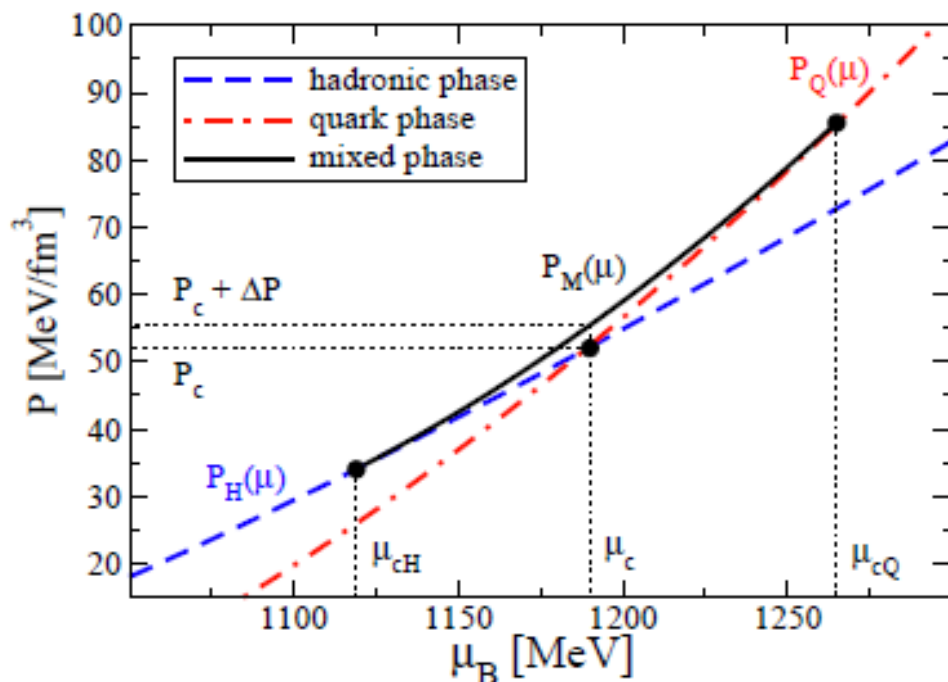
²Institute of Theoretical Physics, University of Wrocław, Max Born place 9, 50-204 Wrocław, Poland

³Bogoliubov Laboratory for Theoretical Physics, Joint Institute

for Nuclear Research, Joliot-Curie street 6, 141980 Dubna, Russia

⁴National Research Nuclear University (MEPhI), Kashirskoe Shosse 31, 115409 Moscow, Russia

(Dated: November 13, 2017)



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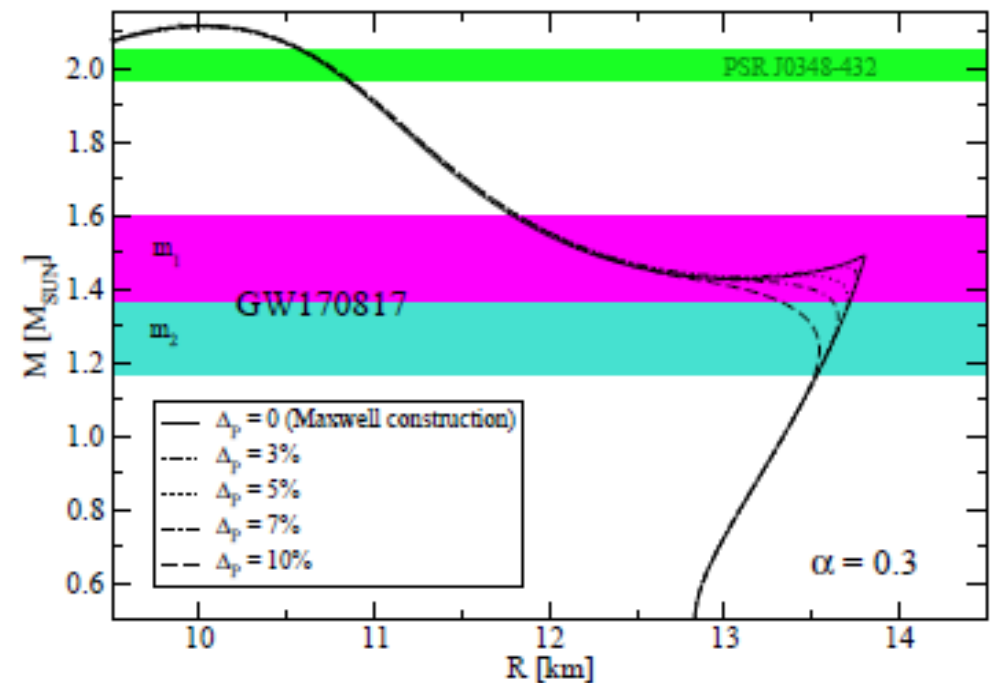
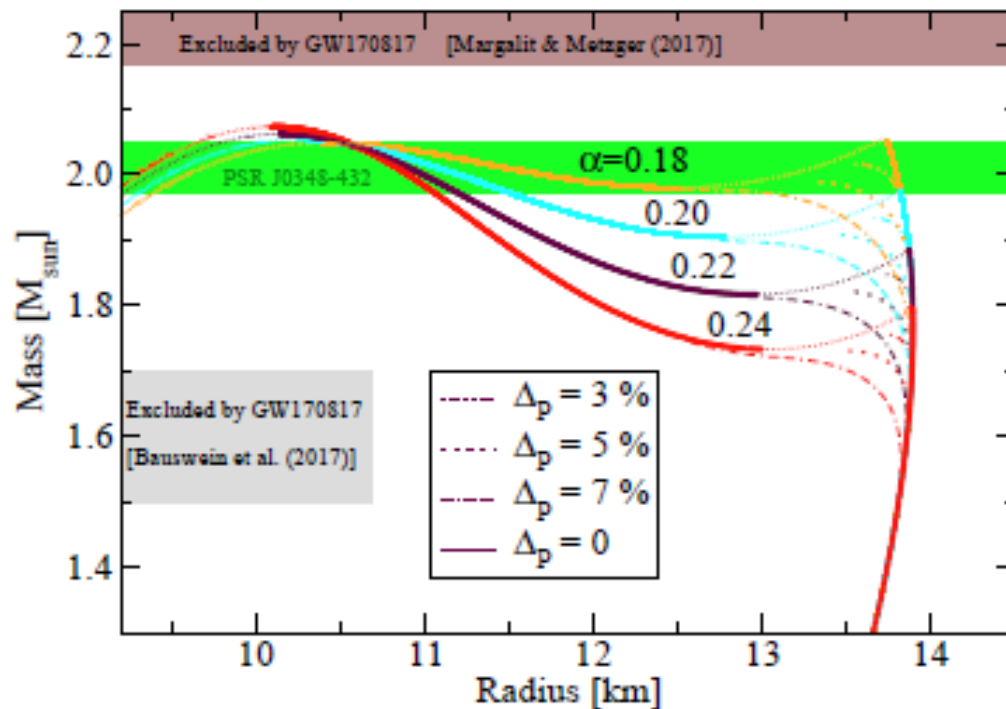
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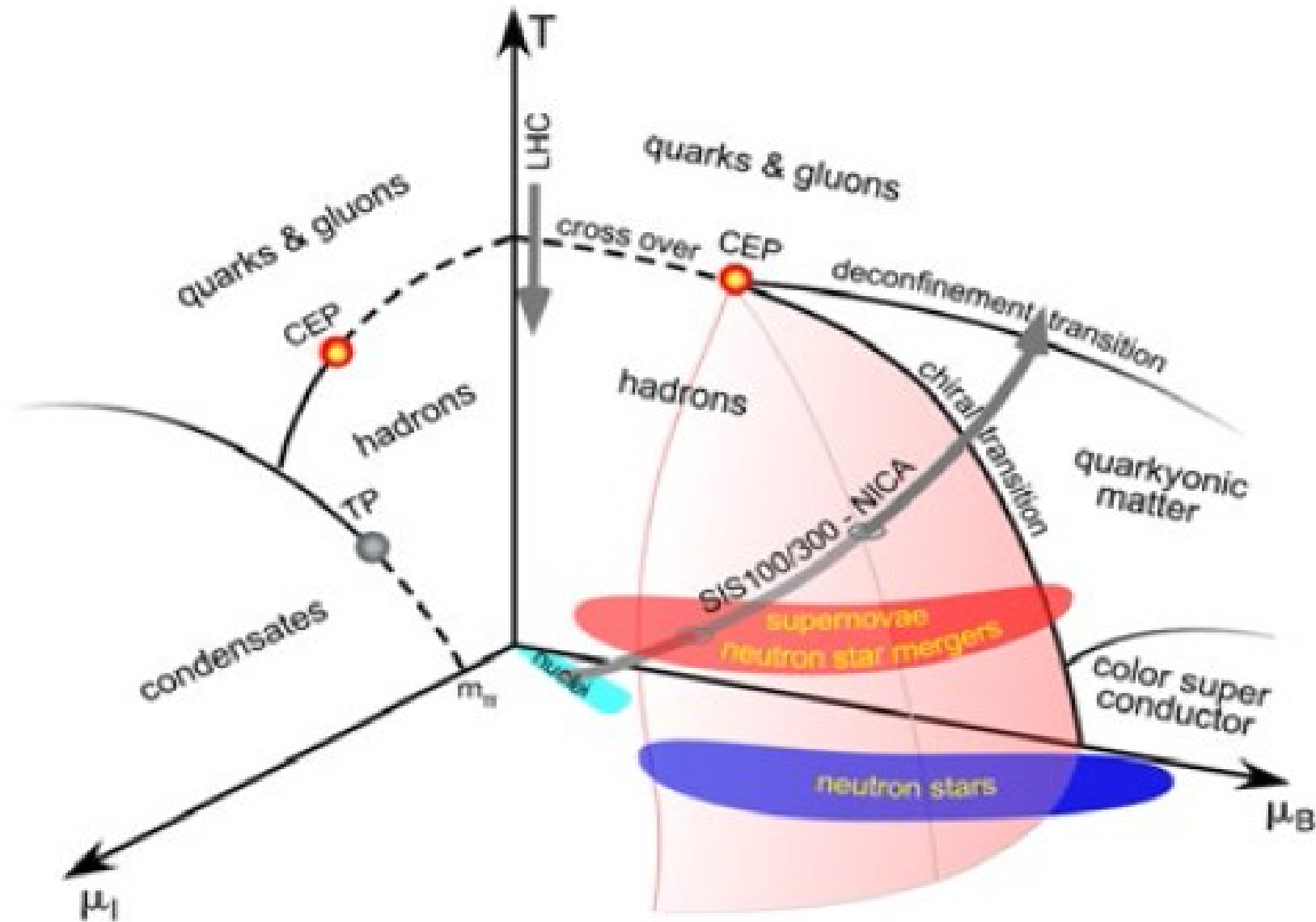
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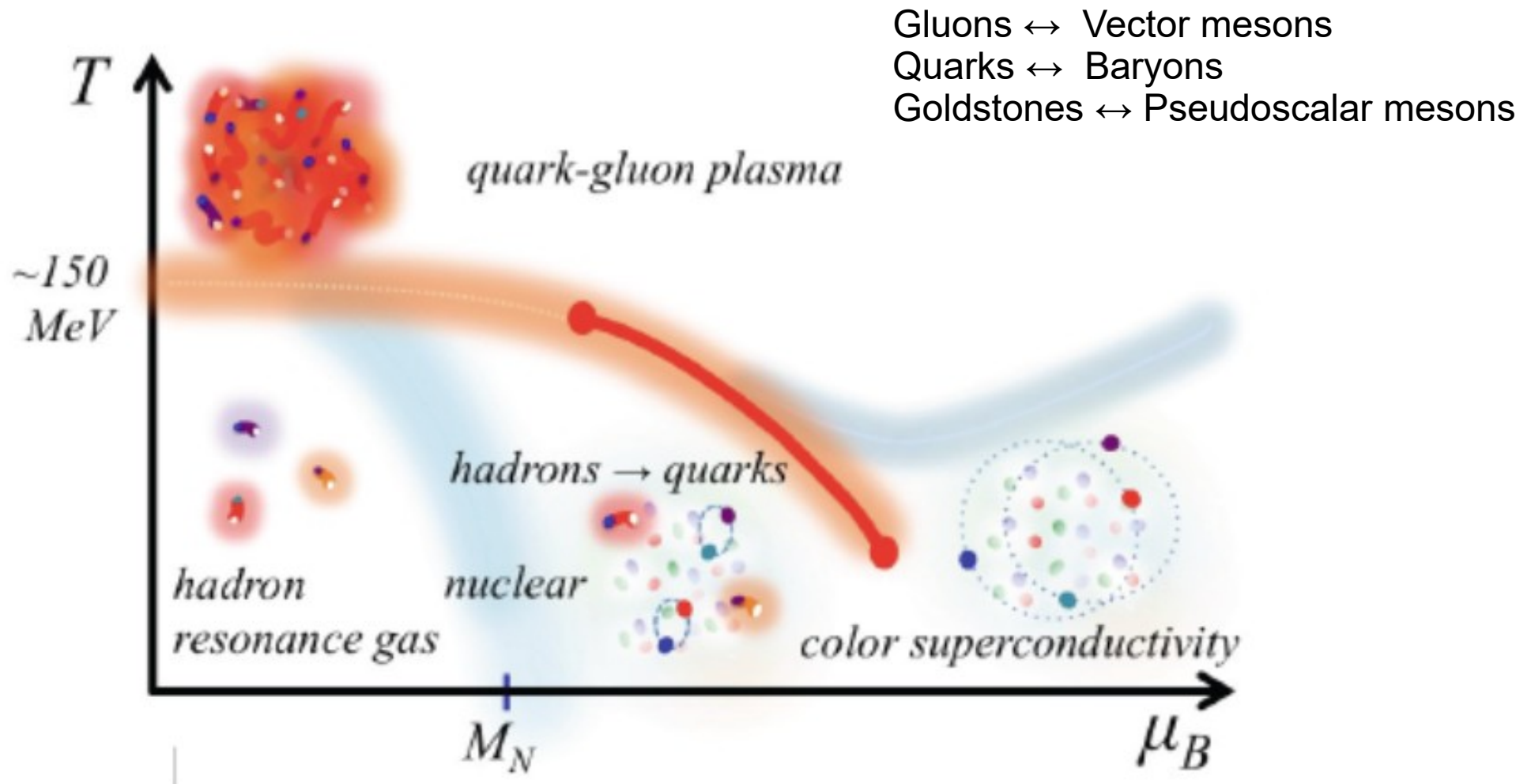
(Dated: November 13, 2017)



CEP in the QCD phase diagram: HIC vs. Astrophysics



2nd CEP in QCD phase diagram: Quark-Hadron Continuity?



T. Schaefer & F. Wilczek, Phys. Rev. Lett. 82 (1999) 3956

C. Wetterich, Phys. Lett. B 462 (1999) 164

T. Hatsuda, M. Tachibana, T. Yamamoto & G. Baym, Phys. Rev. Lett. 97 (2006) 122001

Conclusions:

High-mass twin (HMT) and Typical-mass twin (TMT) solutions obtained within different hybrid star EoS, e.g.,

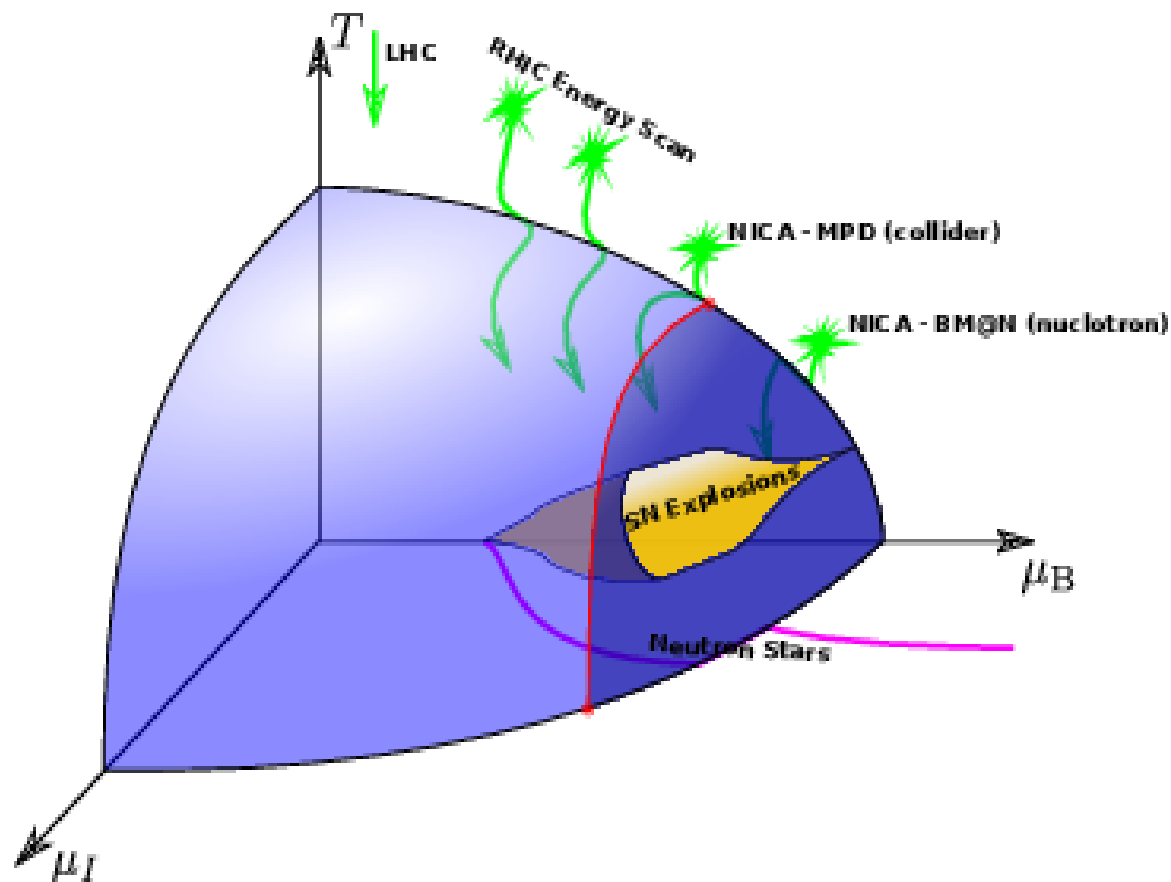
- constant speed of sound
- higher order NJL
- piecewise polytrope
- density functional

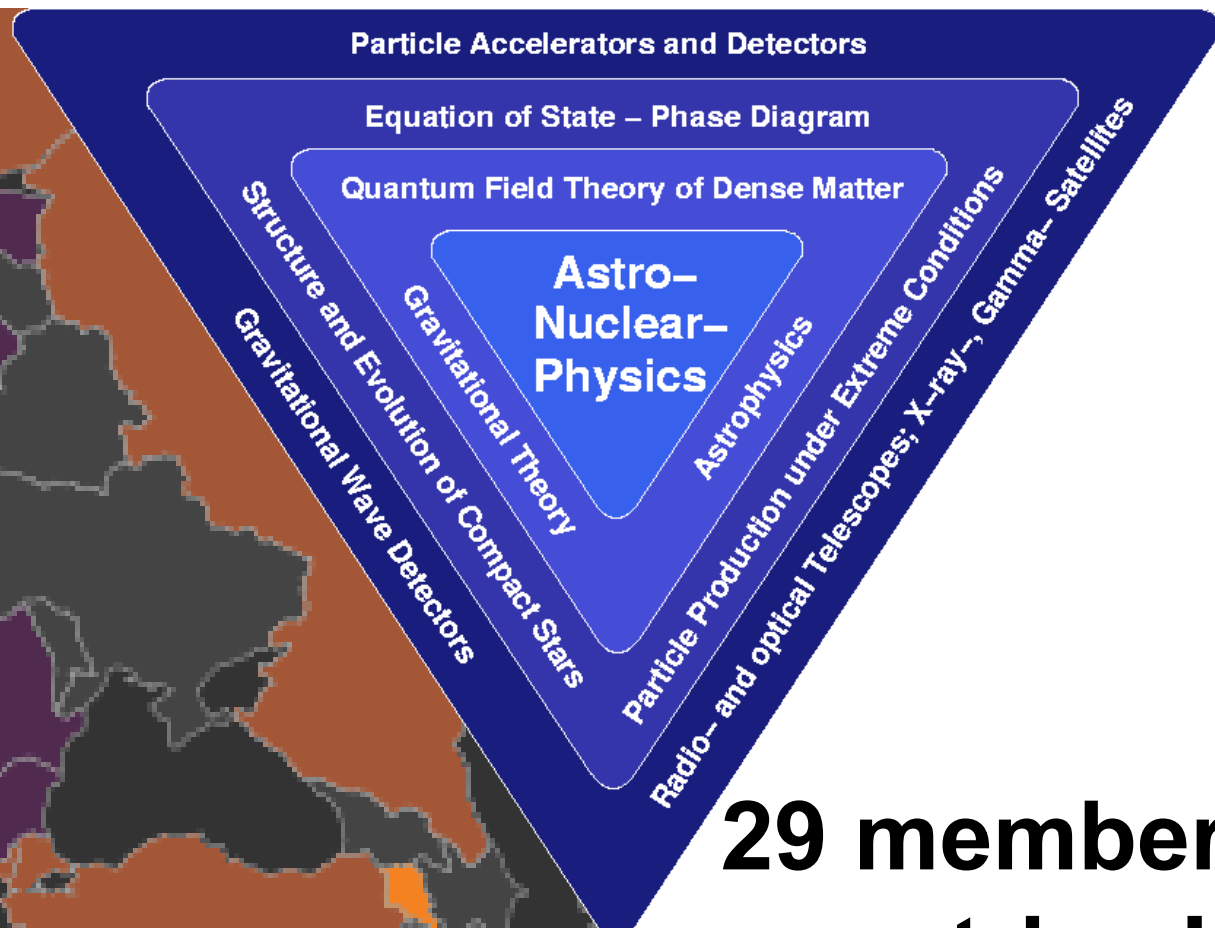
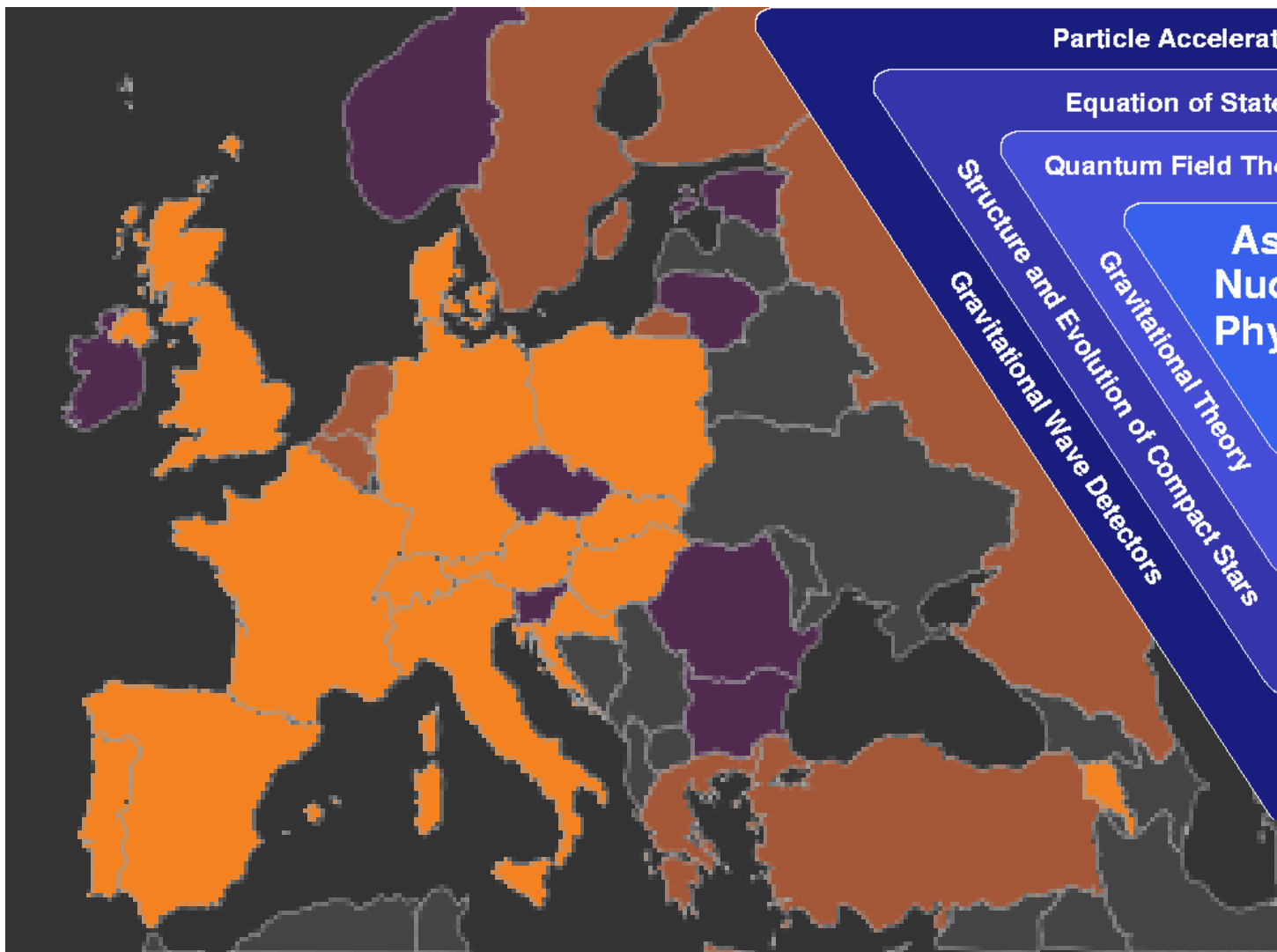
Main condition: stiff hadronic & stiff quark matter EoS with strong phase transition (PT)

Existence of HMTs & TMTs can be verified, e.g., by precise pulsar mass and radius measurements (and good luck) → Indicator for strong PT !!

Extremely interesting scenarios possible for dynamical evolution of isolated (spin-down and accretion) and binary (NS-NS merger) compact stars; GW170817 could be the inspiral of a neutron star – hybrid star binary !

Critical endpoint search in the QCD phase diagram with Heavy-Ion Collisions goes well together with Compact Star Astrophysics



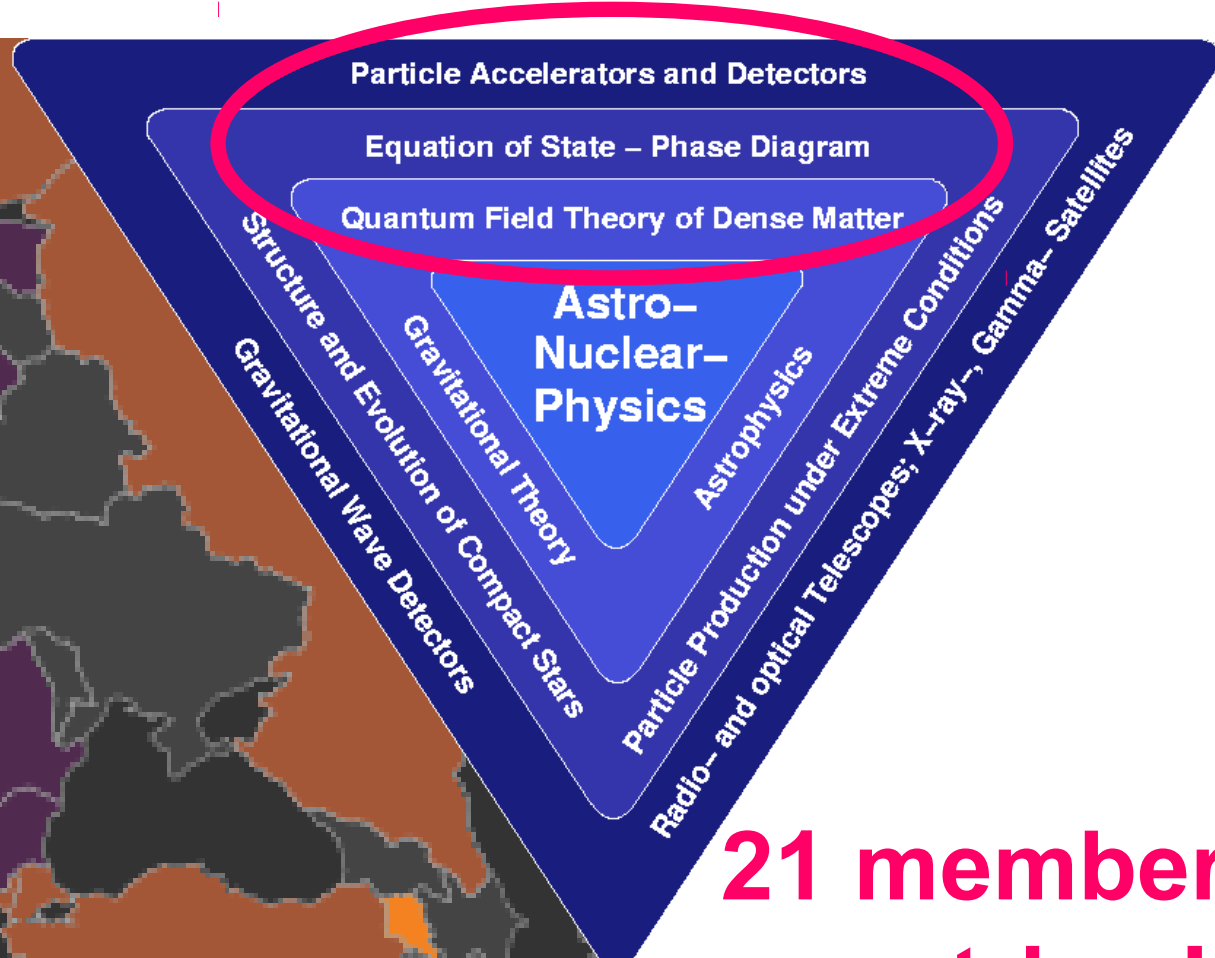
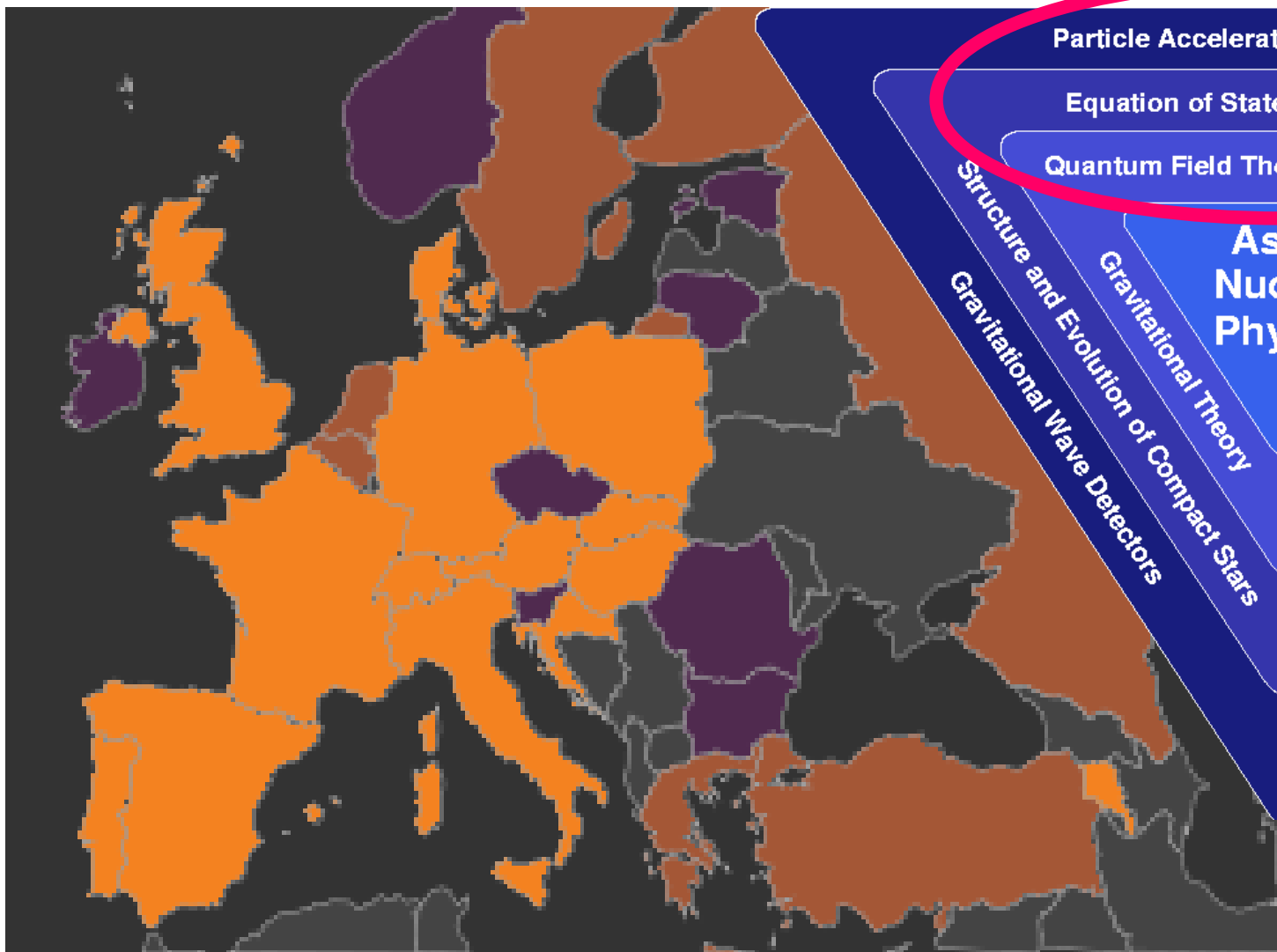


**29 member
countries !!
(MP1304)**

New



Kick-off: Brussels, November 25, 2013



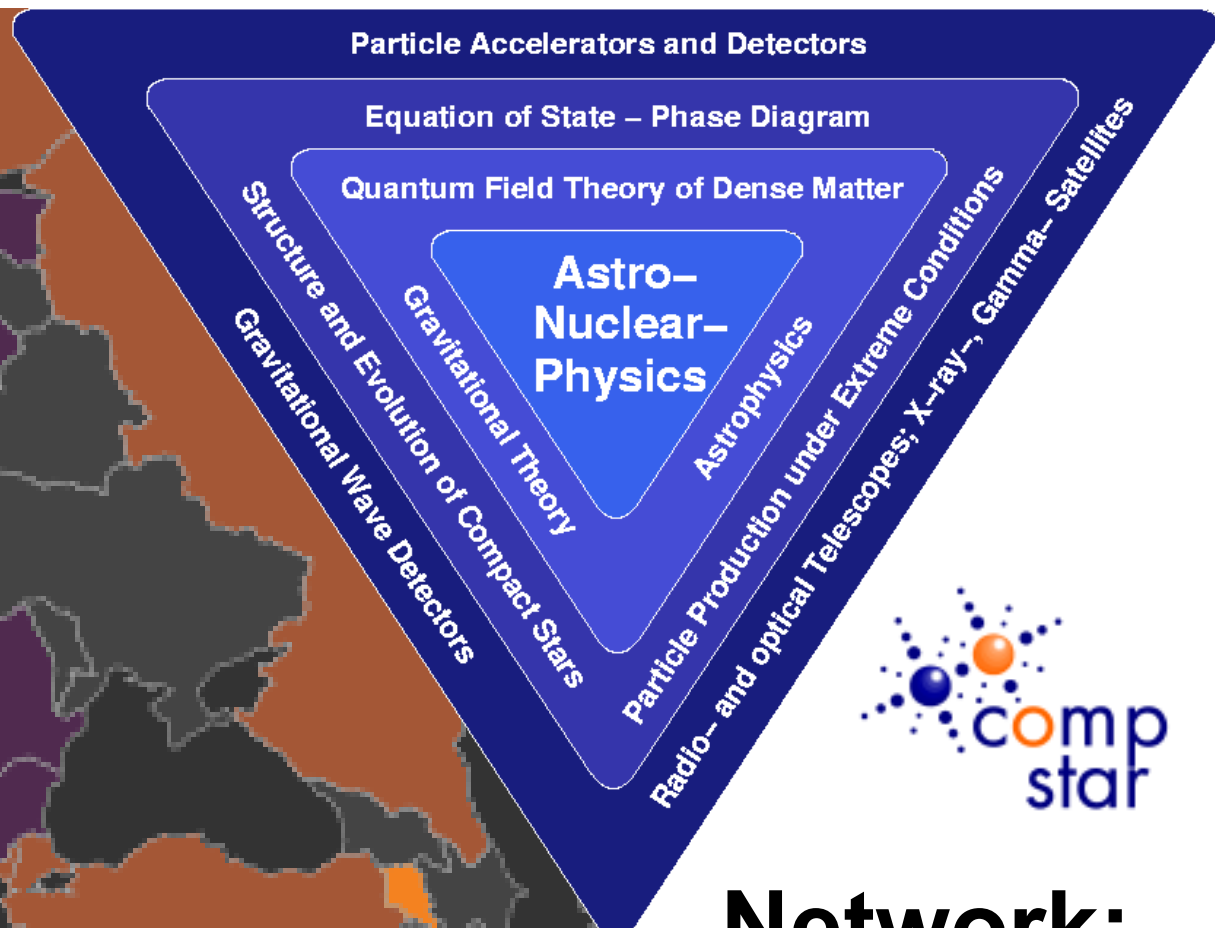
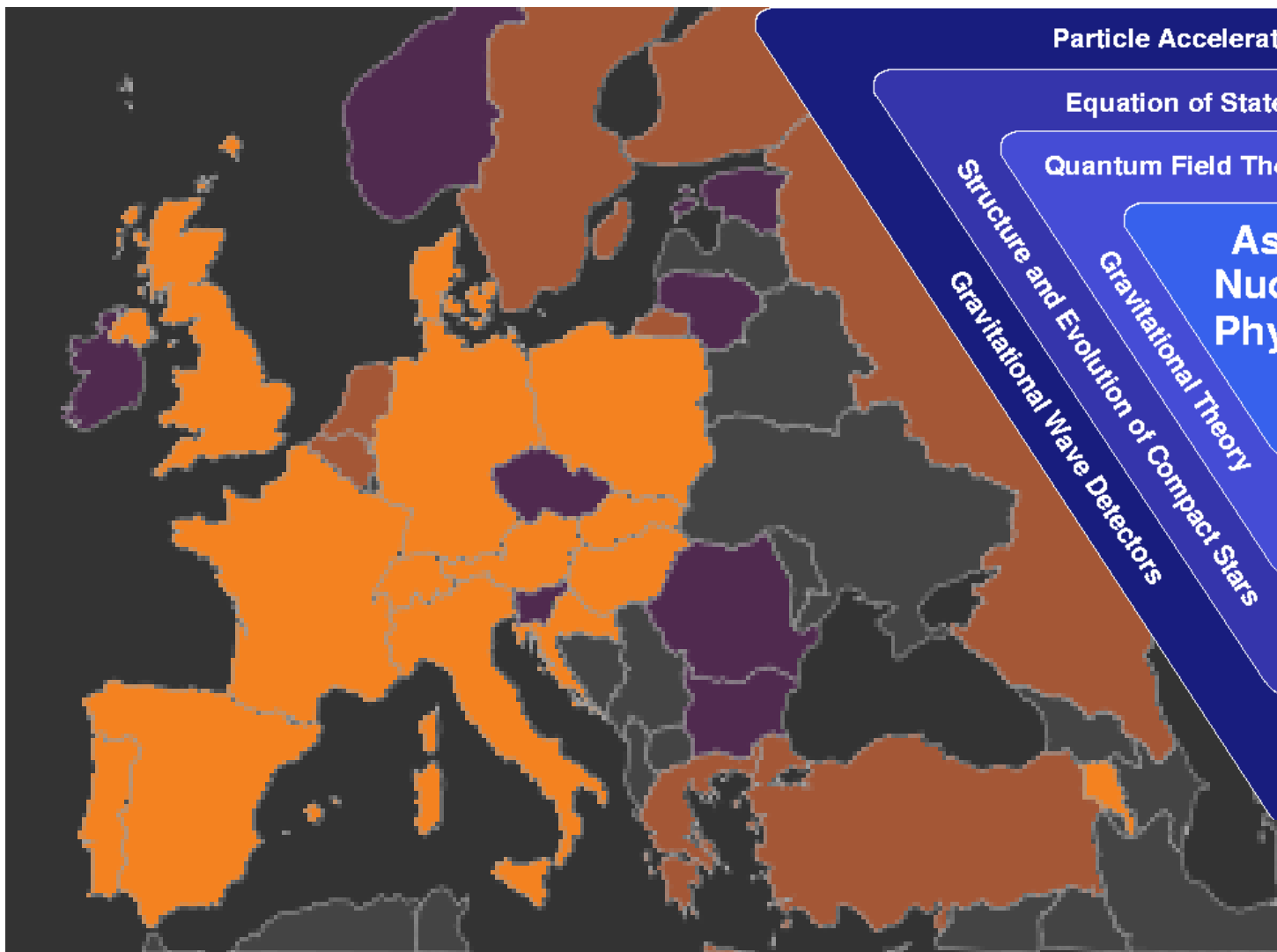
**21 member countries !
(CA15213)**

“**T**heory of **H**Ot Matter in **R**elativistic Heavy-Ion Collisions”

New: THOR!



Kick-off: Brussels, October 17, 2016



**Network:
CA16214**

**Newest:
PHAROS**



http://www.cost.eu/COST_Actions/ca/CA16214

Kick-off: Brussels, 22.11. 2017



International Conference “Critical Point and Onset of Deconfinement”
University of Wroclaw, May 29 – June 4, 2016



Recognized by European Physical Society

Hadrons and Nuclei

Topical Issue on Exploring Strongly Interacting Matter at High Densities - NICA White Paper

edited by David Blaschke, Jörg Aichelin, Elena Bratkovskaya, Volker Friese, Marek Gazdzicki, Jürgen Randrup, Oleg Rogachevsky, Oleg Teryaev, Viacheslav Toneev



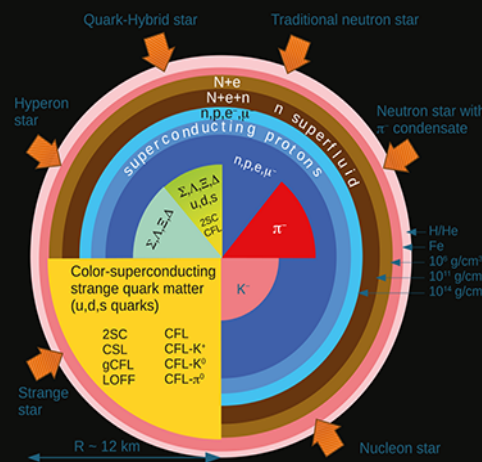
From: Three stages of the NICA accelerator complex by V. D. Kekelidze et al.



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Hadrons and Nuclei

Inside: Topical Issue on Exotic Matter in Neutron Stars edited by David Blaschke, Jürgen Schaffner-Bielich and Hans-Josef Schulze



From: Neutron star interiors: Theory and reality by J.R. Stone (left)

Phenomenological neutron star equations of state: 3-window modeling of QCD matter by T. Kojo (right)

